



Durham E-Theses

Vibration of shells with application to hollow blading

Ucmaklioglu, Mehmet

How to cite:

Ucmaklioglu, Mehmet (1978) *Vibration of shells with application to hollow blading*, Durham theses, Durham University. Available at Durham E-Theses Online: <http://etheses.dur.ac.uk/8400/>

Use policy

The full-text may be used and/or reproduced, and given to third parties in any format or medium, without prior permission or charge, for personal research or study, educational, or not-for-profit purposes provided that:

- a full bibliographic reference is made to the original source
- a [link](#) is made to the metadata record in Durham E-Theses
- the full-text is not changed in any way

The full-text must not be sold in any format or medium without the formal permission of the copyright holders.

Please consult the [full Durham E-Theses policy](#) for further details.

VIBRATION OF SHELLS
WITH APPLICATION TO
HOLLOW BLADING

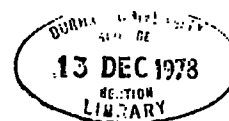
by

MEHMET UÇMAKLIOĞLU, B.Sc.

A thesis submitted for the degree of
Doctor of Philosophy to the
Department of Engineering Science,
University of Durham.

September 1978

The copyright of this thesis rests with the author.
No quotation from it should be published without
his prior written consent and information derived
from it should be acknowledged.



SEVGİLİ ANNEME VE BABAMA

(Dedicated to my Parents)

ACKNOWLEDGEMENTS

I am extremely grateful to my supervisor, Dr. P.A.T.Gill for his guidance and assistance throughout the work.

I would like to thank too Dr. J.M. Wilson, and Dr. D.L. Collins for their help and suggestions about the numerical analysis.

Thanks are also due to the Technicians in the Department of Engineering Science, to the member of staff of the Computer Unit and to the staff of the Science Site Library.

Finally, I wish to express my thanks to the Union Carbide for supplying the financial support for this work.

ABSTRACT

The finite element method was applied to the natural frequency analysis of arbitrary shell structures. A computer program based on the isoparametric thick-shell element was developed.

The program was tested against several plate and shell problems. The results were compared with the experimental and numerical results reported by other researchers with excellent agreement in most cases and fair agreement in others.

An oval cross-section hollow blade was analysed in detail both numerically and experimentally. The experimental model could not match the design geometry due to manufacturing difficulties. The numerical analysis was first performed on the nominal geometry which lead to a regular set of modes. Later, the numerical model was corrected to match the experimental model, and satisfactory agreement was obtained between the results for the lower modes of vibration.

Other topics which could be studied as an extension of this work were pointed out, and some excercises were performed on them without given any experimental verification. Finally a hollow turbine blade was analysed and very good results were obtained.

CONTENTS

	<u>Page</u>
<u>CHAPTER 1. INTRODUCTION</u>	1
<u>CHAPTER 2. THEORY AND NUMERICAL FORMULATION</u>	6
2.1. Theory	7
2.1.1. Introduction	7
2.1.2. Formulation	7
2.2. Isoparametric Shell Element	10
2.2.1. Introduction	10
2.2.2. Element Formulation	10
2.3. Shape Functions	13
2.3.1. Introduction	13
2.3.2. Shape Functions for Cubic- Quadratic Element.	14
2.4. Reduced Integration	17
2.5. Repetitive use of Element Matrices	18
2.6. Eigenvalue Economization	19
<u>CHAPTER 3. TEST EXAMPLES</u>	21
3.1. Uniform Cantilever Plates	21
3.2. Tapered Rectangular Plate	22
3.3. Pretwisted Cantilever Blading	25
3.4. Curved Fan Blade	30
3.5. Circular Cylindrical Shell	34
3.6. Curved Cantilever Beam-Pipe Segments	34
3.7. Conclusion	40

	<u>Page</u>
<u>CHAPTER 4. EXPERIMENTAL METHOD</u>	42
4.1. Determination of the Actual Geometry of the Shell	42
4.2. Apparatus and Procedure	43
<u>CHAPTER 5. OVAL CROSS-SECTION HOLLOW BLADE</u>	46
5.1. Geometry of the blade	46
5.2. Manufacturing the Experimental Model	48
5.2.1. Solder Assembly	48
5.2.2. Weld Assembly	49
5.3. Analysis of the blade Assembled using Solder	51
5.4. Effect of the Root Fixing on the Natural Frequencies	53
5.5. Weld Assembled Blade	54
5.5.1. Finite Element Analysis using Ideal Geometry	54
5.5.2. Finite Element Analysis - Real Geometry	56
5.5.3. Experimental Analysis	58
5.5.4. Comparison of the Results	58
5.6. Effect of the Material Properties and the choice of Degrees of Freedoms on the Natural Frequencies and the Mode Shapes	60
5.7. Discussions of Results	62
<u>CHAPTER 6. FUTURE APPLICATIONS</u>	66
6.1. Sharp Corner Connections	66
6.1.1. Cantilever Plate	67
6.1.2. V-Shape Cross-Section Cantilever	71
6.1.3. Aerofoil Cross-Section Hollow Blading	74
6.2. Stiffeners	78
6.3. Effect of Pretwist	78
6.3.1. Pretwisted Oval Blading	80
6.3.2. Pretwisted Aerofoil Blading	80
6.3.3. Pretwisted Cylinder	81
6.4. A Real Turbine Blade	87

	<u>Page</u>
<u>CHAPTER 7. CONCLUSIONS</u>	94
<u>REFERENCES</u>	96
<u>APPENDIX 1</u>	102
<u>APPENDIX 2</u>	109
<u>APPENDIX 3</u>	118
<u>APPENDIX 4</u>	142
<u>APPENDIX 5</u>	149

CHAPTER 1

1. INTRODUCTION

As a result of the technological developments in civil, mechanical and aeronautical engineering, shell structures gained a popularity in the 20th century. Since many of such structures are subject to dynamic loading, it is important to have some indication of their vibrational characteristics.

The number of exact solutions which satisfy the governing equations and the boundary conditions of shell structures are very limited. Hence, engineers are forced to use some approximate numerical methods for their analysis. Fortunately, the present state of high speed, large storage computers provide the necessary facilities for the development of efficient numerical methods.

In the literature, different methods have been employed for the analysis of shell structures. Warburton outlines Rayleigh-Ritz, finite element, finite difference and numerical integration methods in his review of "Dynamics of Shells" (55). Petyt (48) analysed a singly curved rectangular plate using Rayleigh-Ritz, extended Rayleigh-Ritz, finite element, and Kantorovich methods. Amongst the others, finite element method occupies an outstanding position. The superiority of the method comes from its applicability to arbitrary shapes, different loading and support conditions and many other aspects of practical design.

In general, considerable work has been done on the vibrational analysis of axisymmetric shells (37), and some on the circular cylindrical shell segments. The number of publications available for more complicated shapes are very few. Kurt and Boyd (38) for instance, studied a non-circular cylindrical shell segment by



using power series method. Applicability of the method is restricted by the curvature being expressed as a power series, and with the boundary conditions being simply supported at two opposite edges. McDaniel and Logan (41) studied the panels with exponential curvature using transfer matrix method. Srinivasan and Bobby(52) used a matrix method to analyse the vibrations of clamped non-circular cylindrical shell segments. Cheung and Cheung (14) applied the finite strip method to the analysis of non-circular cylindrical panels. In all these methods either the geometry or the boundary conditions impose some restrictions on their range of application.

Early application of the finite element method to the shell structures was limited to the axisymmetric shells. Finite element analysis of non-axisymmetric shells can be reviewed under three groups. Namely, flat plate elements, curved elements, and three dimensional elements.

Flat plate elements (15,21,58) have the common disadvantage of uncoupled representation of membrane and bending behaviours. Also, the representation of curved geometry using flat elements requires a fine mesh in order to achieve a reasonable degree of accuracy.

Curved shell elements, based on the assumptions of different thin-shell theories, were developed as an attempt to overcome the disadvantages of flat plate elements. Generally, it is possible to divide them into two sub-groups as cylindrical shell elements (12,43), and doubly curved shallow shells (17,39,53). Many of these elements, together with some flat ones are reviewed and compared in references (16,18,25,26).

Three dimensional brick elements were first modified and used for the analysis of thick-shell structures by Ahmad and others (1,3,4,58). The formulation of thick shell element was more attractive than the conventional shell theories because of its conceptual simplicity. Also, it was capable of reproducing shear deformations in the element, and representing any arbitrary geometry. Further modifications on the element (61,47) made it applicable to thin shell structures as well as thick shells. After these modifications it became one of the most accurate and popular element for the analysis of shells.

Most of the shell elements were initially developed for static problems. Dynamic applications of these elements are not as common as the static applications in literature. Specifically for the vibrational analysis of non-axisymmetric shells, there is very little published literature.

A valuable experimental and numerical contribution to the vibrational analysis of shells was done by Olson and Lindberg. They used a cylindrical shell element (43,44) to analyse a curved fan blade. Later, Lindberg et al. (39,45) solved the same problem using a doubly curved triangular element. Their fan blade became a popular test example for many others (2,9,11,30,40,42).

Neale developed a hybrid cylindrical shell element (42) and tested it against several problems available in the literature. His element was non-conforming. He reported that for coarse meshes it is superior to many others.

Martins and Oven (40) used the semiloof element for the analysis of thin arbitrary shells. The element gave very good results for several problems they studied. Unfortunately, the element can not accommodate lateral shear, and is applicable only to thin shell situations.

Ahmad's shell element has been applied to thick shell vibration problems by its original developers (2). The element gave good results in application to thick shell vibration problems, but they pointed out that it was over stiff for the representation of thin shells.

Hofmeister and Evensen (30) have used both the modified and unmodified elements, and the twelve-node cubic element to solve several plate and shell vibration problems. They too, reported that unmodified eight-node element was too stiff, except for simple vibration modes, or large meshes. Both modified eight-node element and the twelve-node element performed very well in most of the cases.

Bossak and Zienkiewicz (9) applied the modified version of the element to thin shell and pretwisted beam problems. They made a comparison between the results of the original and the modified elements indicating the improvement achieved by reducing the order of integration.

In the present study, a complete computer program using eight-node quadratic and ten-node cubic-quadratic elements has been developed. The program was checked against several plate and shell problems. The results were compared with the other numerical and experimental results with excellent agreement.

The study continued with the analysis of an oval cross-section shell, representing a hollow blading. Although the comparison of the experimental and numerical results were very difficult for some complicated modes which were highly effected by the imperfect geometry of the experimental model, very good agreement was observed for simpler mode shapes. Later the program was used on pretwisted oval and aerofoil cross-section bladings and a hollow turbine blade. In general the results indicate that the method used is very efficient for the vibrational analysis of arbitrary shell structures.

CHAPTER 2

2. THEORY AND NUMERICAL FORMULATION

In this chapter the theoretical basis to the computer program which has been written for the numerical analysis in this thesis is reviewed. In addition to the general theory some points to increase the accuracy of the element, or to reduce the computational cost of the program are also mentioned.

Explanations on the program are given in Appendix 1. Appendix 2 contains a detailed explanation for the preparation of input data. A complete listing of the program is included in Appendix 5.

2.1 Theory

2.1.1. Introduction

The solution of any structural problem by finite element displacement method follows more or less the same procedure (58,19,10). First the continuum is divided into a number of "finite elements", and these elements are considered to be interconnected at a discrete number of nodes. Displacements of these nodal points are the basic unknowns of the problem. Then, the element characteristics such as the stiffness and mass matrices are evaluated. Their assembly into the system matrices is followed by a solution procedure.

2.1.2. Formulation

The displacement at any point within a typical element "e" can be defined as:

$$\{f\} = [N] \{\delta\}_e \quad (2.1.1)$$

Where $[N]$ contains the shape functions which will be discussed in section 2.3 and $\{\delta\}_e$ is the vector containing the displacements of the nodes of element "e".

Strains within the element are defined as:

$$\{\epsilon\} = [B] \{\delta\}_e \quad (2.1.2)$$

Where $[B]$ can be obtained from (2.1.1) using strain-displacement relationship.

Stresses, in the elastic range, are related to the strains:

$$\{\sigma\} = [D] \{\epsilon\} \quad (2.1.3)$$

In which $[D]$ is the elasticity matrix which contains the material properties.

If $\{F\}_e$ is the vector of nodal forces equivalent to the boundary stresses and the distributed loads on the

element, and $\{p\}$ is the vector of distributed loads and forces acting on a unit volume of material within the element; principle of virtual work can be applied to equate the external and internal work done due to a virtual nodal displacement $d\{\delta\}_e$. The work done by the nodal forces is:

$$(d\{\delta\}_e)^T \{F\}_e \quad (2.1.4)$$

The internal work per unit volume done by the stresses and the distributed forces is:

$$d\{\epsilon\}^T \{\sigma\} - d\{f\}^T \{p\} \quad (2.1.5)$$

or by substituting (2.1.1) and (2.1.2)

$$(d\{\delta\}_e)^T ([B]^T \{\sigma\} - [N]^T \{p\}) \quad (2.1.6)$$

Now, the external and total internal work done can be equated:

$$(d\{\delta\}_e)^T \{F\}_e = (d\{\delta\}_e)^T \left(\int [B]^T \{\sigma\} dV - \int [N]^T \{p\} dV \right) \quad (2.1.7)$$

Substitution of equation (2.1.3), followed by the substitution of (2.1.2) gives:

$$\{F\}_e = \left(\int [B]^T [D] [B] dV \right) \{\delta\}_e - \int [N]^T \{p\} dV \quad (2.1.8)$$

In which the first term contains the element stiffness matrix:

$$[k]_e = \int [B]^T [D] [B] dV \quad (2.1.9)$$

and the second term is the nodal forces

$$[F]_e^p = - \int [N]^T \{p\} dV \quad (2.1.10)$$

For dynamic problems $\{p\}$ includes inertia and damping forces in it

$$\{p\} = \{\bar{p}\} - \rho \frac{\partial^2 \{f\}}{\partial t^2} - \mu \frac{\partial \{f\}}{\partial t} \quad (2.1.11)$$

where $\{\bar{p}\}$ represents the distributed loads, $-\rho \frac{\partial^2 \{f\}}{\partial t^2}$ represents the inertia force with ρ being the mass per unit volume, and $-\mu \frac{\partial \{f\}}{\partial t}$ represents a linear viscous damping

effect, where μ is a constant which characterizes the damping mechanism.

Substituting (2.1.11) into (2.1.10) and replacing $\{f\}$ with its equivalent in (2.1.1) gives

$$\{F\}_e^p = \{\bar{F}\}_e^{\bar{p}} + \int [N]^T \rho [N] \frac{\partial^2}{\partial t^2} \{\delta\}_e dV + \int [N]^T \mu [N] \frac{\partial}{\partial t} \{\delta\}_e dV \quad (2.1.12)$$

which contains the element mass matrix as:

$$[m]_e = \int [N]^T \rho [N] dV \quad (2.1.13)$$

and element damping matrix as:

$$[c]_e = \int [N]^T \mu [N] dV \quad (2.1.14)$$

Substituting (2.1.12) into (2.1.8) and assembling the element matrices into the system matrices, the general equation for discrete structures is obtained.

$$[K] \{\delta\} + [C] \frac{\partial}{\partial t} \{\delta\} + [M] \frac{\partial^2}{\partial t^2} \{\delta\} + [F] = 0 \quad (2.1.15)$$

Where $\{\delta\}$ lists the nodal displacements (degrees of freedoms), $[F]$ contains the external forces, specified loads and initial stresses, $[K]$, $[C]$, and $[M]$ are the system stiffness, damping and mass matrices respectively.

For the natural frequency analyses equation (2.1.15) reduces to:

$$[K] \{\delta\} + [M] \frac{\partial^2}{\partial t^2} \{\delta\} = 0 \quad (2.1.16)$$

Solution to this equation may be expressed as:

$$\{\delta\} = \{\delta_0\} e^{i\omega t} \quad (2.1.17)$$

which, when substituted into (2.1.16) leads to the eigenvalue problem

$$([K] - \omega^2 [M]) \{\delta_0\} = 0 \quad (2.1.18)$$

In solution, ω 's will be the natural angular frequencies, and

{ δ .} will contain the eigenvectors describing the mode shapes of vibration.

2.2. Isoparametric Shell Element

2.2.1. Introduction

As it was stated in Chapter 1, Ahmad's isoparametric shell element was chosen for the numerical analysis of this work. This element differs from a three dimensional brick element with the assumption that the normals to the midsurface remain straight after deformation and that the strains normal to the midsurface are negligible. It also differs from the elements of thin shell theory with the assumption that the normals to the midsurface are allowed to become inclined after deformation. This assumption permits the element to experience shear deformations.

The two typical elements used in this analysis are shown in figure 2.1. The element in figure 2.1 (a) is known as 8-node isoparametric shell element which assumes a parabolic variation along the edges. The element in figure 2.1(b) is derived by combining a parabolic and a cubic element, for a better representation of sharp changes of curvature along one edge.

2.2.2. Element Formulation

The formulation of the isoparametric shell element (1,3,4, 58) follows the phases given in section 2.1.2.

The geometry of each element is prescribed by top and bottom pairs of coordinates of the nodes, and the shape functions used. ξ, η, ζ are the curvilinear coordinates of the

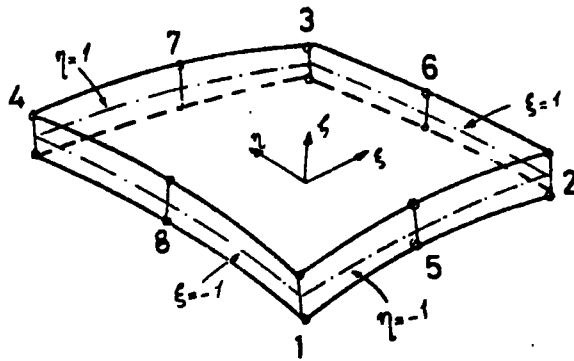


Figure 2.1 (a) : Isoparametric Shell Element with Quadratic Variation

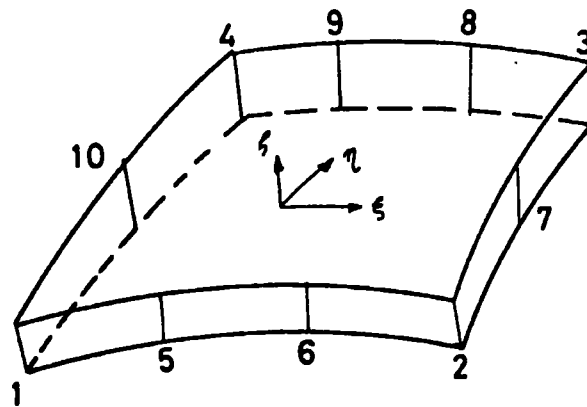


Figure 2.1 (b) : Quadratic - Cubic Shell Element

shell element, and each varies between -1 and 1 on the respective faces of the element. (See figure 2.1 (a)). Global cartesian coordinates of any point of the element are related to the curvilinear coordinates as:

$$\begin{Bmatrix} x \\ y \\ z \end{Bmatrix} = \sum N_i \begin{Bmatrix} x_i \\ y_i \\ z_i \end{Bmatrix}_{\text{mid.}} + \sum N_i \frac{\xi}{2} \underline{V}_{3i} \quad (2.2.1)$$

where N_i is the shape function corresponding to node i and \underline{V}_{3i} is the thickness vector at node i .

Displacements through the element is defined in terms of the midsurface node displacements u_i, v_i, w_i , and two rotations of the \underline{V}_{3i} vector about orthogonal directions normal to it. Two such directions are given by orthonormal vectors \underline{V}_{1i} and \underline{V}_{2i} with the corresponding rotations β_i and α_i .

Displacements within the element is related to the nodal displacements as:

$$\begin{Bmatrix} u \\ v \\ w \end{Bmatrix} = \sum N_i \begin{Bmatrix} u_i \\ v_i \\ w_i \end{Bmatrix} + \sum N_i \xi \frac{t_i}{2} [\underline{V}_{1i}, \underline{V}_{2i}] \begin{Bmatrix} \alpha_i \\ \beta_i \end{Bmatrix} \quad (2.2.2)$$

where t_i is the magnitude of the thickness vector \underline{V}_{3i} .

The components of strains and stresses are defined in the local cartesian coordinates x', y', z' with z' being normal to

$\xi = \text{constant}$ surface. Three dimensional strain relationship is given as:

$$\{\epsilon'\} = \begin{Bmatrix} \epsilon_{x'} \\ \epsilon_{y'} \\ \gamma_{x'y'} \\ \gamma_{x'z'} \\ \gamma_{y'z'} \end{Bmatrix} = \begin{Bmatrix} \frac{\partial u'}{\partial x'} \\ \frac{\partial v'}{\partial y'} \\ \frac{\partial u'}{\partial y'} + \frac{\partial v'}{\partial x'} \\ \frac{\partial w'}{\partial x'} + \frac{\partial u'}{\partial z'} \\ \frac{\partial w'}{\partial y'} + \frac{\partial v'}{\partial z'} \end{Bmatrix} \quad (2.2.3)$$

The local derivatives of the local displacements are given by

$$\begin{bmatrix} \frac{\partial u'}{\partial x'} & \frac{\partial v'}{\partial x'} & \frac{\partial w'}{\partial x'} \\ \frac{\partial u'}{\partial y'} & \frac{\partial v'}{\partial y'} & \frac{\partial w'}{\partial y'} \\ \frac{\partial u'}{\partial z'} & \frac{\partial v'}{\partial z'} & \frac{\partial w'}{\partial z'} \end{bmatrix} = \Theta^T J^{-1} \begin{bmatrix} \frac{\partial u}{\partial \xi} & \frac{\partial v}{\partial \xi} & \frac{\partial w}{\partial \xi} \\ \frac{\partial u}{\partial \eta} & \frac{\partial v}{\partial \eta} & \frac{\partial w}{\partial \eta} \\ \frac{\partial u}{\partial \zeta} & \frac{\partial v}{\partial \zeta} & \frac{\partial w}{\partial \zeta} \end{bmatrix} \Theta \quad (2.2.4)$$

where, Θ is the orthogonal transformation matrix between the local and global system of coordinates, and J is the Jacobian matrix.

Since matrix $[B]$ can be obtained from equation (2.2.4) and matrix $[N]$ is given by the equation (2.2.2), the stiffness and mass matrices of equation (2.1.9) and (2.1.13) can now be evaluated.

2.3. Shape Functions

2.3.1. Introduction

As was mentioned previously, displacements throughout the element are defined as a function of the nodal displacements. In the isoparametric formulation the same functions, namely the shape functions, are used to describe both the geometry and the displacements (58). Conveniently, these functions can be chosen as polynomials.

For the convergence of finite element analysis, there are two criteria to be satisfied (58,19,60):

- a) Continuity of the displacements between the elements should exist.
- b) Elements should be able to reproduce any required state of constant strain. This criteria includes the rigid body

displacements as a special case.

In selecting the polynomials for the shape functions, these two criteria must be taken into account. An additional desirable property of the shape functions is the geometric invariance (22,23,19).

To satisfy the first condition the displacement function along the boundary of the two adjacent elements must only be influenced by the nodes on this boundary, and the order of the polynomial must be uniquely determined by the number of these nodes (60,24). Second criteria is automatically satisfied for the isoparametric shape functions (24). To achieve the geometric invariance it is necessary to choose symmetric polynomials as the shape functions (19).

By definition, shape functions take a unit value at a preferred node and zero at all other nodes. Suitable shape functions, satisfying all the conditions mentioned above are available in the literature (58,59,60,24). Following the tradition, and the recommendations of references (59,54), the Serendipity family of shape functions were used in this work. In the following section, the derivation of shape functions of this family, for a rectangular element with different number of nodes along the parallel sides is demonstrated.

2.3.2. Shape Functions for Cubic-Quadratic Element

For some practical purposes, it is desirable to have elements with varying number of nodes along different sides. The shape functions for such an element, which is shown in

figure 2.2 and will be used in chapters 3 and 5, derived following the procedure given in (59,54).

Shape functions for the midside nodes of this element are the same as a cubic element in ξ and as a quadratic element in η directions.

To derive the shape functions for the corner nodes, one can start from the linear function shown in figure 2.3(a). To obtain the cubic variation of figure 2.3(e), first the variation in figure 2.3(b) is to be subtracted from figure 2.3(a) to give figure 2.3(c), then figure 2.3(d) is to be subtracted from figure 2.3(c) to give the final form in figure 2.3(e).

A similar process, of course, will be applied for the quadratic variation in η direction.

The curves of figures 2.3(b) and 2.3(d) are the cubic shape functions for the midside nodes, multiplied by a constant.

The shape function for node 1 of figure 2.2 for instance will be

$$N_1 = N_L - \frac{2}{3} N_5 - \frac{1}{3} N_6 - \frac{1}{2} N_{10} \quad (2.3.1)$$

where N_L is the linear shape function for node 1. Substitution of the corresponding functions of N_i gives the shape functions as:

$$N_1 = \frac{1}{32} (1-\xi)(1-\eta) [-8\eta - 9(1-\xi^2)] \quad (2.3.2)$$

or the general form for all the corner nodes:

$$N_i = \frac{1}{32} (1+\xi_0)(1+\eta_0) [8\eta_0 - 9(1-\xi^2)] \quad (2.3.3)$$

with $\xi_0 = \xi \xi_i$ and $\eta_0 = \eta \eta_i$

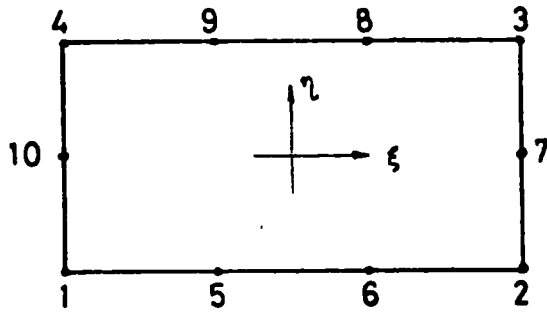


Figure 2.2 : Node numbering system for 10-node element.

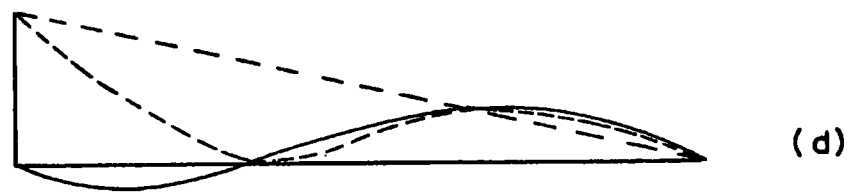
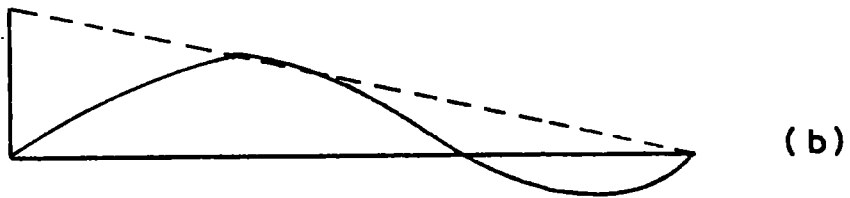
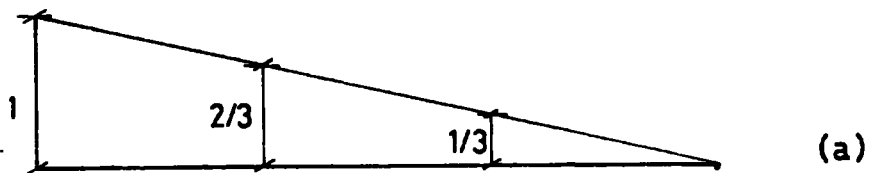


Figure 2.3 : Derivation of Shape Functions for corner nodes.

For the midside nodes the shape functions are:

$$N_i = \frac{1}{2} (1 + \xi_0) (1 - \eta^2) \quad \text{for} \quad \xi_i = \pm 1, \eta_i = 0 \quad (2.3.4)$$

$$N_i = \frac{9}{32} (1 + \eta_0) (1 - \xi^2) (1 + 9\xi_0) \quad \text{for} \quad \xi_i = \pm \frac{1}{3}, \eta_i = \pm 1 \quad (2.3.5)$$

2.4. Reduced Integration

In its original form, the thick shell element was too stiff in representing bending deformations in thin shell and plate applications. Zienkiewicz et al. (61), and Pawsey and Clough (47) had noticed that the over stiffness of the element was due to the displacement function imposing unrealistic restrictions upon the modes of deformation of the element. Following a procedure similar to the one that Doherty et al (20) employed, that is, by reducing the order of integration, it was possible to relax the over stiffness of the element. Accuracy of the reduced numerical integration in finite elements is discussed by Irons (33,34,35,36). The point he emphasises is that the convergence is guaranteed, provided that the determinant of Jacobian does not change sign in the domain, and that the volume of the element is calculated accurately.

To modify the element, Zienkiewicz et al. proposed a uniform low order integration on all stress components, whereas Pawsey and Clough used a selective integration order for different stress components. Pawsey, in a later publication (46) agrees that the eight node quadratic element can be integrated using uniformly reduced integration points. The advantage of the uniform reduction over the selective reduction lies in the

fact that it does not require any additional effort in programming, and reduces the computational time due to the lesser number of integration points.

In both of the references (61,47), the improvement achieved by reducing the order of integration is demonstrated by several examples. Vibrational application of the modified element (using uniform reduction) is reported in (30,9). In all examples but one, the element performed very well. Only in the case of a turbine blade, was it found to be overflexible (30). The reason was the high curvature across the width of the blade. Since only one element was used across the width, it was not able to represent rapid local change of curvature.

In the present study , the uniform reduction of integration points is used to modify the element. In general $2 \times 2 \times 2$ integration points are used for the stiffness matrix, and $3 \times 3 \times 2$ for the mass matrix. The effect of using different integration points has been observed on the results of some of the test examples of chapter 3. A survey was carried out to decide the best choice of order of integration for the elements having a high curvature. For these elements it is found out that unless a refined mesh is used, 3 integration points are needed along the curved edge. The comparison of results using different integration orders is given in the next chapter.

2.5. Repetitive use of Element Matrices

Using a regular mesh, makes it easy to produce the input data, and also helps to reduce the computational time. Element matrices, if they were calculated in a local coordinate system

having the same orientation with respect to the element geometry, would be the same for similar elements. Once the matrices for one element are evaluated, the same matrices can be used for similar elements, provided that the necessary coordinate transformation is performed. If the similar elements have the same orientation with respect to the global system element matrices need not be transformed but could be used repetitively. If the orientation of similar elements is different with respect to the global coordinate system transformation is only required for the entries of the matrices corresponding to u , v , w degrees of freedoms, since α and β are the rotations of the thickness vector, and by definition, they are independent of the global coordinate system.

2.6. Eigenvalue Economization

The time required to solve the eigenvalue problems increases rapidly with the size of the matrices involved. Frequently, finite element idealization of a structure yields a representation with several hundreds of degrees of freedoms. The eigenvalue solution of a matrix of this size would require a great deal of computer time and storage which may be impossible to supply. Fortunately, these requirements may be reduced by a careful selection of a reduced number of degrees of freedoms which is sufficient to give acceptable results in the solution of the dynamic problems. This can be explained by Rayleigh's principle that a first order error in modal shape causes only a second order error in frequency.

Elimination of unwanted variables have been studied by several researchers (31,32,28,5,50,6,29). The method due to Irons (32,5) has been used in this study, because of its advantage of performing the elimination process during assembly.

The unwanted variables which are called slaves are chosen amongst the degrees of freedoms which have the least contribution to the strain energy. When sth degree of freedom is eliminated from the stiffness matrix $[K]$, the new entries of the reduced stiffness matrix $[K]^*$ are given by the formula

$$K^*_{ij} = K_{ij} - K_{is} (K_{js}/K_{ss}) \quad (2.6.1)$$

with row and column s being deleted. The corresponding operation on the mass matrix is:

$$M^*_{ij} = M_{ij} - M_{is} (K_{js}/K_{ss}) - M_{js} (K_{is}/K_{ss}) + M_{ss} (K_{is}/K_{ss})(K_{js}/K_{ss}) \quad (2.6.2)$$

The advantage of the method lies in the fact that the operations (2.6.1) and (2.6.2) can be applied to the M_{ij} and K_{ij} that are incompletely summed, as long as all of the M_{is} and K_{is} are completely summed, and as long as all contributions from later elements will eventually be added in.

CHAPTER 3.

3. TEST EXAMPLES

In order to see the performance of the computer program, it was tested against several plate and shell vibration problems. Either eight-node or ten-node elements were employed for the mathematical modelling of the structures. The effect of the number of Gauss points used in the numerical integration was investigated.

The results are given either in Hertz, or in dimensionless frequency parameter ϕ , which is defined as:

$$\phi = \frac{\omega}{(D/\rho t l^4)^{1/2}} \quad \text{where} \quad D = \frac{E t^3}{12(1-\nu^2)} \quad (3.1.)$$

in which ω is the natural angular frequency, ρ , E and ν are the density, modulus of elasticity and the Poisson's ratio respectively. t and l stand for the thickness and the length of the shell.

3.1 Uniform Cantilever Plates

The first example was a square cantilever plate. The natural frequencies of this plate were determined by using both modified and unmodified elements. The problem was solved for different length/thickness (l/t) ratios. The increase in the stiffness of the unmodified element with the increasing l/t ratio was observed. The dimensionless frequency parameter (ϕ), obtained by using 1×1 and 2×2 ^{uniform} meshes are listed on table together with the results reported in references (5) and (8).

3.1.1 The results of the unmodified element shows a rapid

convergence with the increasing number of elements. The modified element gives acceptable results even with one element for the first two frequencies. In addition to the modes shown on table 3.1., some in-plane modes were also observed. The frequencies corresponding to these modes were independent of the order of integration used.

A similar analysis was performed on a rectangular plate having a length/width ratio of 2. Since the results obtained were very similar to the square plate case, only the results of modified elements with 2×2 ^{uniform} mesh for three different l/t ratios are given on table 3.2, together with the results of references (5) and (49).

Because of the coarse mesh used, the frequencies corresponding to high mode shapes were overestimated on both tables.

3.2. Tapered Rectangular Plate

In this example, a varying thickness cantilever plate was analysed. The plate had a rectangular plan form with dimensions 127 x 63.5mm (5 x 2.5 inches). The cross section of it was an isosceles triangle with an apex angle of 2.4° . The frequencies were calculated for different l/w ratios. The width of the plate was varied by successively shaving it down on the thin side. This plate was first studied experimentally by Plunkett (49). He also calculated the fundamental frequencies using the beam theory.

In the present study, four eight-node modified elements were used to represent the structure. Dimensionless frequency parameters reported in reference (49), and the results of the

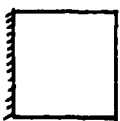
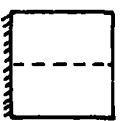
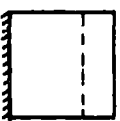
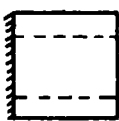
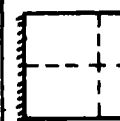
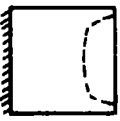
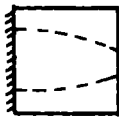
Number of elements	1/t	Order of Integration for stiffness Matrix					
Ref. 8		Ritz Meth.	3.49	8.55	21.44	27.46	31.17
Ref. 5		Finite Element	3.47	8.54	21.45	27.06	-
4	100/2	2x2x2	3.48	8.56	22.40	28.79	32.88
4	100/5	2x2x2	3.47	8.45	22.06	28.02	32.10
4	100/10	2x2x2	3.45	8.17	21.10	26.70	30.06
4	100/2	3x3x2	3.68	9.25	(34.20)	(41.24)	48.00
4	100/5	3x3x2	3.63	9.03	(29.51)	(35.77)	42.90
4	100/10	3x3x2	3.55	8.56	(25.30)	30.17	35.00
1	100/2	2x2x2	3.58	9.04			
1	100/3	2x2x2	3.58	9.02			
1	100/5	2x2x2	3.57	8.97			
1	100/7	2x2x2	3.57	8.89			
1	100/10	2x2x2	3.55	8.72			
1	100/2	3x3x2	4.45	11.43			
1	100/3	3x3x2	4.42	11.35			
1	100/5	3x3x2	4.35	11.12			
1	100/7	3x3x2	4.29	10.83			
1	100/10	3x3x2	4.18	10.30			

Table 3.1. Dimensionless frequency parameters of a square cantilever plate determined by using one and four 8-node elements.

(The mode shapes corresponding to the frequencies in paranthesis are shown at the bottom of the column).


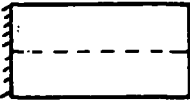


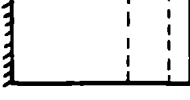
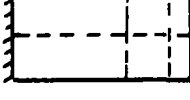
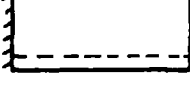

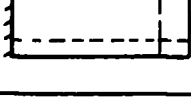
Mode Shape	Measured Ref.49	Calculated Ref.5	4 Elements 1/t = 100/2	4 Elements 1/t = 100/5	4 Elements 1/t=100/10
	3.50	3.44	3.47	3.46	3.45
	14.50	14.77	14.83	14.72	14.38
	21.70	21.50	22.74	22.66	22.38
	48.10	48.19	50.91	50.45	49.00
	60.50	60.54	75.93	75.08	72.36
	92.30	91.76	111.41	108.87	102.23
	92.80	92.78	99.16	98.12	95.04
	118.70	119.37	There are not enough nodes to identify this mode		
	125.1	124.23	137.11	133.12	126.06

Table 3.2 Dimensionless frequency parameter for a rectangular plate of $w/l = 1/2$

finite element analysis are given on table 3.3. Very good agreement was obtained for the simple mode shapes like $0/0$, $1/0$, $0/1$ (m/n , where m indicates the number of nodal lines perpendicular to the support, and n is the number parallel to it). Accurate determination of higher modes would require a finer mesh.

In general, the frequencies were underestimated for small l/w ratios, and overestimated for large ones. This is probably because of the changing element aspect ratio due to the decreasing width.

3.3. Pretwisted Cantilever Blading

The next example was a rectangular cross section pretwisted blading. The experimental values of the frequencies for this blading have been reported by Carnegie (13). He also gave the theoretical values for the fundamental frequencies, and for all torsional frequencies. The blade was 152.4 mm (6 inches) long, 25.4mm (1 inch) wide and 1.6129mm (0.0635 inches) thick.

The calculations in this study were performed for 0° , 30° , 60° and 90° of pretwist angle, using both eight-node and ten-node elements with 1×6 and 2×6 meshes. The results are listed on table 3.4.

In bending type of modes, the modified eight-node element performed very well with 1×6 mesh. But for torsional modes it was overflexible especially for large pretwist angles. In this case, it was possible to increase the accuracy by refining the mesh. Alternatively, by increasing the order of

Mode	l/w	2.00	2.22	2.86	4.00	6.67
0/0	Meas.	2.57	2.57	2.71	2.91	3.15
	Calc.	2.49	2.50	2.58	2.78	3.04
	F.E.	2.43	2.54	2.73	2.97	3.26
1/0	Meas.	11	11.6	15	22	37
	F.E.	10.63	11.50	15.50	23.48	40.79
0/1	Meas.	15.5	16	17	18	19.5
	F.E.	15.85	16.50	18.08	19.96	21.47
1/1	Meas.	30.7	33	49	68	112
	F.E.	30.55	36.34	49.67	75.25	127.42
0/2	Meas.	38	40	42.5	49	54
	F.E.	48.18	53.07	60.35	64.9	72.1

Table 3.3 Comparisons of the frequency parameters for the tapered rectangular plate.

integration points in the direction of the curvature, better results were obtained for torsional modes. Another alternative was to increase the number of nodal points in the direction of the curvature. This was achieved by using ten-node elements. The results obtained by using different number of integration points indicate that this element too, exaggerates the flexibility of the structure when a low order of integration is used. The best choice of integration order for this element seems to be $4 \times 2 \times 2$. Using 3 integration points along the length of the blade gives over stiff representation for both eight-node and ten-node elements.

Two different frequencies corresponding to some of the modes are the result of coupled bending-bending due to the pretwist. For instance, the second bending mode at 60° pretwist angle have two frequencies. One of them is around 260 Hz and second bending mode out of the plane of the root, couples with the second bending mode in the plane of the root. The other frequency which lies around 1320 Hz is the coupling between second bending out-of-plane and first bending in-plane modes. The value of this frequency, though, is more likely to correspond to the second value of the third bending frequency (1200 Hz) found experimentally. A similar discrepancy, in the identification of the modes, is seen at 90° pretwist angle. The second values corresponding to the fourth bending frequency almost coincide with the experimental fifth bending frequency.

MODE	Angle of Twist 0°				Angle of Twist 30°									
	CARNegie		FINITE ELEMENT		CARNegie		FINITE ELEMENT							
	Exper.	Theor.	8 Node Elements 1x6 2x2x2†	10 Node Elements 1x6 3x2x2 4x2x2	Exper.	Theor.	8 Node Elements 1x6 2x2x2	10 Node Elements 1x6 3x2x2 4x2x2						
BENDING	1	57.5	61.5	58.7	58.2	58.7	58.7	59	58	62	58.9	58.4	58.8	58.9
	2	362	386*	364	367	365	367	380	320	—	326	324	325	325
	3	1031	1081*	1026	1038	1031	1037	1128	1000	—	924 1084	923 1087	928 1092	929 1093
	4	1987	2119*	2061	2102	2075	2097	2423	1950	—	2041	2051	2072	2076
	5	3283	3502*	3613	3706	3645	3692	4620	3275 3212	—	3257	3587	3629	3640
TORSION	1	712	700	689	689	685	686	700	730	725	718	709	662	712
	2	2156	2100	2104	2106	2093	1937	2145	2200	2150	2192	2167	2024	2175
	3	3690	3600	3659	3670	3645	3383	3755	3800	3750	3810	3767	3525	3782
	4	5373	5380	5532	5564	5523	5160	5774	5500	5500	5745	5893	5342	5708

Table 3.4. Carnegie Pretwisted blading comparison
 The frequencies given in c/sec.
 * Mesh used.
 + Order of integration for the stiffness matrix
 X Calculated using beam theory.

MODE	ANGLE OF TWIST 60°										ANGLE OF TWIST 90°											
	CARNEGIE					FINITE ELEMENT					CARNEGIE					FINITE ELEMENT						
	Theor.		Exper.		8 Node Elements	10 Node Elem.		8 Node Elements		10 Node Elements		Theor.		Exper.		8 Node Elements	10 Node Elements		8 Node Elements		10 Node Elements	
	1 x 6 2x2x2	3x2x2	1 x 6 3x2x2	2 x 6 2x2x2		1 x 6 3x2x2	1 x 6 4x2x2	1 x 6 3x2x2	2 x 6 2x2x2	1 x 6 3x2x2	1 x 6 3x3x2	1 x 6 2x2x2	1 x 6 2x2x2	1 x 6 2x2x2	1 x 6 3x2x2		1 x 6 3x2x2	1 x 6 3x2x2	1 x 6 3x2x2	1 x 6 4x2x2	1 x 6 4x3x2	1 x 6 4x2x2
1	59	63	59	59.5	59	59.3	59.4	61	65	60	60.4	63.5	60	60.2	60.2	60.2	60.2	60.2	60.2	60.2	60.2	63.4
2	263	-	258 1321	295 1324	257 1316	258 1321	259 1322	210	-	207	208	227	206	207	207	207	207	207	207	207	207	226
3	890 1200	-	873	881	873	879	879	812 1500	-	790 1525	796 1532	907 1621	788 1518	793 1527	793 1527	793 1527	793 1527	793 1527	793 1527	793 1527	793 1527	905 1618
4	1980	-	2002	2028	2005	2023	2026	2000	-	2047 3319	2063 3339	2367 4366	2046 3299	2058 3323	2058 3323	2058 3323	2058 3323	2058 3323	2058 3323	2058 3323	2058 3323	2364 4359
5	3318	-	3465	3503	3454	3488	3498	3218	-	-	-	-	-	5411	5411	5411	5411	5411	5411	5411	5411	6415
1	800 794	790	691	797	778	742	784	900	910	692	913	940	879	857	857	857	857	857	857	857	857	915
2	2500	2450	2113	2430	2373	2263	2392	2800	2650	2126	2775	2878	2673	2604	2604	2604	2604	2604	2604	2604	2604	2810
3	4200	4100	3683	4197	4105	3912	4142	4700	4600	5597*	4761	5025	4595	4463	4463	4463	4463	4463	4463	4463	4463	4927
4	5900	5950	5559	6247	6125	5840	5998	6000	6600	3710*	6989	7701	6761	6555	6555	6555	6555	6555	6555	6555	6555	7582

Table 3.4 (Continued)

* Based on mode recognition

3.4. Curved Fan Blade

First shell problem which was solved to test the program was the fan blade, originally studied by Olson and Lindberg. They performed the experiments on a model which was constructed by rolling a piece of sheet steel 3.048mm (0.12 inches) thick, to a radius of curvature of 609.6mm (24.0 inches). This curved sheet was then cut to size 304.8 x 304.8 mm (12x12 inches) and welded to a steel block to simulate the clamped boundary condition.

Initially Olson and Lindberg (43,44) used a cylindrical shell element, with four nodes and 28 degrees of freedoms, to predict the natural frequencies of this fan blade. Later Lindberg et al (39,45) solved the same problem by using a curved triangular shallow shell element which was more accurate than the previous one. They predicted the first 25 frequencies within a few per cent of error.

The problem was also solved by Bridle (11) using power series method, by Neale (42) using a hybrid shell element, and by Martins and Owen (40) using the semi-loof element. The original form of the thick shell element was used by Ahmad et al. (2) to solve the problem. They noticed that the element was too stiff. After the modifications Hofmeister and Evensen (30), and Bossak and Zienkiewicz (9) applied the element to the same problem obtaining very good results.

Table 3.5 lists the first five natural frequencies reported in the references mentioned above, together with the results of a 4 x 4 mesh of this study which is in excellent agreement with the experimental values. Table 3.6 contains the first seventeen frequencies reported in (43) and (39), and the results of the

Ref. Grid	39 Exper.	43 4 x 4	39 4 x 4	42 6 x 6	11 Power Series	40	2 4 x 4	30 2 x 3	9 3 x 3	Present 4 x 4
1	85.6 86.6 (43)	93.5	86.6	89	85.8	85.3	113.0	87.0	88.3	86.2
2	134.5	147.6	139.2	143.2	138.3	138.1	147.0	143.0	142.8	139.5
3	259	255.1	251.3	252.6	246.7	245.1	296.0	252.0	257.6	249.8
4	351	393.1	348.6	380.5	342.5	340.4	440.0	367.0	369.2	347.9
5	395	423.5	393.4	428.2	386.8	383.8	475.0	412.0	441.8	405.0
D.O.F.	-	175	300	390	108	275	325	145	288	280

Table 3.5. Predictions of the first five natural frequencies by different references for the curved fan blade.

Freq. No.	Exp. * (39)	Ref. (39)	Present 4 x 4	Ref. (43)
1	85.6 86.6(43)	86.6	86.2	93.5
2	135	139	139.5	148
3	259	251	249.8	255
4	351	348.6	347.9	393
5	395	393	405	424
6	531	533	549	534
7	743	752	771	782
8	751	746	756	792
9	790	790	817	863
10	809	813	911	862
11	997	1008	1100	1002
12	1216	1231	1383	1175
13	1252	1246	1266	
14	1241	1266	1371	
15	1281	1286	1583	
16	1310	1303	-	
17	1706	1652	1762	

Table 3.6. Comparison of the first 17 Natural Frequencies for curved fan blade.

Mode	Exp. (39)	2 x 2	3 x 2	2 x 3	3 x 3	4 x 4
1	85.6 86.6(43)	94.1	91.66	86.6	87.11	86.2
2	135	145	145	143.2	141.8	139.5
3	259	250.6	258	252.5	255.6	249.8
4	351	402.4	385	366	359	347.9
5	395	418	448	411	430	405
6	531	828	595	717	584	549
7	751	807	801	736	755	771
D.O.F.*		80	110	120	165	280

Table 3.7 Frequencies of the fan blade using different meshes.

* Total number of degrees of freedom after the boundary conditions.

Mode	2 x 3 (30) ₊	2 x 3* 3x2x2**	2 x 3 4x2x2	3 x 3 3x2x2	3 x 3 4x2x2	3 x 3 4x3x2
1	91	90	91	88	88	91
2	149	144	148	143	144	146
3	310	270	298	261	264	282
4	383	374	380	362	364	407
5	556	455	535	461	479	582

Table 3.8. Frequencies of the fan blade using 10-node elements

+ Results for 12-node element ref (30)

* Mesh used

** Order of integration

modified eight-node element with 4×4 mesh. The values on table 3.7 demonstrate the importance of the location of the nodal points and the choice of the mesh used. Finally, on table 3.8 the frequencies obtained by using ten-node elements are compared with the results of the twelve-node element of reference (30).

3.5. Circular Cylindrical Shell

This example was chosen to test the program on a cantilevered cylindrical shell. The configuration is shown on figure 3.1. The experimental values of the frequencies for this cylinder were reported by Gill (27). He calculated the frequencies using the method given in (51). Wilson (57) also solved the problem using five axisymmetric shell elements. The experimental and the calculated values of the frequencies of the references (27,57), and the results of the present study are listed on table 3.9.

In the present study, the frequencies were calculated by using modified eight-node elements with different meshes. Due to the symmetry only half of the cylinder was considered. Even with 2×2 mesh reasonable results were obtained for the simple modes like $1/1$, $2/1$, $2/2$. A finer mesh was required for the accurate determination of higher frequencies.

3.6. Curved Cantilever Beam-Pipe Segments

The examples solved in sections 3.4 and 3.5, showed that the program is capable of dealing with shallow and closed shell structures. The hollow blading of chapter 5 consists of two

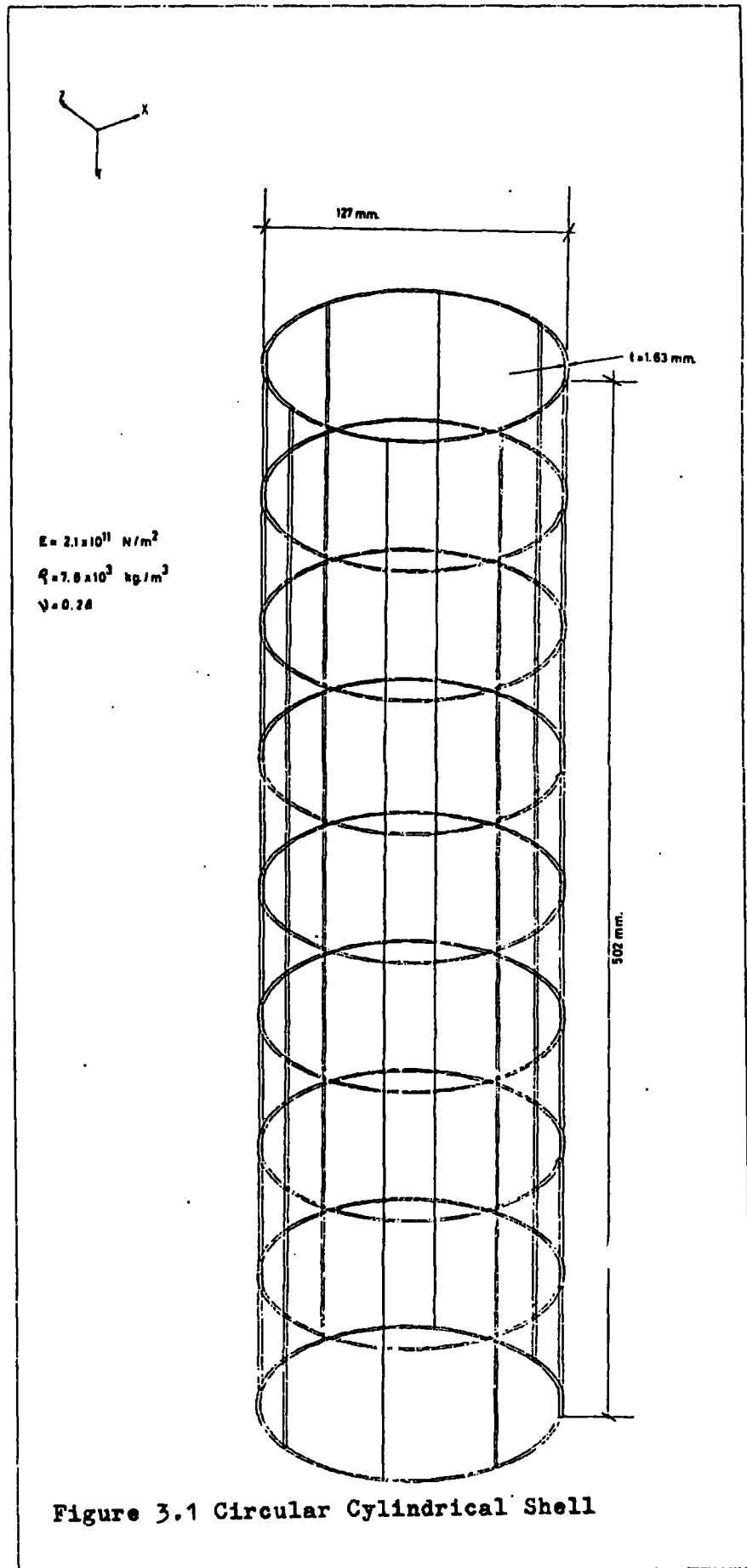


Figure 3.1 Circular Cylindrical Shell

Mode n/m*	Exp.	Ref. (27)	Ref. (57)	**				
				2 x 2	3 x 5	4 x 8	4 x 10	6 x 8
1/1	364	-	470	466	470	468	468	468
2/1	293	319	315	289	334	320	320	316
3/1	740	767	767	1274	916	821	816	779
4/1	1451	1462	1461	-	-	-	1782	1527
5/1	2236	2361	2359	-	-	-	-	-
1/2	-	-	2061	-	-	-	2025	2055
2/2	827	1017	943	990	960	996	951	941
3/2	886	928	914	-	-	1061	986	932
4/2	1503	1521	1517	-	-	-	1832	1597
2/3	1894	2393	2212	-	-	-	2230	2213
3/3	1371	1511	1459	1632	-	-	1593	1505
4/3	1673	1726	1712	-	-	-	1992	1874
4/4	2045	2158	2122	-	-	-	2338	2494

Table 3.9 Natural frequencies of circular cylindrical shell.

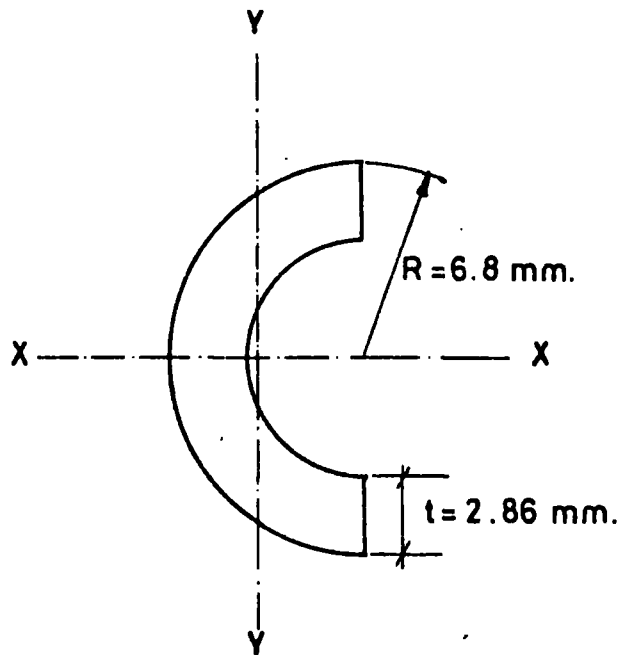
* n/m where n is the number of nodal diameters
and m = number of nodal circles + 1

** Mesh used

shallow shells, similar to the fan blade of section 3.4., connected by two circular pipe segments. The analysis in this section was performed in order to decide on a reasonably accurate representation of these pipe segments.

The cantilever beam with the cross-section shown on figure 3.2 was analysed. Being a slender beam, all of its lower modes of vibration were of bending type. First six natural frequencies - first three bending in \bar{X} - X plane and first three bending in Y - Y plane - were calculated by using the closed form formula for the natural frequencies of a beam. These modes of vibration coincided with the first six modes of the finite element analysis which was performed by using different meshes and integration points. The natural frequencies were calculated by using both the eight-node and the ten-node elements. The meshes used were either 1×8 or 2×8 . Eight elements along the length were used in order to have compatible results for the analysis of chapter 5 where two beams similar to the present one were used as parts of the oval cross-section blading. The results are given on table 3.10.

The eight-node element with 1×8 mesh and $2 \times 2 \times 2$ integration points was too flexible and did not give any reasonable results. Increasing the number of integration points from two to three in the direction of curvature improved the results and brought them within acceptable limits. An additional integration point along the length caused small increases in the higher frequencies. A refined mesh with 2×8 elements gave good results with $2 \times 2 \times 2$ and $3 \times 2 \times 2$ integration points,



$$I_{xx} = 745.012 \text{ mm}^4$$

$$I_{yy} = 154.144 \text{ mm}^4$$

$$\text{Area} = 48.249 \text{ mm}^2$$

$$\text{Length} = 380.0 \text{ mm}$$

$$E = 2.11 \times 10^{11} \text{ N/m}^2$$

$$\text{Density} = 7.85 \times 10^3 \text{ kg/m}^3$$

Figure 3.2 Assumed cross section of the beam segment.

Mode No.	Beam Theory	8 - node Elements				10 - node Elements.				
		1 x 8 2x2x2	1 x 8 3x3x2	2 x 8 2x2x2	2 x 8 3x2x2	1 x 8 3x2x2	1 x 8 4x2x2	1 x 8 4x3x2	2 x 8 3x2x2	2 x 8 4x2x2
1	79	64	71	71	79	80	82	82	79	79
2	495	394	440	445	483	491	508	502	485	485
3	1385	1079	1215	1246	1333	1348	1416	1382	1336	1337
1	36	11	34	34	36	34	35	35	36	36
2	225	70	214	214	226	218	224	223	227	227
3	630	201	615	618	639	627	646	644	641	642
x-section Area mm ²	483	411	411	411	476	488	488	488	483	483

Table 3.10. Results for the cylindrical pipe segment.

later being more accurate. The results of ten-node element was much better for both 1×8 and 2×8 meshes. Order of integration did not change the results much, but the use of three integration points along the length caused slight over-stiffness of the element. For this particular problem, the order of integration did not effect the calculated volume of the elements, but in general the volume was calculated more accurately by refined meshes or when ten-node elements were used. More accurate results were obtained when the volume was determined accurately.

The results indicate that the ten-node element is more suitable for the finite element modelling of this beam. A 1×8 mesh with $4 \times 2 \times 2$ integration points seems to be the best choice when the economy is taken into consideration.

3.7 Conclusion

A modified eight-node element with the stiffness matrix integrated at $2 \times 2 \times 2$ Gauss points gave excellent results for all plate and shallow shell analysis. Only in two cases, those of pretwisted blading and of pipe segments, was this element found to be too flexible. Increasing the number of integration points from two to three in the direction of curvature, increased the accuracy. In these two cases, the performance of the ten-node element was very good.

For the numerical integration of the mass matrix $3 \times 3 \times 2$ Gauss points were used for both eight-node and ten-node elements. Also $3 \times 3 \times 3$ integration points were tried for integrating the mass matrix of the eight-node element and

4 x 3 x 2 integration points for the ten-node element. Their effects on the results were unnoticeable.

For all the examples the mode shapes determined were in a very good agreement with the mode shapes given in the original references.

CHAPTER 4

4. EXPERIMENTAL METHOD

The majority of the experiments were undertaken to confirm the numerical predictions of the natural frequencies and the mode shapes of an oval cross-section hollow blade. In addition to these, some subsidiary experiments were performed to see the effects of various assumptions on the experimental and numerical results.

Geometry and construction of the shells, and the results of the experiments are given in the related sections of chapters 5 and 6. In this chapter, measurement of the geometry of the shell, apparatus and the procedure will be given.

4. 1 Determination of the Actual Geometry of the Shell

After the blade was manufactured (see 5.2), some measurements were taken in order to determine its actual geometry. A grid, which was later used in the finite element analysis, was drawn on the surface of the shell. Relative distances of each node of the grid from a base line was measured by means of a micrometer. The measurements were taken both along $y = \text{const.}$ and $x = \text{const.}$ lines. Then, taking one of the nodes as the reference point z coordinates of all the nodes were calculated by using the two sets of data, thus cross-checking the results. At the end, the average values were taken, and these were checked against the direct measurement of the distance between the faces of the blade.

In spite of the careful measurements there were some differences between coordinates that were determined by different methods. So, some adjustment had to be made by close visual

inspection of the shell. Also, there was no way of measuring the coordinates of the points at the bottom faces of the elements, since they were trapped inside the shell. Coordinates of these points were calculated analytically.

4.2 Apparatus and Procedure

All the experiments were conducted on a vibration table which consisted of a cast-iron bedplate, mounted on a concrete block. The flanged channels on the surface of the bedplate was used to bolt the base of the shells on to the table. To prevent the shell from rocking, it was found to be useful to use as many bolts as possible. In some instances, it was observed that the extra bolts were effecting the bending type of frequencies, increasing them by up to 4 per cent.

The shells were excited over a large range of frequencies, using an oscillator which was driving a coil through an amplifier. In spite of the amplifier, the magnetic field created was not sufficient above 500 Hz., and it was necessary to have the coil very close to the shell surface. The position of the coil was often changed for a better excitation of different modes. The improvement gained was small, although it marginally helped in the torsional type of modes.

The response of the shell was indicated by piezoelectric strain gauges cemented on the surface of the shell. Up to six gauges were bonded to the shell, different ones being useful for different modes.

Output from the gauges was displayed on the lower trace of an oscilloscope where it was compared with the forcing frequency

displayed on the upper trace. Usually the shell responded at some whole number multiple of the forcing frequency. When the core of the coil was not a permanent magnet, the response of the shell was at twice the driving frequency. The output signal of the gauges reached their maximum when a natural frequency was excited. This frequency was accurately determined by using a Muirhead-Wigan Decade oscillator.

Once a frequency was detected, the identification of the mode shape was done using a hand-held piezoelectric gauged probe. The gauge on the probe converted the displacements of any vibrating surface into an electrical signal which was displayed on the upper trace of the oscilloscope instead of the forcing frequency. By touching a particular point on the shell with the probe, the signals were compared as either in or out of phase with the signals from the strain gauge. A thorough examination of the shell surface yielded the mode shape for the particular frequency.

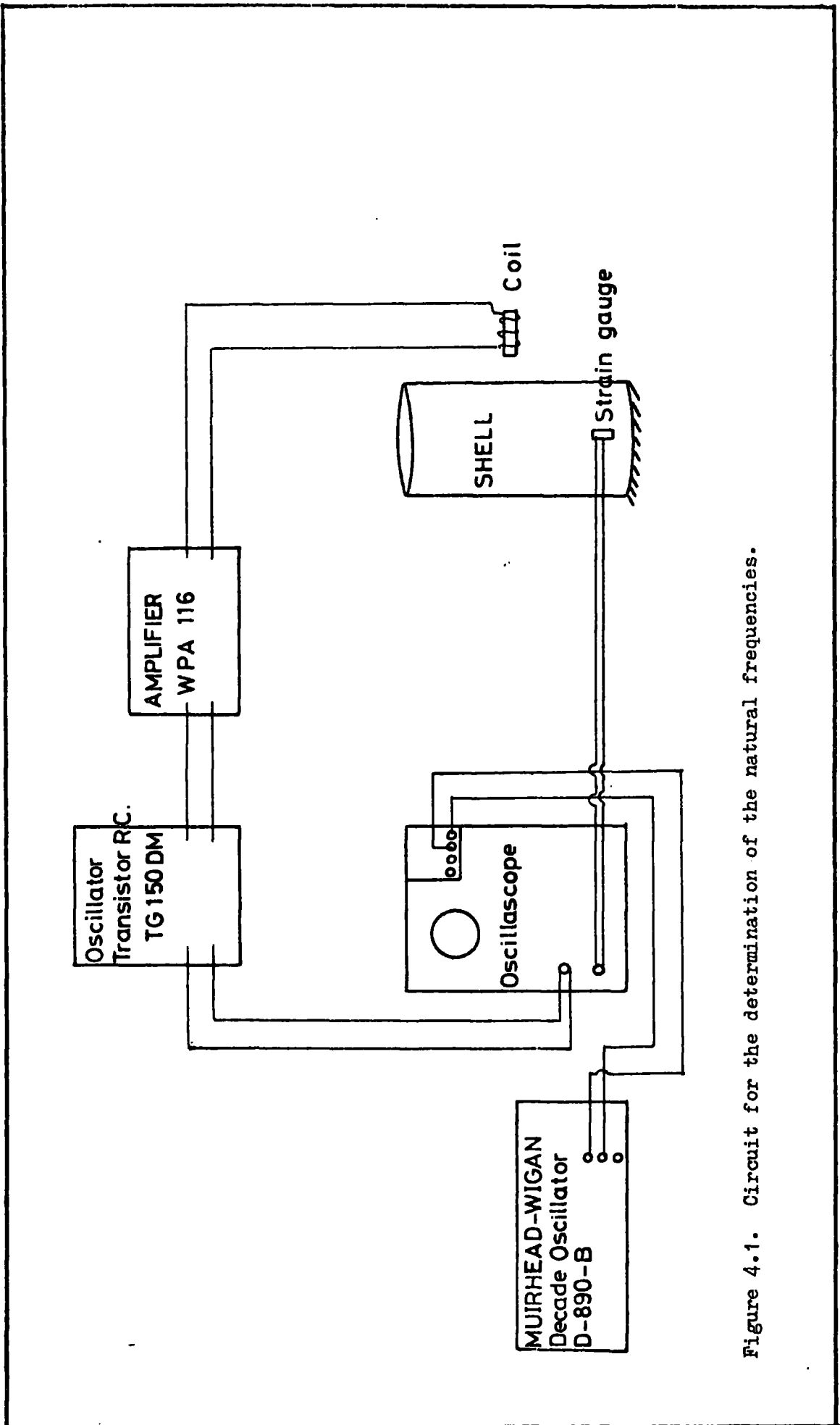


Figure 4.1. Circuit for the determination of the natural frequencies.

CHAPTER 5

5. OVAL CROSS-SECTION HOLLOW BLADE

In this chapter, the computer program which was explained in chapter 2, and checked against several test examples in chapter 3, was used to analyse hollow blading. The natural frequencies predicted by the program were compared with the results of the experiments performed on a model blade, and the causes of some of the differences were studied.

The geometrical descriptions of the experimental and the mathematical models of the blade are given in the following section.

5.1 Geometry of the Blade

The blade was designed to have a constant cross-section of oval shape, consisting of four circular segments with two different radii of curvature along the whole length. (Fig.5.1). To have a continuous surface, the circular segments were so located that they would be tangent at the intersecting points. This condition required the following relationship to be satisfied:

$$R = \frac{a^2 + b^2 - 2ar}{2(b - r)} \quad (5.1)$$

The parameters of this relationship are as shown in Fig. 5.1. and they were chosen as:

$$\begin{array}{ll} a = 105 \text{ mm} & r = 6.76 \text{ mm} \\ b = 20 \text{ mm} & R = 377.6 \text{ mm} \end{array}$$

The length of the blade was taken to be 420 mm.

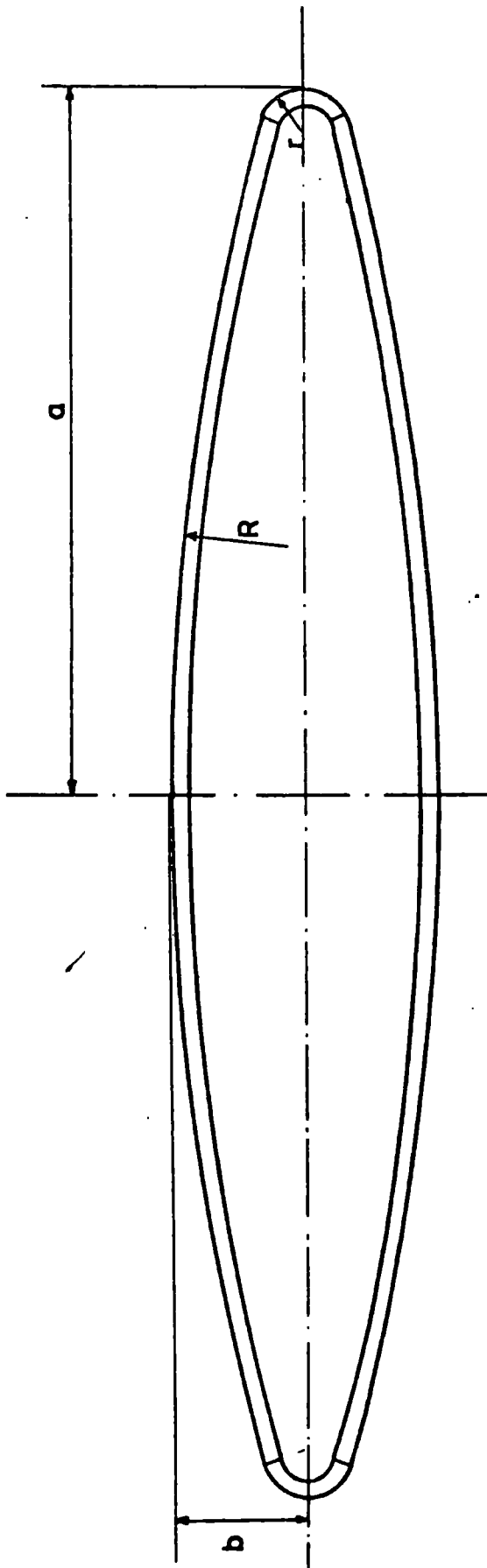


Figure 5.1. The cross-section of the design geometry of oval blade.

Due to the difficulties encountered during the manufacturing process (see Sec. 5.2.), the experimental model of the blade had some geometrical imperfections. To achieve a smooth surface then, was only possible by filing it along the intersecting lines, and the true mathematical description of the experimental model became impossible.

5.2 Manufacturing the Experimental Model

A model of the blade described in the previous section was manufactured for laboratory testing. Two larger cylindrical segments were obtained by rolling two 200 x 420 mm steel sheets of 2.54 mm thickness, to a radius of curvature of 380 mm. The two small segments were cut out of a steel pipe of 13.5 mm outer diameter. The wall thickness of the pipe was 2.75 mm.

Two successive different methods were employed to connect the four pieces together and to fix the root of the assembly. The experiments were performed for both cases.

5.2.1 Solder Assembly

In the first case brazing was used to connect the pieces. The idea was to avoid welding, thus to prevent the pieces from warping under high temperatures. But the pipe segments were inevitably bent during the cutting process due to stress relaxation, so the final assembly was a doubly curved shell with an unwanted curvature along the length, (see fig. 5.2). Since this curvature was non-uniform the symmetry of the blade was lost. Measurements taken on the blade showed that the geometric parameter b was varying from 16.5 mm at the tip to 23 mm at the middle cross-section.

The root fixing was achieved by embedding 40 mm of the blade into a solder base. The clear length of the blade for this case was 380 mm.

The analysis in Section 5.3 has shown that the brazing was not stiff enough to assume the blade as a single piece of metal. Also the solder which was used to fix the root was too flexible to simulate the clamped boundary condition. These difficulties were overcome in the second method at the expense of using welding.

5.2.2 Weld Assembly

In the second case the aim was to have the inter-connections of the pieces stiff enough to treat the whole blade as a single piece of metal, also to have the root fixing as stiff as possible to be able to assume that it was clamped. For this purpose the brazing along the inter-connecting lines was replaced by welding. Also the solder base was removed and the blade was welded onto a steel plate. As expected, the high temperature caused some irrecoverable local distortions in the geometry. The measurements that were taken to determine the final shape of the blade indicated that, in addition to the unwanted curvature now, there was also a slight pretwist on the blade.

To use for numerical analysis the coordinates of each node were calculated. The finite element model of the blade shown in figure 5.2 was drawn using these coordinates. This geometry of the blade will be referred as the real geometry of the blade in the following sections.

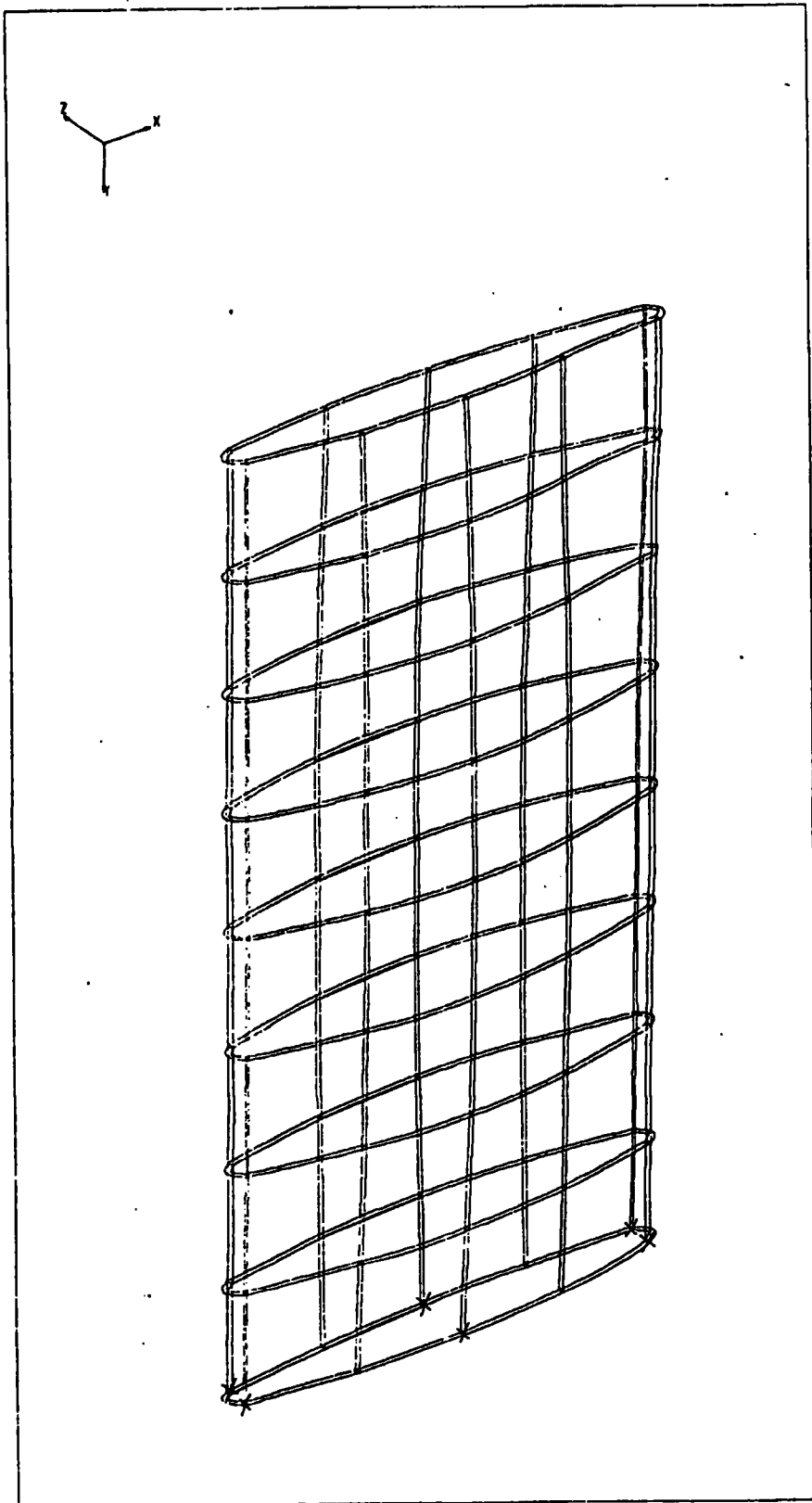


Figure 5.2 Real Geometry finite element model of the blade.

X See footnote on page 56.

The average values of these coordinates yield to the following approximate parameters for the geometry:

$$\begin{array}{ll} a = 105 \text{ mm} & r = 5.5 \text{ mm} \\ b = 19 \text{ mm} & R = 410 \text{ mm} \end{array}$$

The geometry of the blade described by these parameters was used to determine the ideal mode shapes in Section 5.5.1. and will be referred as ideal geometry.

5.3 Analysis of the Blade Assembled using Solder

In this section, the experimental model of the blade described in section 5.2.1. is analysed. First three natural frequencies obtained from the analyses are listed on table 5.1. Column 1 contains the experimental frequencies, columns 2 and 3 give the results of the finite element analyses. As the results indicate the experimental model was too flexible to give any comparable results. The frequencies listed in column 3 were obtained by representing the brazing between the pieces as hinges* in the finite element model. The effect of this assumption on the ovalling and torsional modes was to reduce them considerably, although it did not change the bending mode at all.

The results showed that a rigid connection between the pieces was not achieved in the experimental model, and overflexibility of these connections were resulting low torsional and ovalling frequencies. Bending mode of vibration which was independent of the type of connections between the pieces, depends greatly on the type of root fixing which is discussed in the next section.

* See addendum on page 52.

Mode	Exp.	8 x 10 F.E.*	8 x 8 F.E.**
Bending	151	232	233
Ovalling	197	345	198
Torsion	256	317	135

Table 5.1: First three frequencies obtained experimentally and by finite element analysis.

- **rigid connections**
- ** **hinged connections**

Frequencies given are in Hertz.

APPENDIX

The representation of the hinged connections was achieved in a simple way during the elimination process. The rotational degrees of freedom, at the nodes along the line connecting the pipe segment to the shallow shell, were eliminated as soon as the contribution of the elements along the pipe segment were completed but before any contribution from the elements of the shallow shell was included. The process was reversed when crossing the connecting line from shell to pipe segment. The elimination of u, v and w degrees of freedom was done following the usual procedure.

5.4 Effect of the Root Fixing on the Natural Frequencies

To achieve the idealized clamped boundary conditions in the experiments is very difficult. In this part of the study, the behaviour of different experimental simulations of the clamped boundary conditions was examined.

A set of experiments were performed on two identical cylindrical steel pipes. The pipes had an outer diameter of 50.8 mm and a wall thickness of 1.59 mm.

The first piece of pipe was 260 mm long and it was welded on a steel block of dimensions 305 x 82.5 x 25.4 mm. The pipe responded to the first bending mode (which also was the fundamental mode of vibration) at two different resonant frequencies. When it was excited to vibrate parallel to the length of the base plate the resonant bending frequency was found to be 602 Hz, whereas the same pipe responded to the same mode of vibration at 564 Hz when it was excited to vibrate normal to the previous direction. The second mode of vibration was the first ovaling mode and it was found to be 1783 Hz, independent of the direction of excitation. The third natural frequency was the second ovaling mode at 2300 Hz.

The root fixing of the second piece of pipe was achieved by embedding it into a block of solder having dimensions 84 x 75 x 26 mm. The average length of the pipe was 251 mm. The resonant frequencies corresponding to the first bending mode were determined to be 395 Hz and 417 Hz in two normal directions. The ovaling mode responded to a resonant frequency of 1796 Hz.

Table 5.2 lists these experimental results, together with the results obtained from the computer program. For the numerical analysis a 6 x 8 grid was used, and due to the symmetry, only half of the cylinders were considered. Clamped boundary conditions were assumed for both cases.

Case	Exp.	Num.	Exp.	Num.
length (mm) Mode	260	260	251	251
1	564/602	706	395/417	755
2	1783	1773	1796	1776
3	2300	2297	-	2358

Table 5.2. Comparison of the experimental and numerical frequencies of two similar cylinders. Different methods for root fixing were used for the experiments.

As expected the results of the finite element analysis gave a higher bending frequency for the shorter cylinder. But in the experimental case, due to the effect of the different methods of root fixing, the situation was reversed. The solder base is definitely not suitable to represent the clamped boundary conditions for the experiments.

5.5 Weld Assembled Blade

5.5.1 Finite Element Analysis using Ideal Geometry

The numerical analysis was first performed on a regular geometric shape, approximated by the average geometric

parameters given in section 5.2.2., to determine the ideal mode shapes for the blade.

A 10 x 8 grid (same as on figure 5.2) was used in the analysis, and the length of the blade was taken to be 420 mm. To represent the pipe segments at the sides, the cubic-quadratic element of section 2.2. was used. The number of integration points for these elements were 3 along the curvature, 2 along the length and through the thickness.

The sketches of first 20 mode shapes, and deflection curves for some of these modes are given in appendix 3. Since these mode shapes were obtained by using a regular geometry, and they show a regularity themselves, they are referred as ideal mode shapes. In general, it is possible to classify them in different groups except for the last, mode number 20. Inconsistent shape of this mode could be due to the coupling of bending mode with two nodal diameters, and the mode type 8 with three nodal diameters.

Table 5.3 lists the frequencies obtained by using different number of master degrees of freedoms with 10 x 8 grid. In addition the frequencies for some of the symmetric modes which were calculated by using only half of the blade (making use of the symmetry around z -axis) are included.

The frequencies listed in column 1, were calculated by keeping both u and w displacements as master degrees of freedoms at 92 nodes. Two different frequencies corresponding to mode 6 are the result of the coupling with the bending mode in x - direction. Detailed tip deflections for these two frequencies are given in figure A.3.1. in appendix 3.

Frequencies given in column 2 were obtained by keeping only w displacements at 102 nodes. Since the master degrees of freedoms were distributed more evenly over the blade than they were in the previous case, they gave better results especially for the mode shapes with three nodal diameters. A similar comparison can be made between the second and third columns where only 88 w displacements were kept as masters.

Columns 5 and 6 contain the frequencies which were calculated by using only half of the blade. The retained master degrees of freedoms which gave the frequencies in column 5 were equivalent to the ones of column 3.

The frequencies given in columns 4 and 7 were calculated by restraining only six* and four* nodes respectively at the root of the blade, to simulate more flexible root fixing condition. The greatest effect of this assumption was reflected on the bending frequencies which were reduced by about 30 per cent.

5.5.2. Finite Element Analysis - Real Geometry

The analysis in this section was performed on the geometry that was described by the coordinates of nodes as measured directly on the blade (Section 5.2.2). Thus, a closer approximation to the geometry of the experimental model was achieved.

The finite element idealization used was very similar to the ideal geometry case of section 5.5.1, with the grid shown in figure 5.2. The y -coordinates of the tip nodes were changing between 413 mm and 417 mm, giving an average total length of 415 mm to the blade. In all the analyses in this section the

* See figure 5.2. on page 50 for the location of these nodes.

Column	1	2	3	4*	5	6	7*
Grid	10 x 8	10 x 8	10 x 8	10 x 8	5 x 8	5 x 8	5 x 8
DOF Mode	184	102	88	104	48	88	48
1	196.5	195.6	197	142	197	197	139
2	302.5	302.5	304	299	304	303	299
3	335.5	335.7	337	327	-	-	-
4	562	562	590.4	564	-	-	-
5	553	553	589.8	551	589.8	581	567
6	809/879	821.5	837	830	-	-	-
7	804.7	804.5	873	809	873	866	809
8	953.6	955	962	961	961.5	961	961
9	950	951	1040	1029	-	-	-
10	1062	1064	1148	1144	1148	1143	1144
11	1047	1047	1372	1338	-	-	-
12	1042.7	1042	1378	-	1378	1356	1363
13	1191	1191	1485	-	-	-	-
14	1359	1359	1754	-	-	1696	1696
15	1418.7	1436	1434	-	-	-	-
16	1510	1502	-	-	-	-	-
17	1589	1583	-	-	-	-	-
18	1604	1598	-	-	-	-	-
19	-	1733	1755	-	-	-	-
20	1514	1513	-	-	-	-	-

Table 5.3. Predicted natural frequencies for the ideal geometry oval blade.

* Partial restraining of the root.

master degrees of freedoms were chosen among the w displacements only.

First seven natural frequencies are listed on table 5.4 with the sketches of the tip deflections of the modes, and the results of experimental and ideal geometry analyses. Higher frequencies, corresponding mode shapes, and some of the deflection curves are included in appendix 3.

5.5.3. Experimental Analysis

Some of the natural frequencies and the corresponding mode shapes of the experimental model of the blade were determined following the procedure given in chapter 4.

The location of the nodal lines were very much effected by the imperfections of the geometry of the blade, especially for high natural frequencies. Also, for these natural frequencies the magnetic field created by the coil was relatively weak. These, together with the suppressive effect of the hand-probe when in contact with the shell, made it extremely difficult to identify the mode shapes of high frequencies.

First seven frequencies are listed on table 5.4. Sketches of the mode shapes of these and some higher frequencies are given in appendix 3.

5.5.4. Comparison of the Results

Table 5.4 lists the first seven natural frequencies obtained from experimental and finite element analyses. These frequencies and the corresponding mode shapes are in a good agreement in all three analyses.

The biggest difference between the frequencies is seen in the bending mode. This is an expected situation since the effect of the simulation of the clamped boundary condition is greatest for this mode (this was pointed out both experimentally and numerically in sections 5.4 and 5.5.1.). Comparison of the frequencies listed for real geometry and ideal geometry analyses also indicate that this mode is very sensitive to the geometry used in idealization. Although the nodal coordinates of the real geometry case were taken directly from the experimental model of the blade, the accuracy of these measurements is subject to discussion. Considerations of these facts make the difference in the frequencies of the bending mode tolerable.

Fifth mode, which is the second ovaling mode, of real geometry and experimental analyses seems to couple with the second torsion mode (figure A.3.20, and A.3.32). Also slight coupling of sixth mode with the second bending can be seen on figure A.3.20.(a). These are the first signs of the affect of imperfections of the geometry on the mode shapes. Similar couplings were observed in section 6.3.1, where the pretwist of the blading was considered. Although the pretwist of the geometry considered in this section was very small, the frequencies corresponding to these two sets of modes were fairly close (within about 10%) and the overall irregularity of the geometry may account for these couplings.

During the experiments a frequency at 708.5 Hz was detected, but the identification of the mode shape was not

possible due to the difficulties explained in section 5.5.3. This frequency is inserted into the table as the experimental second bending mode which is in a good agreement with the finite element analyses.

Effects of the imperfections of the geometry on the mode shapes become visible at fifth mode, and increases as the mode shapes get more and more complicated. Coupling occurs between two or even three close modes, and the identification of the resulting mode shape becomes very difficult, sometimes impossible. In appendix 3, the shapes of some of the modes above 7 are examined, and possible matchings between the results of different analyses are pointed. Some of the mode shapes that were obtained either by the finite element analysis of the real geometry, or by the experimental analysis are not included at all, because of the extreme difficulties of identifying them.

5.6 Effects of the Material Properties and the Choice of Degrees of Freedoms on the Natural Frequencies and the Mode Shapes.

In all the numerical calculations performed in this chapter the following nominal material data was used.

$$\begin{aligned} \text{Young's Modulus} &= 0.211 \times 10^{12} \text{ N/m}^2 \\ \text{Density} &= 0.785 \times 10^4 \text{ kg/m}^3 \\ \text{Poissons Ratio} &= 0.285 \end{aligned}$$

In order to observe the variation of the natural frequencies with the changes in material properties, the numerical calculations were repeated for the real geometry of the blade using 10×4 grid for different material data.


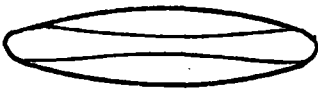

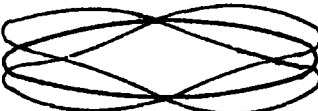



Mod. No.	Cross-section	Experiment.	R.G.	I.G.
1		142	158	196.5
2		296	306.8	302.5
3		355	371.5	335.5
4		561	588.4	562
5		595	648	553
6		812	836	821.5
7		(708)*	739	804.5

Table 5.4. Comparison of the results of the experimental and finite element analyses. The frequencies given are in

* Mode shape was not identified experimentally.

The results obtained are listed on table 5.5. An increase of 2.03% in Young's Modulus caused 1.014% increase; and 0.637% decrease of the density caused 0.32% increase in all frequencies. These changes are exactly as expected, since the frequencies are directly proportional to $(E/\rho)^{1/2}$. Changing the poissons ratio, in the possible range of 0.25 to 0.33 effected different frequencies in different ways, but in all cases the change in the frequencies was less than 2%.

A similar observation was made by changing the master degrees of freedoms retained. The results are given on table 5.6. Again 4 x 10 grid was used, and there were total 640 degrees of freedom after the insertion of boundary conditions. In general keeping more w displacements distributed over the blade gives better results than any other combination of master degrees of freedoms. An exception to this reasoning was seen only for mode 6, and it was the effect of the bending mode in x-direction.

5.7. Discussions of Results

Natural frequency analysis of an oval cross-section hollow blade has been presented. The geometry chosen is interesting from the point of view that it contains both thin-shallow and thick-deep shell properties in it. The interesting part of the numerical analysis on the other hand, is that the same type of finite element was used to represent both types of shell properties.

Difficulties that arose in manufacturing the experimental model, imposed some limitations on accurate representation of the mathematical model. Two different mathematical representations

E	0.211×10^{12}	0.2153×10^{12}	0.211×10^{12}	0.211×10^{12}	0.211×10^{12}	0.211×10^{12}
ρ	0.785×10^4	0.785×10^4	0.780×10^4	0.785×10^4	0.785×10^4	0.785×10^4
ν	0.285	0.285	0.285	0.250	0.300	0.330
1	162.23	163.874	162.749	162.745	162.029	161.664
2	320.111	323.356	321.136	318.677	320.837	322.497
3	375.274	379.078	376.474	374.120	375.883	377.317
4	621.429	627.730	623.418	621.060	621.733	622.614
5	660.429	666.754	662.174	659.237	660.541	661.738
7	782.532	790.466	785.036	782.689	782.534	782.674
6	870.081	878.902	872.866	864.331	872.888	879.153
8	1011.85	1022.11	1015.09	1006.29	1014.62	1020.85

Table 5.5. Variation of natural frequencies with the material properties.

D.O.F.	118 w	60w	60w+60u	60w+60v	60w+60 α	60w+60 β
MODE						
1	162.23	162.24	162.24	162.24	162.238	162.243
2	320.111	320.261	320.227	320.257	320.160	320.253
3	375.274	375.953	375.580	375.945	375.355	375.939
4	621.428	624.996	624.275	624.830	623.348	624.601
5	660.062	667.445	665.841	666.719	666.546	666.330
6	870.081	876.025	827.181 937.434	856.149	871.007	875.804
7	782.532	790.022	786.725	789.004	780.832	789.476
8	1011.85	1031.49	1028.26	1021.58	1018.83	1030.10

Table 5.6. Variation of the Natural Frequencies with different choice of master degrees of freedoms.

were used for the finite element analysis. First, using a regular but approximate geometric shape, general trend of the idealized mode shapes were determined. Seven of these predicted modes and the natural frequencies were confirmed by the experiments. Later, the accuracy of the predictions were further increased by describing the geometry of the mathematical model closer to the experimental one.

Some of the causes of the differences between the results of various analyses were studied experimentally and numerically. The studies indicated that the clamped boundary conditions were not simulated properly on the experimental model. In spite of that, the results obtained were quite satisfactory.

CHAPTER 6

6. FUTURE APPLICATIONS

The range of application of the computer programs which are developed to serve a particular purpose is usually limited. The limitations that isoparametric shell element has, are the difficulties of modelling the sharp corners and multiple junctions in the structures. For a structure where these type of connections are dominant, it might be better to employ some other element, for instance semiloof, for which they do not represent any difficulty. On the other hand this new element may not offer the same facilities as the original one, for instance it may not be able to represent the properties of a thick shell.

In this chapter, some simple modifications are introduced to extend the range of applicability of the isoparametric shell element. Some assumptions are made for very sharp corner connections and multiple junctions. Also the stiffeners are considered for hollow shells. However, if the nature of the problem in hand changes considerably, it may be more practical to employ a general purpose program for the solution.

6.1. Sharp Corner Connections

Instead of an oval cross-section, when an aerofoil cross section is assumed for the hollow blading, as shown in figure 6.1, immediately a difficulty arises in representing the sharp corner. To satisfy the continuity of the displacement between the elements joining at this corner, the nodes which are common to these elements must be defined uniquely. A unique

definition of the nodes at the sharp corner is achieved by assuming the nodal connections between the elements as shown in figure 6.2. Application of this assumption to the section shown in figure 6.1 introduces another difficulty. When one defines the top and bottom points of the thickness vector at the corner, and goes around the aerofoil cross-section by defining the thickness vectors at every node, it is apparent, when reaching the corner node again, that the top and bottom points of the node have to be reversed. This does not affect the continuity of the u , v and w displacements since they are defined at the midpoint of the node, but may affect the rotational degrees of freedoms, α and β since their definition depends upon the thickness vector. Defining two different thickness vectors for the same node is discussed and illustrated on a cantilever plate in the next section.

6.1.1. Cantilever Plate

The rectangular cantilever plate shown in figure 6.3 is represented using three elements. The direction of the curvilinear axes of each element are so chosen that the thickness vectors corresponding to nodes 6,7 and 8 are pointed in opposite directions for the elements 1 and 2. Since the displacements u , v and w are the displacements of the mid-point of the thickness vector in the global system, they do not represent any difficulty in joining the element matrices. On the other hand, α and β are defined as the rotation of the thickness vector \underline{V}_{3i} around two orthogonal axis, \underline{V}_{1i} and \underline{V}_{2i} normal to it.

For a unique definition of \underline{y}_{1i} and \underline{y}_{2i} the following convention was used.

$$\underline{y}_{1i} = \underline{j} \times \underline{y}_{3i} \quad (6.1)$$

where \underline{j} is the unit vector in the global y direction.

$$\underline{y}_{2i} = \underline{y}_{3i} \times \underline{y}_{1i} \quad (6.2)$$

For the simple geometry of the figure 6.3, the local orthogonal axis \underline{y}_{1i} , \underline{y}_{2i} and \underline{y}_{3i} are shown on figure 6.4. In every case, with the scheme given above, it is only the \underline{y}_{1i} vector which will change direction when the sense of the thickness vector \underline{y}_{3i} is changed. Thus, only the β displacement will be affected when top and bottom points of a thickness vector are interchanged for different elements. So, the displacements u, v, w, α and β of \underline{y}_{3i} vector of figure 6.4(a) correspond to the u, v, w, α and $-\beta$ of \underline{y}_{3i} vector of figure 6.4.(b). When the necessary transformation is done for β degree of freedom, the assembly process can continue in the usual form.

The transformation requires the multiplication of the entries of the matrix corresponding to β degrees of freedom of the nodes 6,7 and 8 by (-1). Physically, m_{ij} entry of an element matrix represent the force acting on degree of freedom i due to the unit acceleration (or displacement) of degree of freedom j . Transformation, therefore, is achieved by keeping the i (and then j) entry fixed for the β degree of freedom of the reversed node, and multiplying the j (and then i) entries of the whole matrix by (-1).

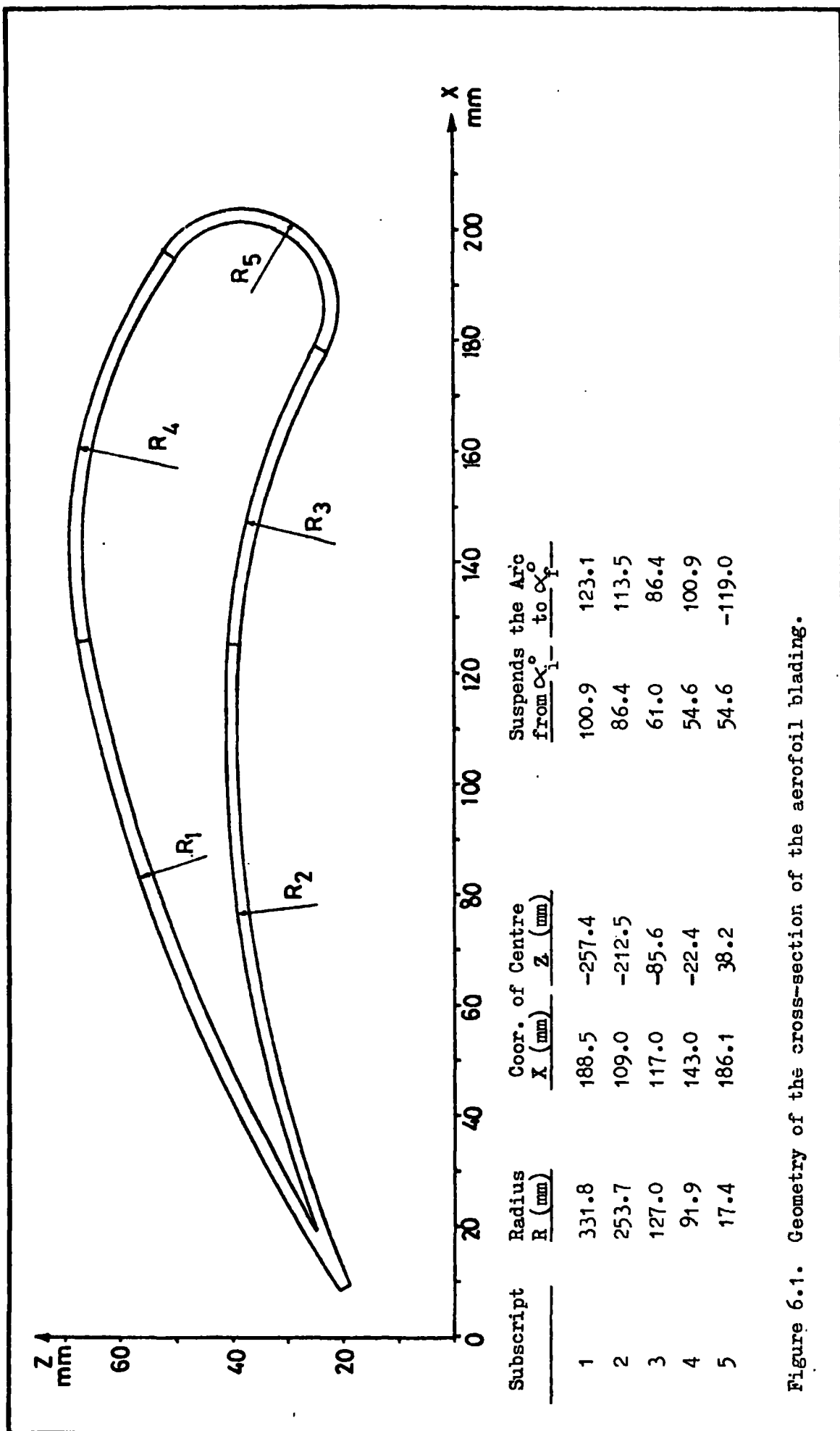


Figure 6.1. Geometry of the cross-section of the aerofoil blading.

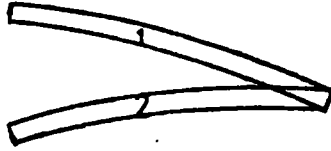


Figure 6.2. Assumed connection of two isoparametric elements at a sharp corner.

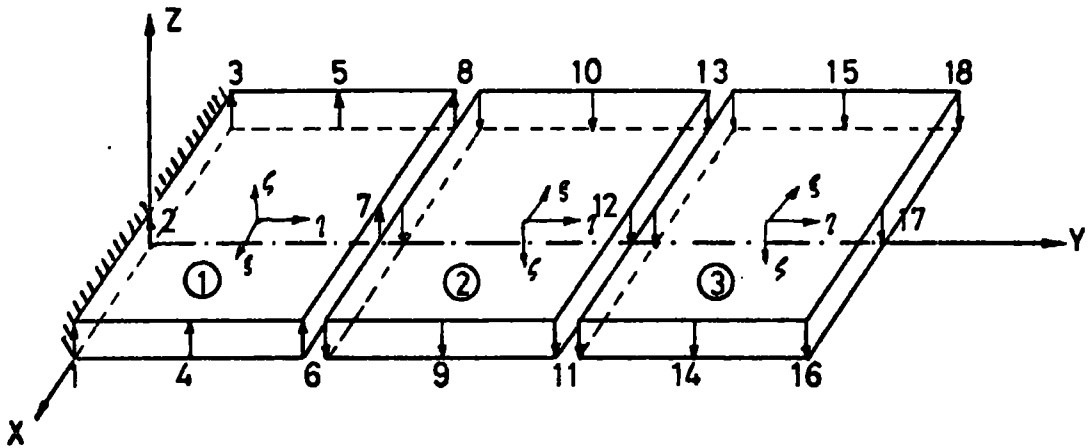


Figure 6.3. Finite element idealization of a rectangular cantilever plate with the thickness vectors mis-matching between elements 1 and 2.

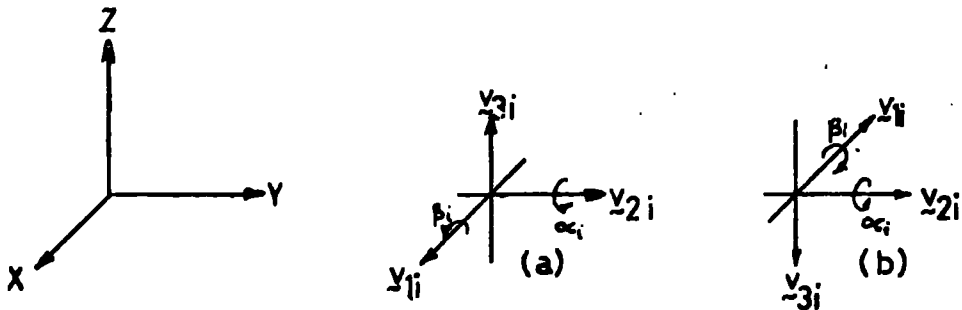


Figure 6.4. Local orthogonal axis and the rotational degrees of freedoms for the elements 1,(a), 2 and 3,(b) of figure 6.3.

6.1.2. V- Shape Cross-Section Cantilever

To assess the validity of the assumption made for the sharp corner connections, the cantilever having the V-shape cross-section, shown on figure 6.5, was studied both experimentally and numerically.

The experimental model was constructed by welding two 200 x 100 x 2.54 mm steel plates together, with an angle of 27° between them. To simulate the clamped boundary condition, it was then welded on a steel block. The determination of the frequencies and the identification of the mode shapes followed the procedure explained in chapter 4.

For the finite element idealization two different meshes, one with two, the other with twelve* elements, were used. The connections of the sharp corner were idealized using both top-to-top and top-to-bottom matching of the thickness vectors. In either assumption the same results were obtained.

The experimental and the numerical results are given on table 6.1. The sketches of some of the typical mode shapes are shown on figure 6.6. The agreement between the results is quite good, and suggests that the assumption made for the sharp corner connections is acceptable.

* This was a uniform mesh with three elements along the length and two across each plate.

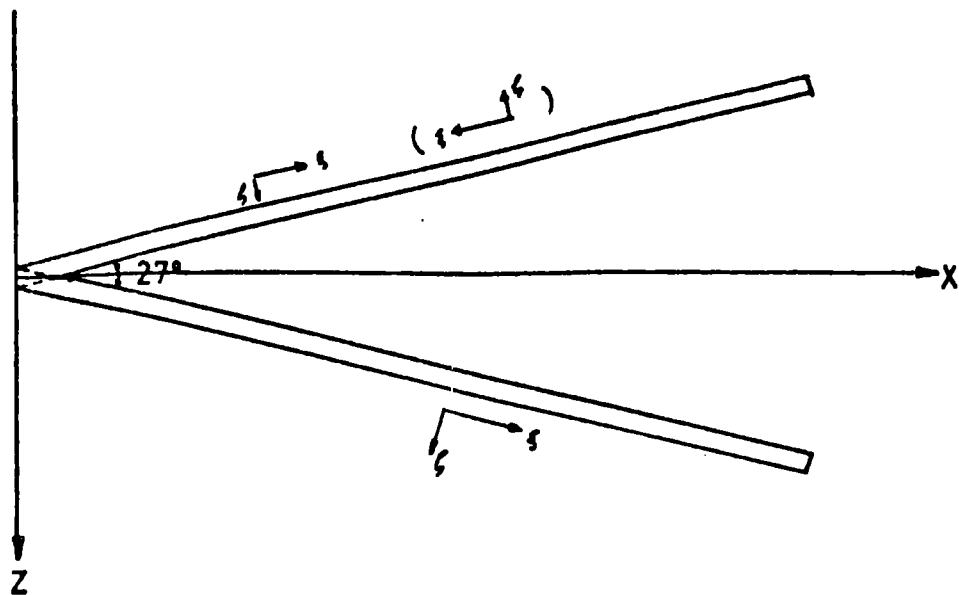
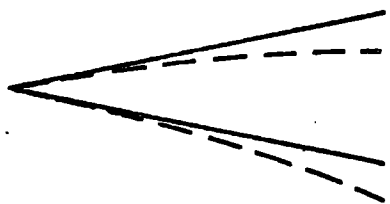
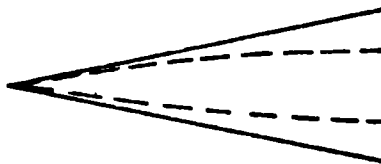


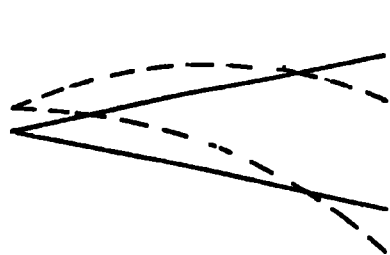
Figure 6.5. Cross section of the cantilever with sharp corner connection.



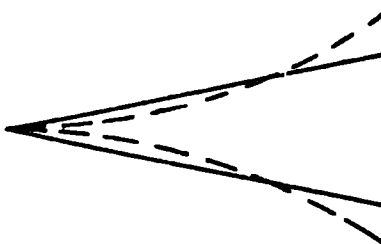
MODE TYPE B



MODE TYPE BR



MODE TYPE T



MODE TYPE 2BR

Figure 6.6. Some of the typical mode shapes of the V- cross-section cantilever.

<u>Mode No</u>	<u>Mode Shape</u>	<u>Frequencies (Hz)</u>		
		<u>Experimental</u>	<u>12 Elem.</u>	<u>2 Elem.</u>
1	B/0	130.9	127.1	134.8
2	BR/0	235.0	254.3	304.9
3	B/1	459.3	471.3	545.7
4	BR/1	545.9	616.8	
5	T/0	772.5	822.9	891.1
6	B/2	1065.9	1156.2	
7	BR/2	1128.2	1252.7	
8	2BR/0	1194.4	1357.0	
9	T/1	1284.5	1385.4	

Table 6.1. Comparison of the experimental and numerical frequencies for the V-Shape cross section cantilever. The number given after / in the mode shapes indicates the number of nodal lines parallel to the root.

6.1.3. Aerofoil Cross-Section Hollow Blading

Sharp corner connection idealization which was studied in the previous section has been applied to two hollow bladings having the aerofoil cross-section shown on figure 6.1. One of the bladings was 400 mm long, shown on figure 6.7, and the other was 200 mm long. Each used a 9x10 mesh.

The first five of the predicted frequencies and the corresponding mode shapes of these blades are given in figures 6.8 and 6.9. The broken lines show the relative displacements of the tip cross-sections. The displacements of a lower cross-section is included if the mode had a nodal line parallel to the base and they are shown by dotted lines. The frequencies given are in Hertz.

The results given in this section are not supported by any experimental analysis, since the foreseen difficulties of manufacturing an experimental model have discouraged any such attempt. Yet, the analyses performed on the individual parts of the structure, like sharp corners, shallow shells or cylindrical segments, have given good results previously. Also, the studies reported in Section 6.4 where the analysis of a hollow turbine blade with a similar geometry gave very good results, imply that the results presented here should be reliable.

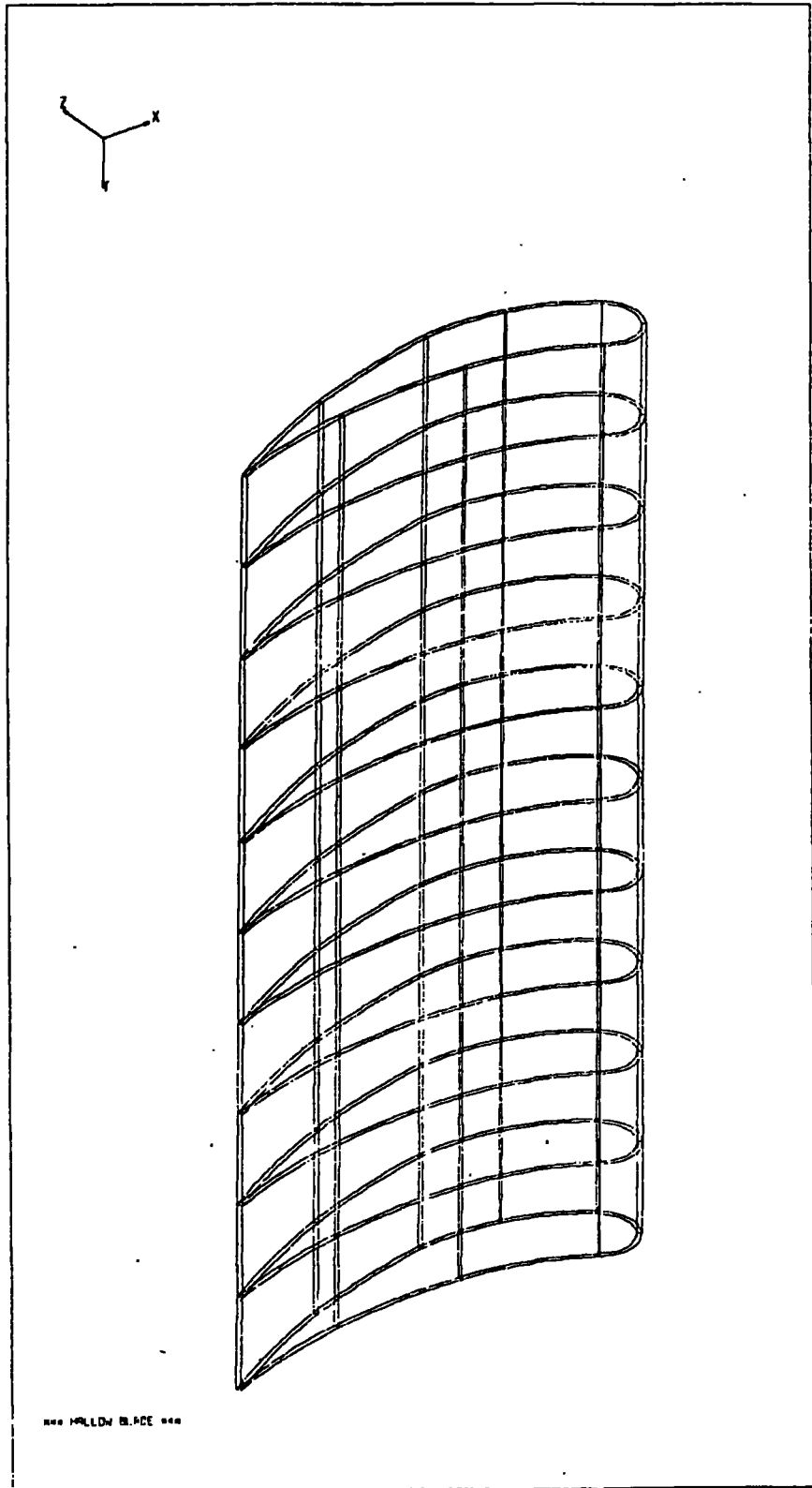
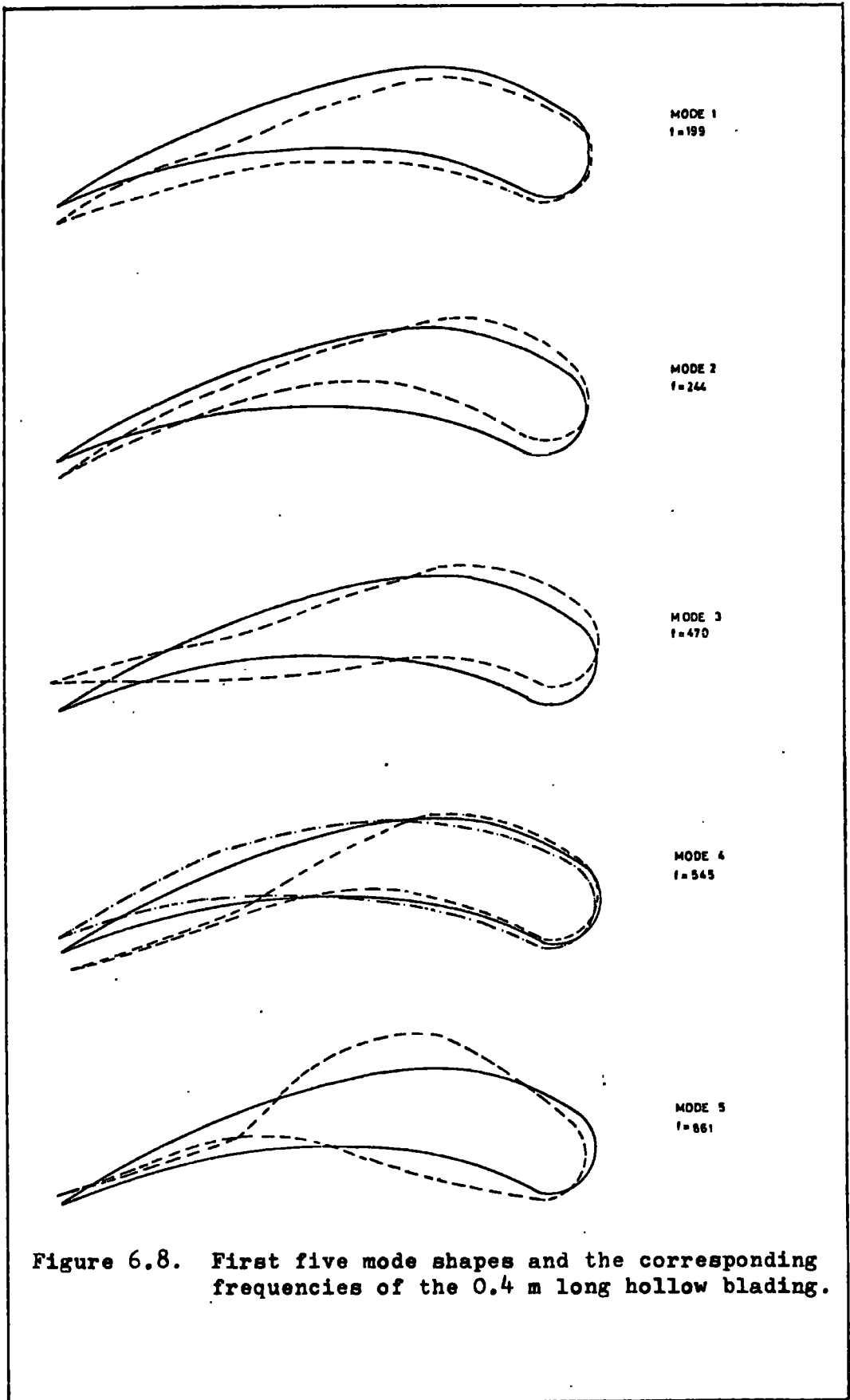
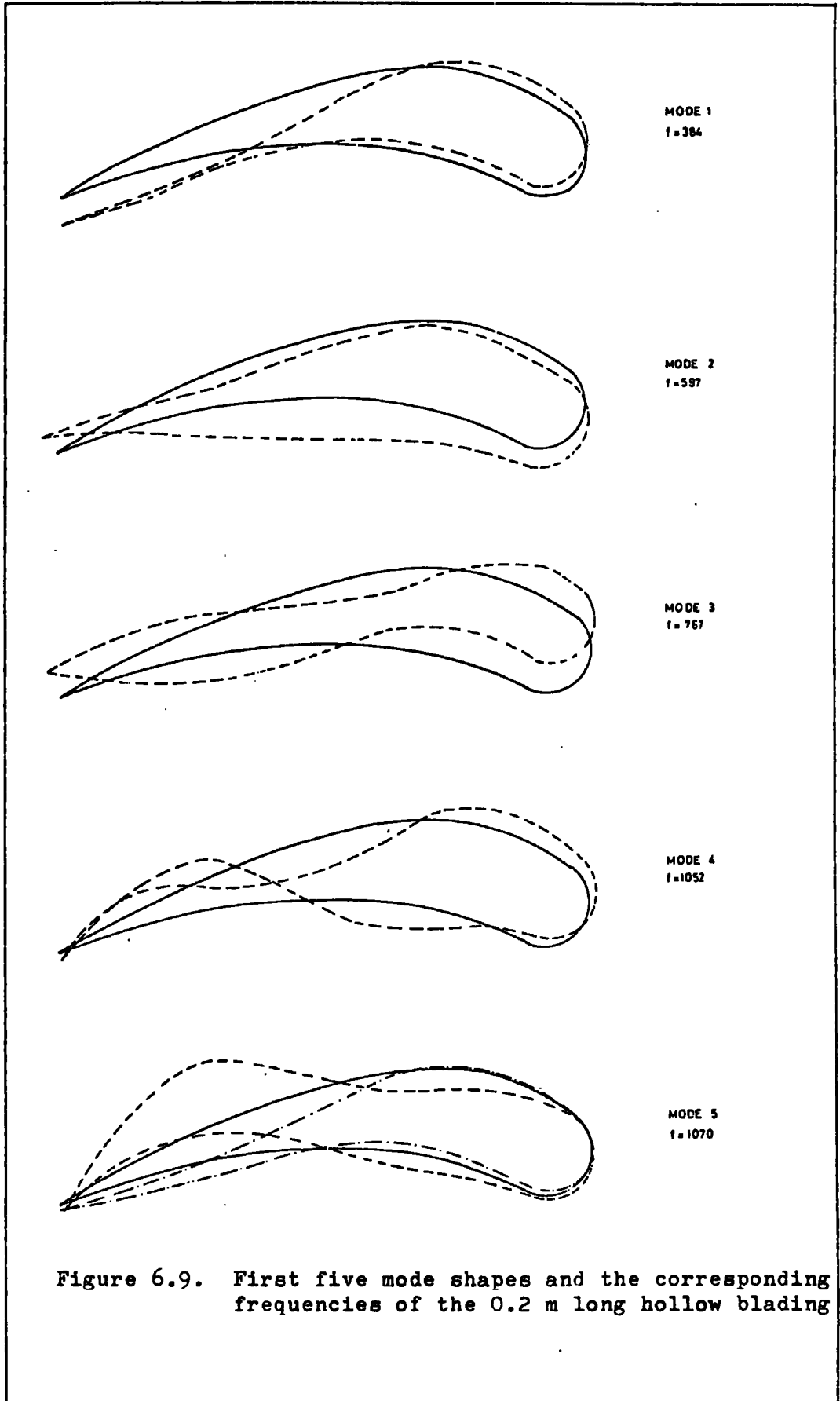


Figure 6.7. Finite element idealization of the 0.4 m long aerofoil cross-section hollow blading.





6.2. Stiffeners

The stiffeners that may be used in a hollow blading are approximated by massless rods connecting the nodes as shown on figure 6.10. The rods are assumed to carry only the axial forces. Thus they have only three degrees of freedom at each end and they do not create any problem in the assembly process.

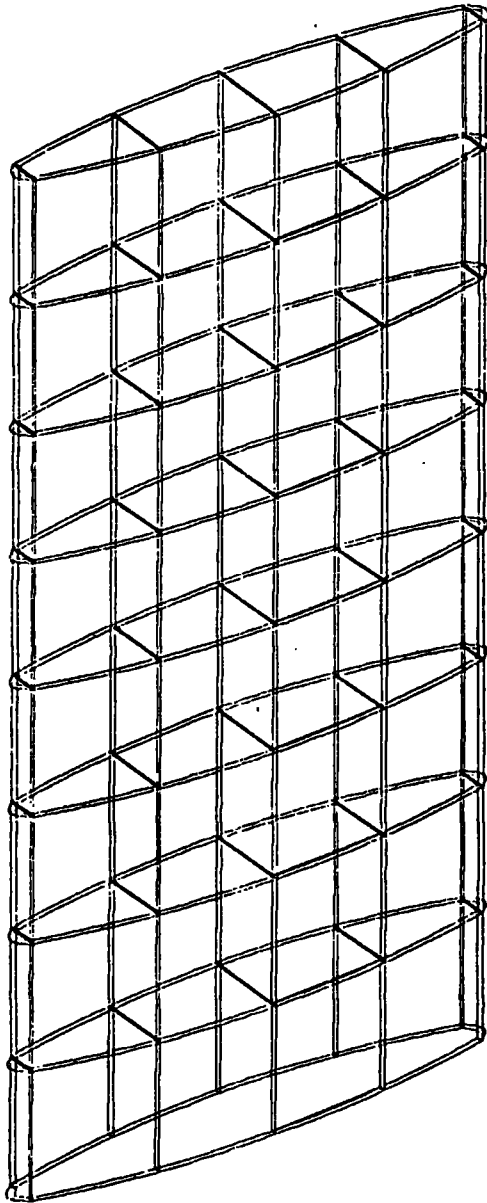
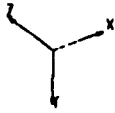
The configuration shown on figure 6.10 was solved with and without stiffeners. The effect of the stiffeners was to suppress the modes which require the relative displacements of the opposite faces. The frequencies corresponding to other modes were identical for both cases.

The only apparent difference was seen at mode 6 (see figures A.32 and A.33), where two modes were coupling. The frequencies corresponding to this mode for the non-stiffened case were 809 Hz and 879 Hz. With the stiffeners the mode became pure in-plane bending with a frequency equal to 866 Hz

6.3. Effect of Pretwist

In section 3.3, the program was used to calculate the natural frequencies of a pretwisted plate, and good results were obtained. In this section, the effect of the pretwist on the natural frequencies and the mode shapes of oval and aerofoil cross-section blades have been considered. The analysis was performed only numerically, and the results still require experimental verification.

Figure 6.10. Hollow Blading with Stiffeners.



*** IDEALIZED GEOMETRY ***

6.3.1. Pretwisted Oval Blading

The oval cross-section blading having the idealized geometry described in Section 5.2.2. was assumed to be twisted linearly about the centroid of its cross-section up to an angle of 60 degrees. The results are presented on figure 6.11 as graphs of natural frequencies in Hz , plotted against the pretwist angle.

The numbers next to each line correspond to the mode number, for the non-twisted blade, as shown on figure A.3.3.

The mode shapes were much influenced by the pretwist. Coupling took place between several modes, and this made it difficult, especially for higher frequencies, to decide the origin of the contributing modes. All frequencies but one, increased with increasing pretwist angle. In some modes the increase was preceded by a slight decline of frequencies at the beginning. In order to reduce the effect of coupling, the same problem was also solved with the stiffeners fitted between the faces. The results are shown with dotted lines on figure 6.11.

The rapid increase of the frequencies was suspected to be due to the stiffening of the deformed elements and this possibility is discussed in section 6.3.3.

6.3.2. Pretwisted Aerofoil Blading

An analysis, similar to the previous one was performed for the aerofoil blading of section 6.1.3. The mode shapes were effected by pretwist even more than the oval blading.

Only two of the modes were identified relatively easily through different pretwist angles, and they are plotted on figure 6.12. This time a rapid decrease of the frequencies was observed. Since a relatively fine mesh was used, the possibility of not being able to integrate the stiffness terms properly due to the reduced integration order (30) seems to be unlikely.

6.3.3. Pretwisted Cylinder

In order to see the effect of the distortions of the elements on the natural frequencies, the cylinder of section 3.5 was reanalysed by assuming different pretwist angles. It was observed that the frequencies corresponding to the simple modes, like bending and torsion, were totally unaffected by the amount of pretwist assumed. For more complex mode shapes some stiffening was observed in the elements, increasing with the complexity of the mode in consideration.

The mode shapes, with a large number of nodal lines and circles, of a circular cylinder are unsuitable to be reproduced by isoparametric shell elements. The difficulty comes from the assumed displacement functions being relatively simple compared with the mode shapes. In order to reproduce a figure similar to the mode shape with the mathematical model, the elements have to be over strained. As a result, the elements become over stiff for certain modes, and to be able to obtain accurate results, one has to use a very fine mesh. When the elements are distorted due to the pretwist, the shape that the elements have to take, in order to describe a mode shape, becomes even

Figure 6.11. Frequencies of oval cross-section blade as function of the pretwist angle.

Broken lines show the coupling modes. Dotted lines are the results with the stiffeners. B/X stands for bending in X direction.

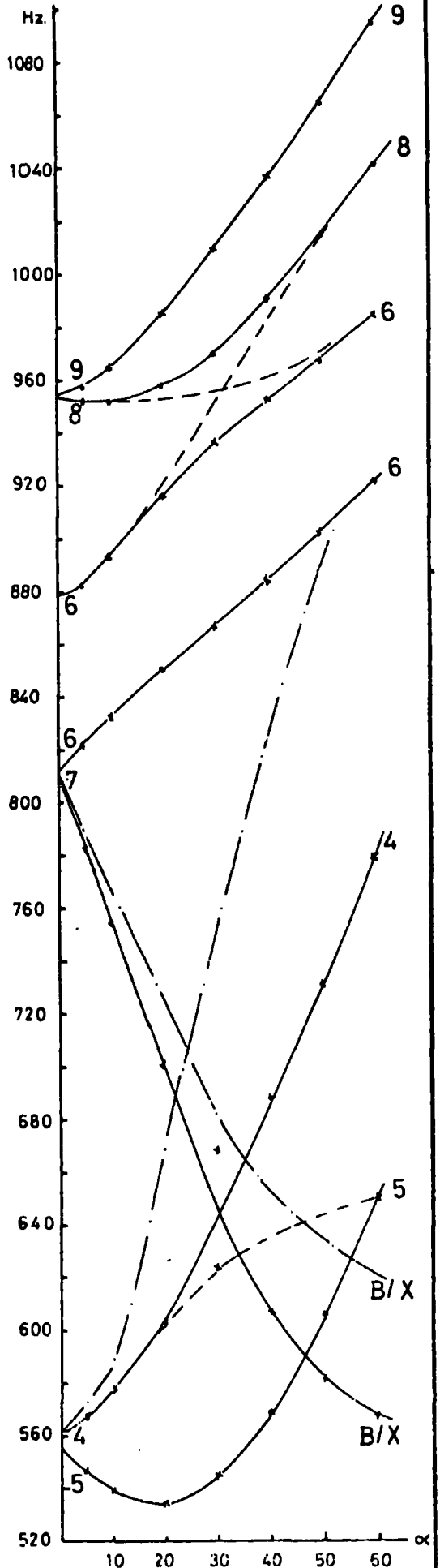
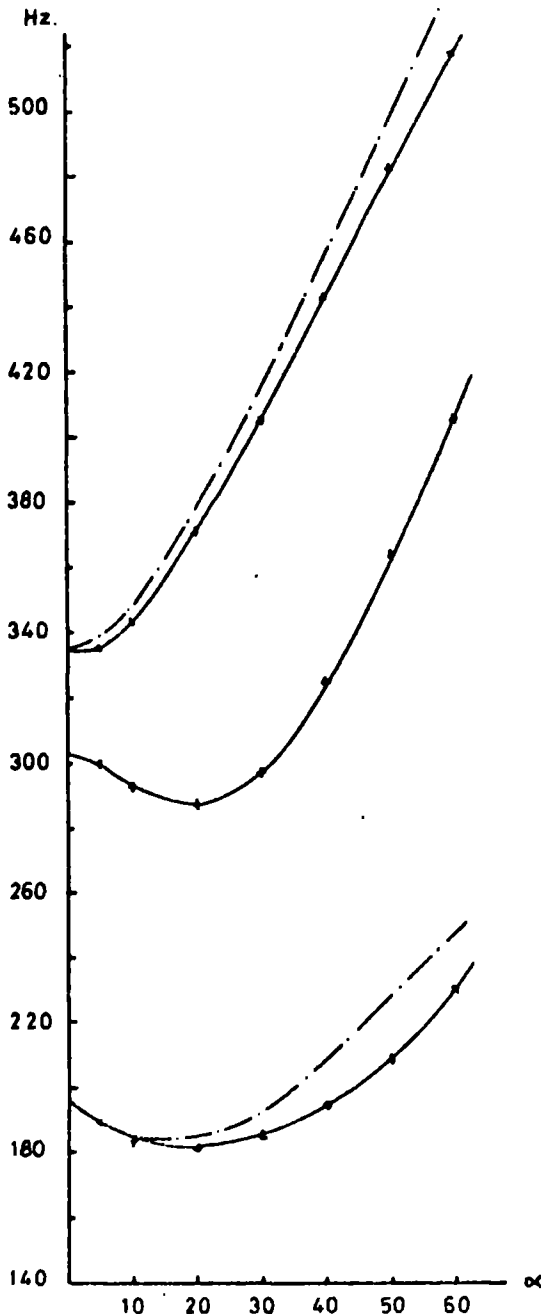
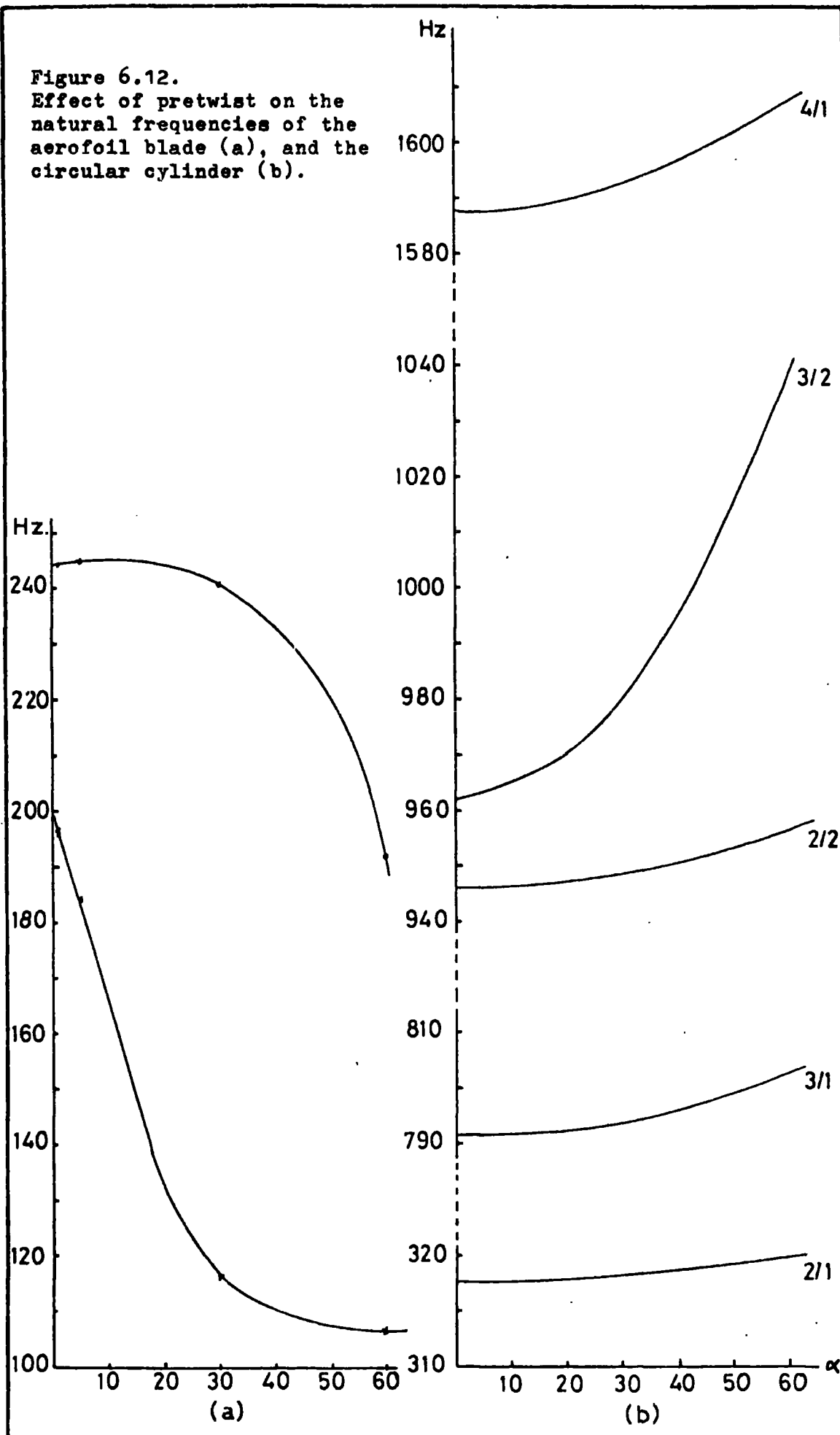


Figure 6.12.
Effect of pretwist on the
natural frequencies of the
aerofoil blade (a), and the
circular cylinder (b).



more complicated. So in fact one should expect the stiffening of the elements with the increasing angle of pretwist. The fact that the bending and torsion modes are unaffected by the pretwist should again be expected because of the regularity of the displacements of individual nodes.

It is clear, from figure 6.12 (b), that the effect of the distortions of the elements on the natural frequencies of simpler modes is very small. The largest difference is only 8% and it corresponds to the mode 3/2. The changes in all the other frequencies are well under 1.5%.

Figure 6.13 shows the computer plotting of the cylinder which was twisted 60 degrees. Figure 6.14 shows the oval blading again with 60 degrees of pretwist angle. The meshes shown are the meshes used in calculations. A close examination of the two figures give the impression that the elements of the oval blading are not distorted as much as the elements of the cylinder. Also, a review of the mode shapes, shown on figure A.3.3, shows that they are much simpler in the sense of the relative displacements of the nodes, than the mode shapes of a cylinder.

As a result of this discussion it may be said that the stiffening of the elements should effect the results of the pretwisted oval blading less than it does the results of the pretwisted cylinder. Still, the lack of a convincing explanation for the common trend of some of the lines in figure 6.11 - i.e. a negative slope followed by a sharp increase - and the lack of any experimental confirmation make it necessary that care should be taken in accepting the results, especially for large pretwist angles, and complicated mode shapes.

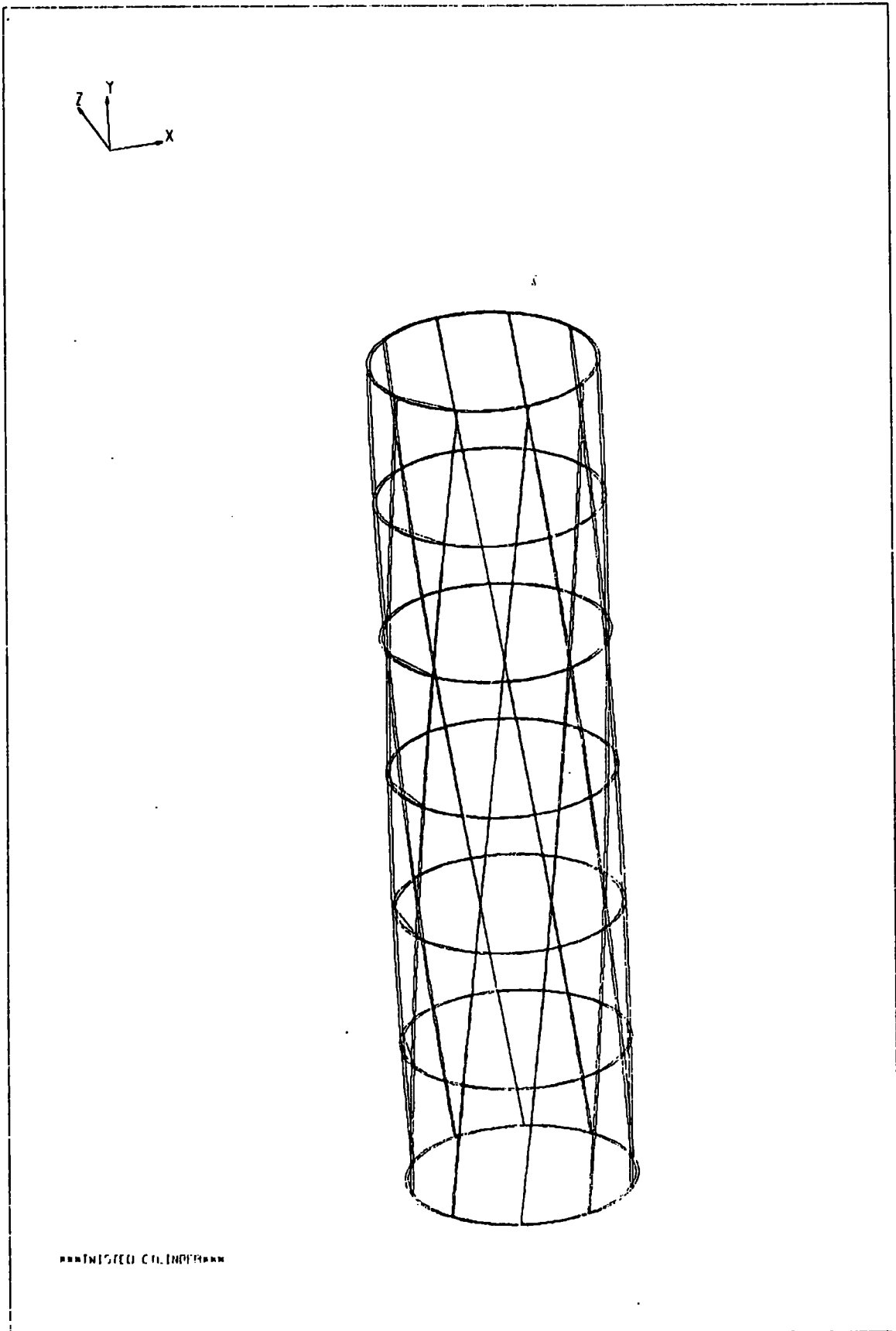


Figure 6.13. Finite element idealization of the cylinder with 60° of pretwist angle.

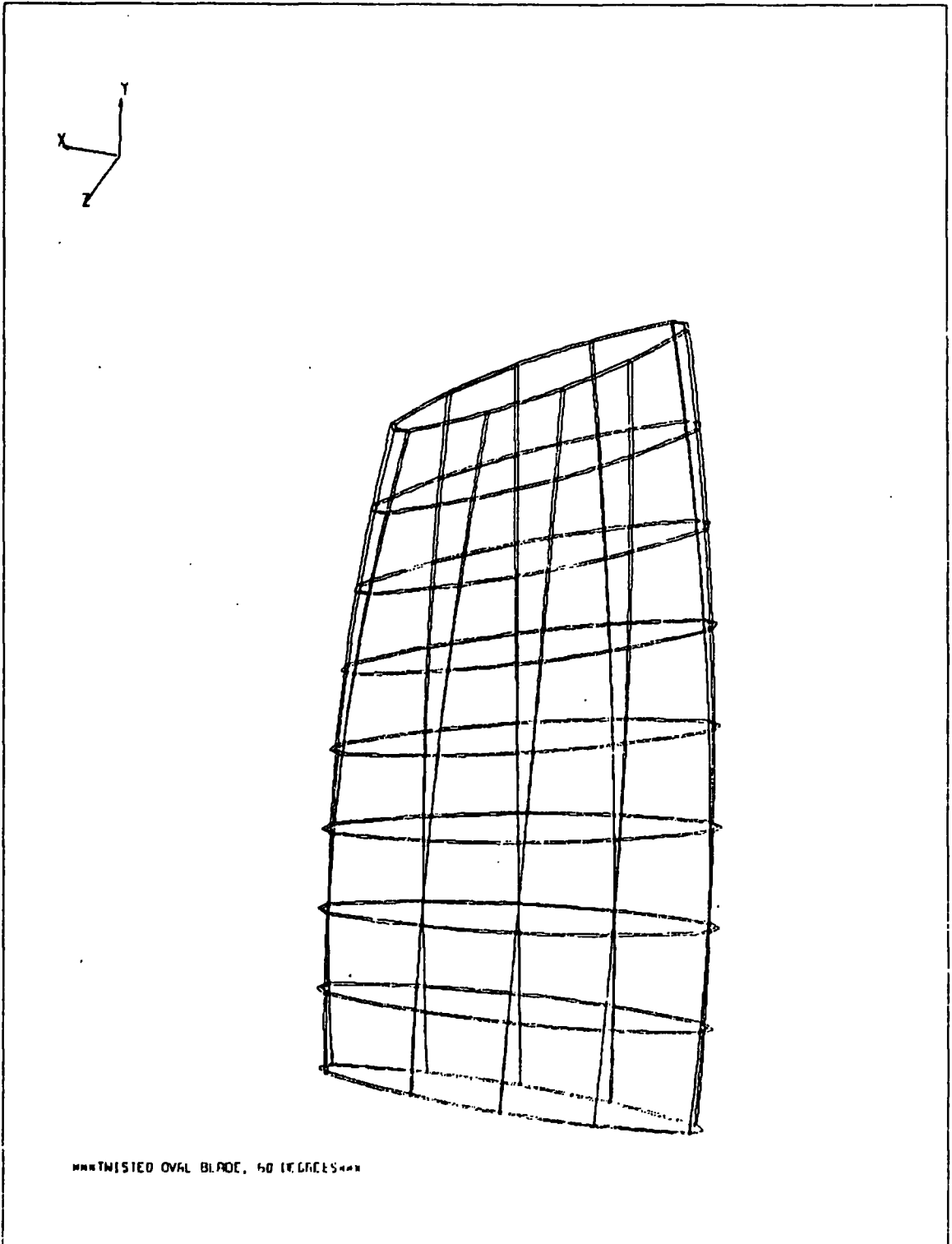


Figure 6.14. Finite element idealization of oval blading with 60° of pretwist angle.

6.4. A Real Turbine Blade

The final exercise was performed on a real hollow turbine blade (supplied by G.E.C.).

The experimental analysis of this blade was done by G.E.C. They encased the root of the blade in lead alloy surrounded by a large mass of cast iron. The method of excitation was electro-mechanical. The natural modes of vibration, between the frequency limits of 0-6 kHz were reported.

For the finite element analysis of the blade the idealization shown on figure 6.15 was used. Nodal coordinates for this idealization were determined by taking direct measurements on the blade. With these measurements it was possible to determine the geometry of the blade only approximately. In the idealization two sharp connections were used, and they were approximated as discussed in section 6.1. The root of the blade was assumed to be clamped at the platform level. The mesh contained 40 elements with 600 degrees of freedoms, 160 of them were retained as masters.

The frequencies and the mode shapes obtained from the finite element analysis are given on figures 6.16 - 6.19, together with the matching experimental results. The positive and negative signs on the figures indicate the transverse displacements of the faces, and "s" shows the stationary regions. The sketches on the left show the pressure side, and those on the right show the suction side. Plottings of the tip and the middle cross section displacements are included in appendix 4.

In spite of the difference between the root fixing of the experimental and the finite element models excellent agreement was obtained for most of the modes. First torsional mode was found to be exactly the same in both analysis. Results for the second mode were very close. Third mode did not exactly meet the experimental mode shape where the areas indicated with negative sign were much lower along the blade length. The big difference seen between the frequencies of the fourth mode was possibly due to the different assumptions of root fixing. Fifth and eighth modes are again in very good agreement both mode shape and frequencywise. Modes seven and twelve show slight differences from the experimental mode shapes. Modes 6,9,10 and 11 have not been detected during the experiments, and one experimental mode found at 3005 Hz was not encountered in finite element analysis.

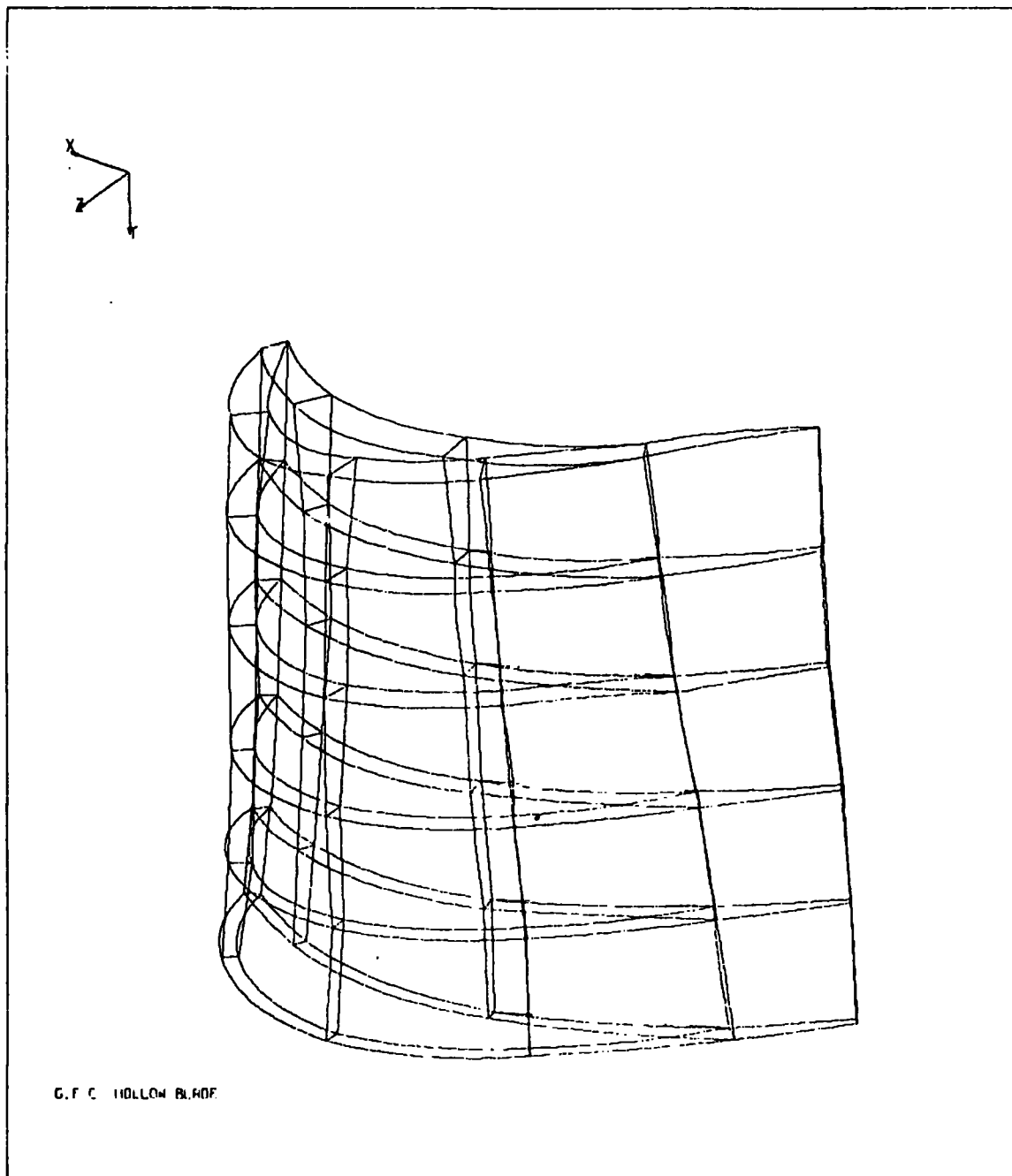
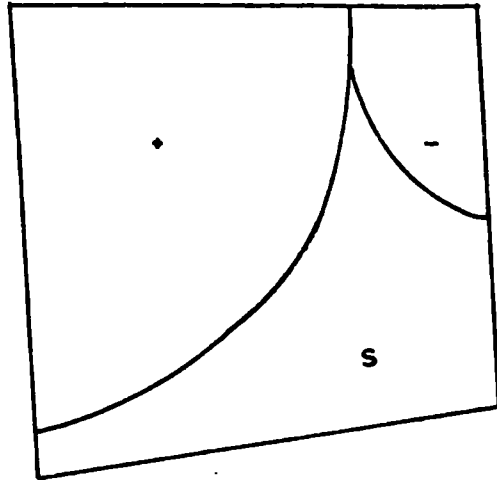
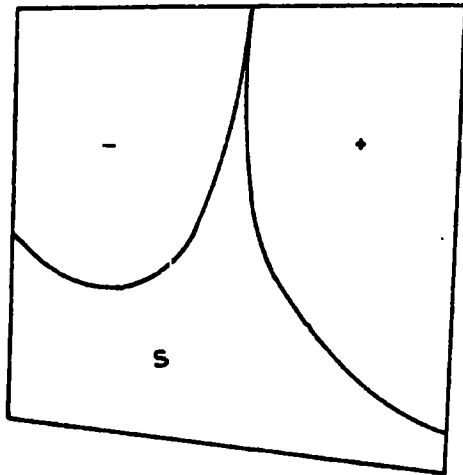
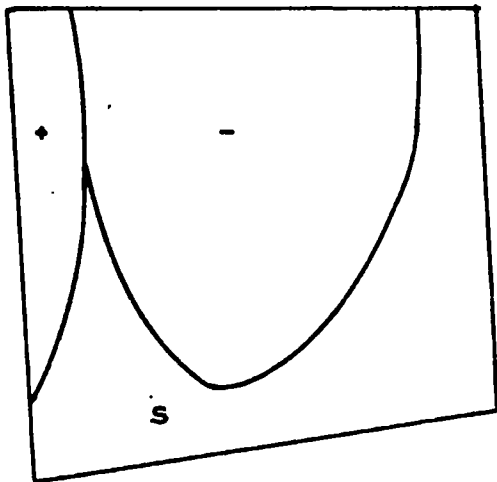
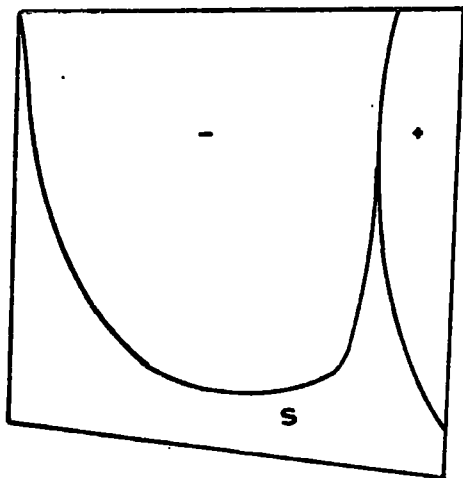


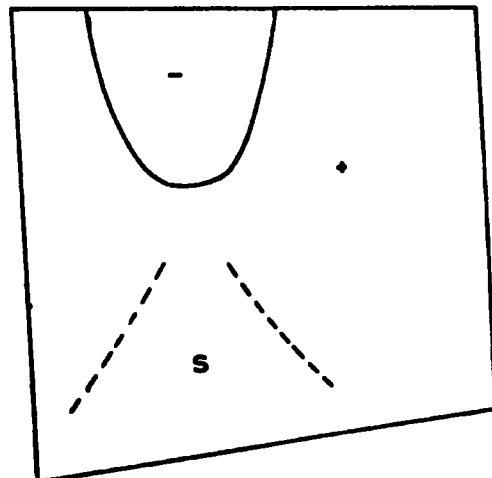
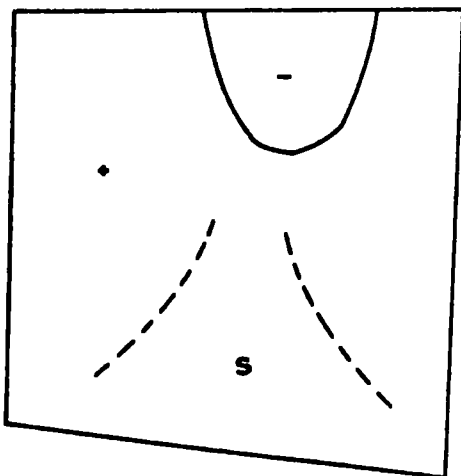
Figure 6.15. Finite element representation of the hollow Turbine Blade.



Mode 1 FE: 1064 EX:1064

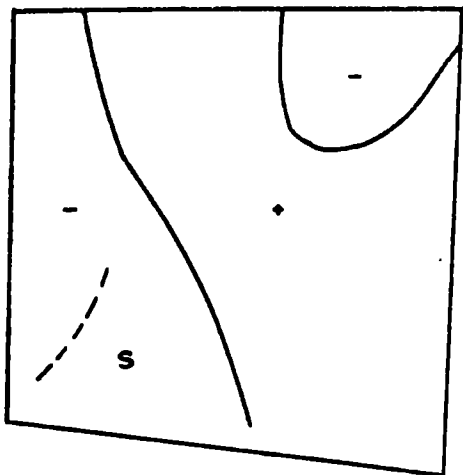


Mode 2 FE: 1713 EX:1717

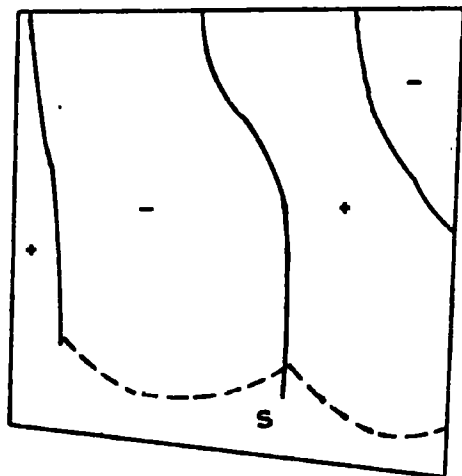
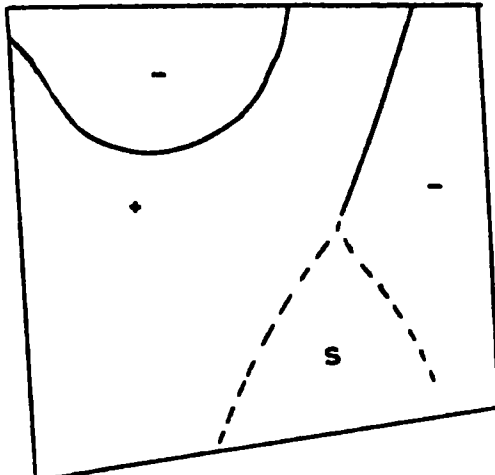


Mode 3 FE:1915 EX:1875

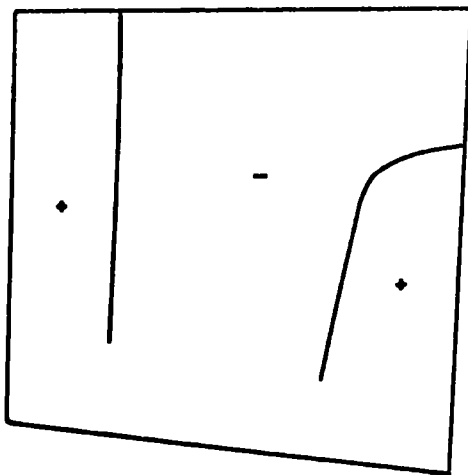
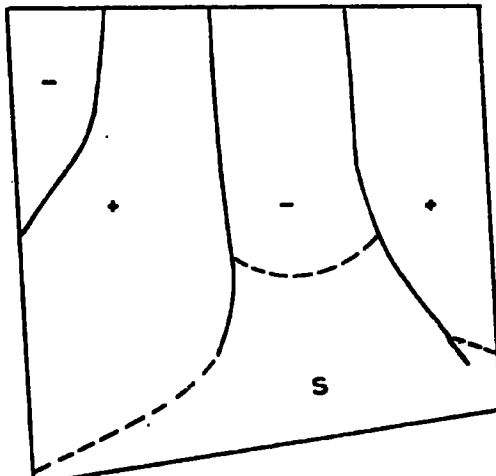
Figure 6.16. Mode shapes and the natural frequencies of the hollow turbine blade.



Mode 4 FE: 2787 EX: 1533



Mode 5 FE: 3310 EX: 3274



Mode 6 FE: 3864

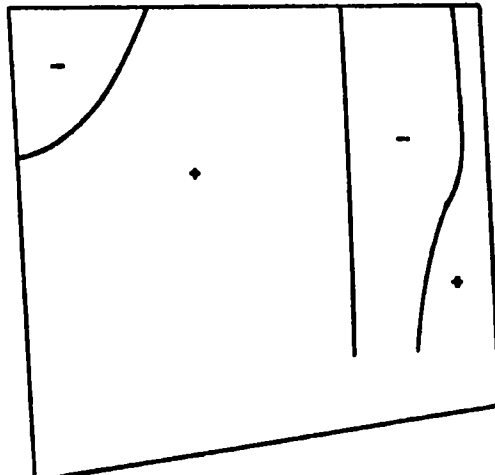
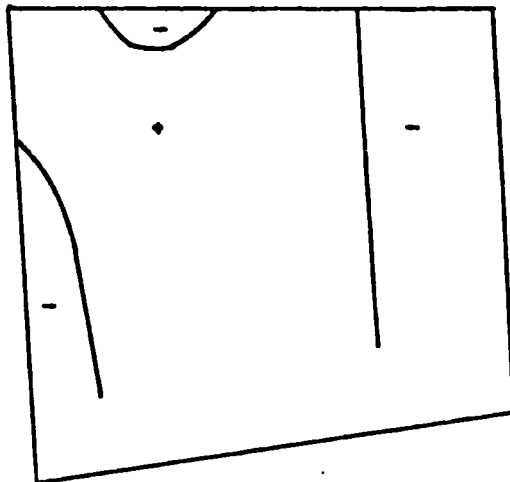
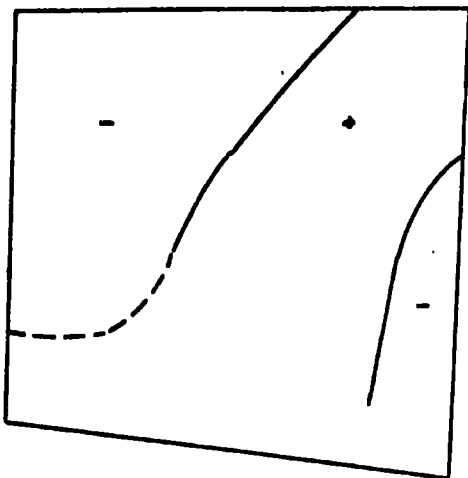
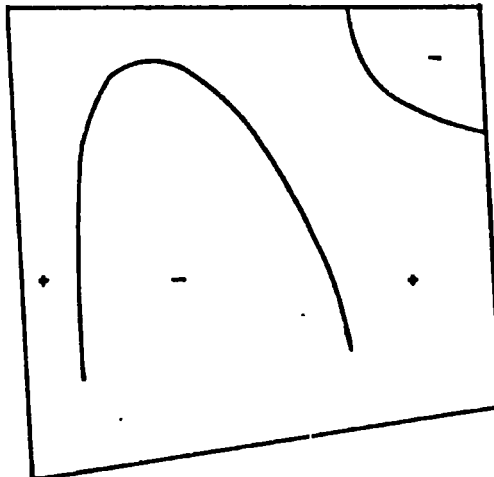
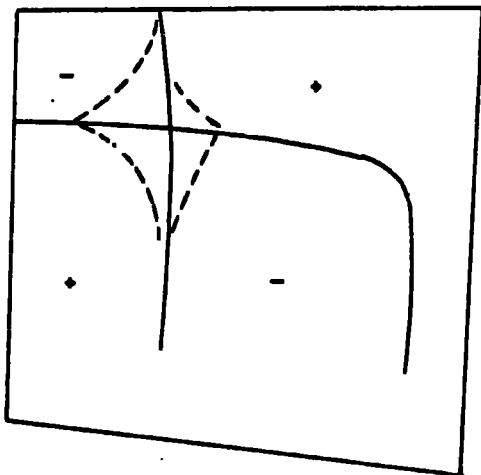


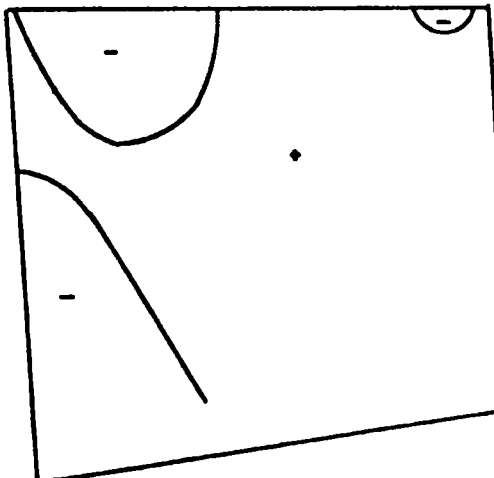
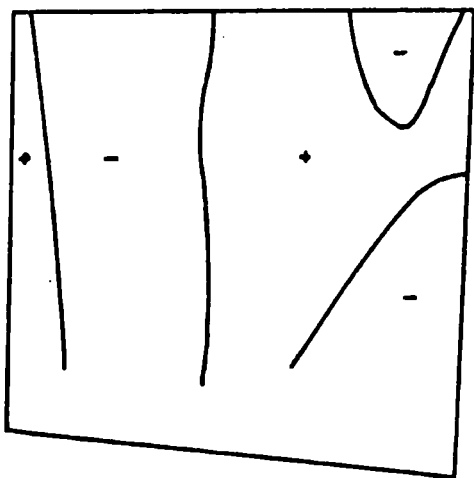
Figure 6.17. Mode shapes and natural frequencies of the hollow turbine blade .



Mode 7 FE: 4209 EX: 4065

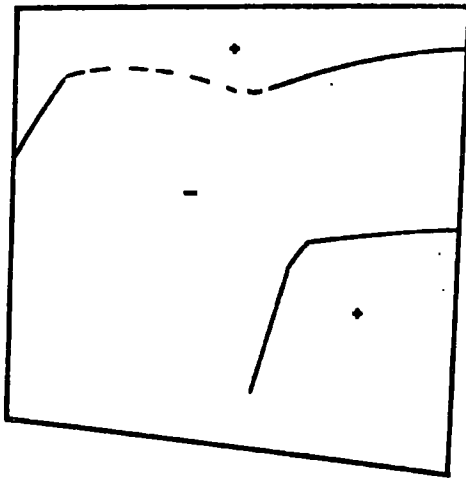


Mode 8 FE: 4414 EX: 4589

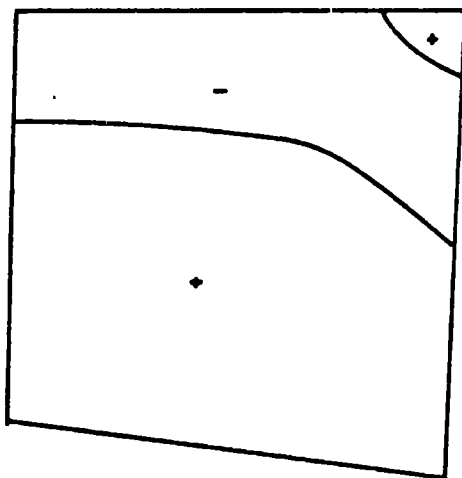
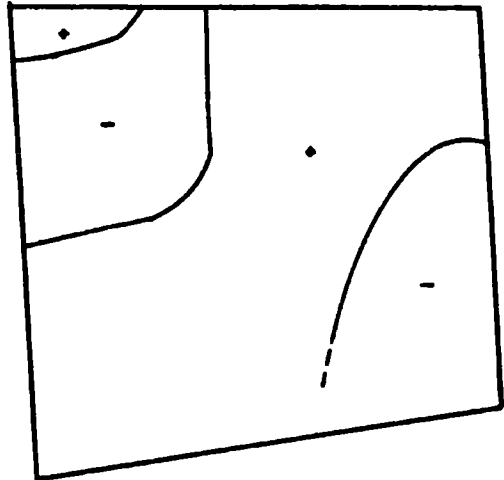


Mode 9 FE: 4658

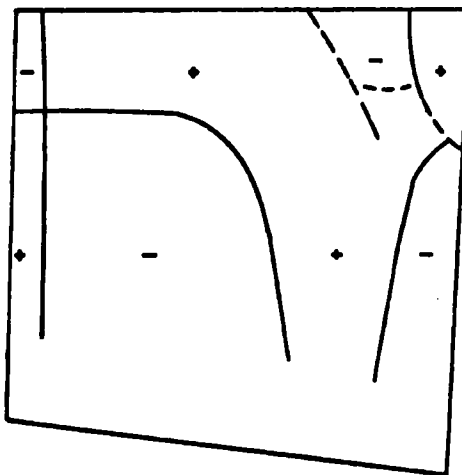
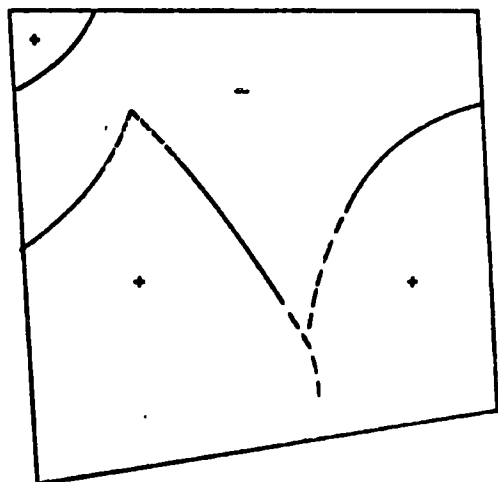
Figure 6.18. Mode shapes and the natural frequencies of the hollow turbine blade.



Mode 10 FE: 5129



Mode 11 FE: 5540



Mode 12 FE: 5919 EX: 5530

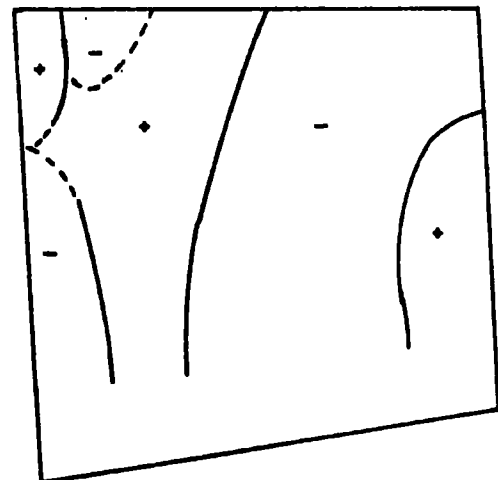


Figure 6.19. Mode shapes and the natural frequencies of the hollow turbine blade.

CHAPTER 7

7. CONCLUSIONS

A numerical and experimental study of the vibrations of arbitrary shell structures is presented. In particular, closed, non-circular cylindrical cantilever shells have been studied, as a preliminary step towards the understanding of the dynamical behaviour of hollow bladings.

Finite element method was employed for the numerical analysis, and a computer program based on the isoparametric thick-shell element has been developed. At the beginning, the element was inefficient in representing sharp curvatures. Later, improved efficiency was achieved by increased the order of integration and with the help of the additional nodes inserted along the curvature. Some assumptions were suggested to represent the sharp corners and multiple junctions in the structure. As a result, the program has become applicable to a large range of vibration problems. It has been used to solve the problems of uniform and variable thickness plates, pretwisted plates and shells, axi-symmetric and non-axisymmetric cylinders, shallow shells and hollow bladings. The results obtained agreed very well with the available experimental and theoretical results of other researchers.

The imperfect geometry of the experimental model of the oval blade was an inevitable result of the difficulties encountered during manufacturing. In spite of that the experimental results agreed reasonably well with the finite element predictions. The difference between the results was within 5 per cent for most of the lower modes. Sensitivity of some of the modes to the geometry and the boundary conditions was observed and the difference in the results for these modes was within 11 per cent.

Experiments have shown that by embedding the blade into a solder base, it is not possible to achieve a root fixing which is stiff enough to simulate the clamped boundary condition. A similar observation was made for brazing which was found to be too flexible to join the components of the oval blading together. Welding, on the other hand, provided stiffer joints both for connecting the components and for fixing the root.

Further experiments with a better experimental model would be interesting to perform for pretwisted and stiffened oval blading.

The method and element used proved to be satisfactory for the dynamic analysis of arbitrary shell structures. The program is well tested and reliable.

REFERENCES

- (1) Ahmad, S., "Curved Finite Elements in the Analysis of Solid, Shell and Plate Structures". Ph. D. Thesis, University College of Swansea, 1969.
- (2) Ahmad, S., Anderson, R.G., Zienkiewicz, O.C., "Vibration of Thick Curved Shells, with Particular Reference to Turbine Blades". J. Strain Analysis, Vol. 5, No. 3, 1970
- (3) Ahmad, S., Irons, B.M., Zienkiewicz, O.C., "Curved Thick Shell and Membrane Elements with Particular Reference to Axisymmetric Problems". Proc. 2nd Conf. on Matrix Meth. in Structural Mech. AFFDL-TR-68-150, part 1 of 2, pp.539-572 1968
- (4) Ahmad, S., Irons, B.M., Zienkiewicz, O.C., "Analysis of Thick and Thin Shell Structures by Curved Finite Elements". Int. J. Num. Meth. Eng., Vol. 2, pp. 419-451, 1970.
- (5) Anderson, R.G., Irons, B.M., Zienkiewicz, O.C., "Vibration and Stability of Plates Using Finite Elements". Int. J. Solids Structures, Vol.4. pp. 1031 - 1055, 1968.
- (6) Appa, K., Smith, G.C.C., Hughes, J.T., "Rational Reduction of Large-Scale Eigenvalue Problems". AIAA. J., Vol.10, No.7, pp. 964-965, 1972.
- (7) Barth, W., Martin, R.S., Wilkinson, J.H., "Calculation of the Eigenvalues of a Symmetric Tridiagonal Matrix by the Method of Bisection". Numerische Mathematik, Vol.9, pp. 386-393, 1967
- (8) Barton, M.V., "Vibration of Rectangular and Skew Cantilever Plates". J. Applied Mechanics, Vol.18, pp.129 - 134, 1951.
- (9) Bossak, M.A.J., Zienkiewicz, O.C., "Free Vibration of Initially Stressed Solids, with particular reference to centrifugal-force Effects in Rotating Machinery". J.Strain Analysis, Vol. 8, No. 4, pp. 245 - 252, 1973.
- (10) Brebbia, C.A., Connor, J.J., "Fundamentals of Finite Element Techniques". Butterworth & Co. Ltd., 1973.

- (11) Bridle, M.D.J., "Vibration of Thick Plates and Shells". Ph. D. Thesis, University of Nottingham, 1973.
- (12) Cantin, G., Clough, R.W., "A Curved, Cylindrical-Shell, Finite Element". AIAA.J., Vol. 6, pp. 1057-1062, 1968.
- (13) Carnegie, W., "Vibrations of Pre-twisted Cantilever Blading". Proc. Inst. of Mech. Eng., Vol.173, No.12, pp 343-374, 1959.
- (14) Cheung, Y.K., Cheung, M.S., "Vibration Analysis of Cylindrical Panels". J. Sound Vib., Vol.22 (1), pp 59-73, 1972.
- (15) Clough, R.W., Johnson, C.P., "A Finite Element Approximation for the Analysis of Thin Shells". Int. J. Solids Structures, Vol.4, pp.43-60, 1968.
- (16) Clough, R.W., Johnson, C.P., "Finite Element Analysis of Arbitrary Thin Shells". ACI publication SP-28, Concrete Thin Shells, pp.333-363, 1971.
- (17) Connor, J.J., Brebbia, C.A., "Stiffness Matrix for Shallow Rectangular Shell Element", Proc. ASCE, Vol. 93, EM5, pp. 43-65, 1967.
- (18) Cowper, G.R., Lindberg, G.M., Olson, M.D., "A Shallow Shell Finite Element of Triangular Shape". Int. J. Solids Structures, Vol.6, pp.1133-1156, 1970.
- (19) Desai, C.S., Abel, J.F., "Introduction to the Finite Element Method". Van Nostrand Reinhold Co., 1972.
- (20) Doherty, W.P., Wilson, E.L., Taylor, R.L., "Stress Analysis of Axisymmetric Solids Utilizing Higher-order Quadrilateral Finite Elements". Structural Engineering Lab. Report No.69-3, Univ. of California, Berkeley, 1970.
- (21) Dungar, R., Severn, R.T., Taylor, P.R., "Vibration of Plate and Shell Structures Using Triangular Finite Elements". J. Strain Analysis, Vol. 2, No.1, pp 73-83, 1967.

- (22) Dunne, P.C., "Complete Polynomial Displacement Fields for Finite Element Method". The Aeronautical Journal of the Royal Aeronautical Society, Vol.72, pp.245-246, 1968.
- (23) Dunne, P.C., "Comments on Complete Polynomial Displacement Fields for Finite Element Method". The Aeronautical Journal of the Royal Aeronautical Society, Vol.72, pp.709-710, 1968.
- (24) Ergatoudis, I., Irons, B.M., Zienkiewicz, O.C., "Curved, Isoparametric, Quadrilateral Elements for Finite Element Analysis". Int. J. Solids Structures, Vol.4, pp. 31-42, 1968.
- (25) Gallagher, R.H., "Shell Elements". World Cong. on Finite Elem. Meth. in Structural Mech., Bournemouth, Dorset, Vol.1., pp E.1 - E.35, 1975.
- (26) Gallagher, R.H., "Problems and Progress in Thin Shell Finite Element Analysis". Finite Elements for Thin Shells and Curved Members, Edited by D.G. Ashwell and R.H. Gallagher, John Wiley & Sons, pp. 1-14, 1976.
- (27) Gill, P.A.T., "Vibrations of Clamped-free Circular Cylindrical Shells", J. Sound Vib. Vol. 25 (3), pp.501-503, 1972.
- (28) Guyan, R.J., "Reduction of Stiffness and Mass Matrices". AIAA. J., Vol.3, No.2, p.380, 1965
- (29) Henshell, R.D., Ong, J.H., "Automatic Masters for Eigenvalue Economization". Earthquake Engineering and Structural Dynamics, Vol.3, pp375-383, 1975
- (30) Hofmeister, L.D., Evensen, D.A., "Vibration Problems Using Isoparametric Shell Elements". Int. J. Num. Meth. Eng. Vol.5, pp.142-145, 1972.
- (31) Irons, B.M., "Eigenvalue Economisers in Vibration Problems". Journal of the Royal Aeronautical Society, Vol. 67, pp. 526-528, 1963.
- (32) Irons, B.M., "Structural Eigenvalue Problems: Elimination of Unwanted Variables". AIAA. J. Vol.3, No.5 pp. 961-962, 1965.
- (33) Irons, B.M., "Engineering Applications of Numerical Integration in Stiffness Methods". AIAA. J. Vol.4, No.11, pp. 2035-2037, 1966.

- (34) Irons, B.M., "Finite Element Techniques in Structural Mechanics, Discussions". H. Tottenham and C. Brebbia (Eds.), Southampton University Press, pp. 328 - 331, 1970.
- (35) Irons, B.M., "Quadrature Rules for Brick Based Finite Elements", Int. J. Num. Meth. Eng. Vol.3. pp. 293-294, 1971.
- (36) Irons, B.M., Razzaque, A., "A Further Modification to Ahmad's Shell Element". Int. J. Num. Meth. Eng., vol. 5, pp.588-589, 1972-73.
- (37) Jones, R.E., Strome, D.R., "A Survey of Analysis of Shells by the Displacement Method". Matrix Methods in Structural Mechanics, Wright-Patterson Air Force Base, Ohio, AFB, AFFDL-TR-80, 1966.
- (38) Kurt, C.E., Boyd, D.E., "Free Vibrations of Noncircular Cylindrical Shell Segments". AIAA. J., Vol. 9., No.2, pp 239-244, 1971.
- (39) Lindberg, G.M., Olson, M.D., Sarazin, A.C., "Finite Element Dynamic Analysis of Shallow Shell Structures". National Research Council of Canada, Aeronautical Report, LR-540, 1970.
- (40) Martins, R.A.F., Owen, D.R.J., "Structural Instability and Natural Vibration Analysis of Thin Arbitrary Shells by Use of the Semiloof Element". Int. J. Num. Meth. Eng., Vol. 11, pp. 481 - 498, 1977.
- (41) McDaniel, T.J., Logan, J.D., "Dynamics of Cylindrical Shell with Variable Curvature". J. Sound Vib. Vol. 19, pp. 39-48, 1971.
- (42) Neale, B.K., "Vibration of Shell Structures". Ph.D.Thesis University of Nottingham, 1971.
- (43) Olson, M.D., Lindberg, G.M., "A Finite Cylindrical Shell Element and the Vibrations of a Curved Fan Blade". National Research Council of Canada, Aeronautical Report LR-497, 1968.
- (44) Olson, M.D., Lindberg, G.M., "Vibration Analysis of Cantilevered Curved Plates Using a New Cylindrical Shell Finite Element". Proc. 2nd Conf. on Matrix Meth. in Structural Mech., Part 1 of 2, AFFDL-TR-68-150, pp.247-269, 1968

- (45) Olson, M.D., Lindberg, G.M., "Dynamic Analysis of Shallow Shells with a Doubly-Curved Triangular Finite Element". J. Sound Vib., Vol. 19 (3), pp.299-318, 1971.
- (46) Pawsey, S.F., "Discussion of Papers by O.C.Zienkiewicz, R.L. Taylor and J.M. Too and S.F.Pawsey and R.W. Clough". Int. J. Num. Meth. Eng., Vol.4, pp.449-450, 1972
- (47) Pawsey, S.F., Clough, R.W., "Improved Numerical Integration of Thick Shell Finite Elements". Int. J. Num. Meth. Eng., Vol.3, pp. 575-586, 1971.
- (48) Petyt, M., "Vibration of Curved Plates". J. Sound Vib., Vol. 15 (3), pp. 381-395, 1971.
- (49) Plunkett, R., "Natural Frequencies of Uniform and Non-Uniform Rectangular Cantilever Plates". J. Mech. Eng. Sci., Vol.5, No.2, pp. 146-156, 1963.
- (50) Ramsden, J.N., Stoker, J.R., "Mass Condensation: A Semi-Automatic Method for Reducing the Size of Vibration Problems". Int. J. Num. Meth. Eng., Vol.1, pp 333-349, 1969.
- (51) Sharma, C.B., Johns, D.H., "Vibration Characteristics of Clamped-free and Clamped-Ring-Stiffened Circular Cylindrical Shell". J. Sound Vib., Vol.14, pp.459-474, 1971.
- (52) Srinivasan, R.S., Bobby, W., "Free Vibration of Non-circular Cylindrical Shell Panels". J. Sound Vib., Vol. 46(1) pp.43-49, 1976.
- (53) Strickland, G.E., Loden, W.A., "A doubly Curved Triangular Shell Element". Proc. 2nd Conf. Matrix Meth. in Structural Mech., Part 1 of 2 AFFDL-TR-68-150, 1968.
- (54) Taylor, R.L., "On Completeness of Shape Functions for Finite Element Analysis". Int. J. Num. Meth. Eng., Vol. 4, pp 17-22, 1972.
- (55) Warburton, G.B., "Dynamics of Shells". Proc. of a Symposium on Structural Dyn., Loughborough University, 1970.
- (56) Wilkinson, J.H., "Calculation of the eigenvalues of a symmetric tridiagonal matrix by the method of bisection". Numerische Mathematik, Vol.4, pp.362-367, 1962.

- (57) Wilson, J.M., Private communication, Durham University, 1978.
- (58) Zienkiewicz, O.C., "The Finite Element Method in Engineering Science", McGraw-Hill, 1971.
- (59) Zienkiewicz, O.C., "Isoparametric and Allied Numerically Integrated Elements - A Review". Numerical and Computer Methods in Structural Mechanics, Edited by Fenves, Perrone, Robinson, Schnobrich, Academic Press 1973.
- (60) Zienkiewicz, O.C., Irons, B.M., Ergatoudis, I., Ahmad, S., Scott, F.C. "Isoparametric and Associated Element Families for Two- and Three-Dimensional Analysis". Finite Element Methods in Stress Analysis, Edited by I. Holond and K. Bell, Tapir, 1970.
- (61) Zienkiewicz, O.C., Taylor, R.L., Too, J.M., "Reduced Integration Technique in General Analysis of Plates and Shells". Int. J. Num. Meth. Eng., Vol.3, pp. 275-290, 1971.



APPENDIX 1

A.1. Explanations on the Program

A.1.1. General Outline.

The computer program developed for the numerical analysis in this thesis consists of four main parts.

- i) Data preparation.
- ii) Construction of element matrices.
- iii) Assembly and reduction.
- iv) Eigenvalue solution.

The program uses three devices. Device numbers 5 and 6 are reserved for input and output respectively. Device number 7 is used to store the topology and the nodal coordinates, so that they can be referred whenever necessary during the computations.

A complete listing of the program is given in Appendix 5.

A.1.1.1. Data Preparation

Preparation of the input data to the program is given in detail in Appendix 2. This data consists of two parts.

The first part contains the control integers, material properties and the geometric parameters.

The second part corresponds to the finite element idealization of the structure. According to the control integers, the topology of the elements and the nodal coordinates are either read in through device number 5, or calculated by the

data generation subroutines. In either case they are stored on device 7.

Data generation is performed by three subroutines called GENOD, GENTOP and DOF (See Section A.1.2). By using them, the amount of the input data can be reduced considerably. They can be employed for any structure having a constant cross-section in one direction, provided that a regular mesh is accepted in that direction. The input data describes the nodal coordinates at one cross-section, and they are repeated (with twisting if desired) along the length at different levels.

A.1.1.2 Construction of Element Matrices

This is simply achieved by programming the formulation given in Sections 2.1 and 2.2 of Chapter 2. Element stiffness and mass matrices are evaluated simultaneously in the main program. After the first evaluation, if the consecutive elements are similar, this step is either left out or replaced by a coordinate transformation procedure (see Section 2.5 of Chapter 2). In either case the computation time is reduced considerably (See Table A.1.1.).

A.1.1.3. Assembly and Reduction

Assembly of the element matrices into the system matrices, and the condensation of the system matrices are performed by subroutine REDUCE. To reduce the required space, only the lower triangles of the system matrices are stored. They are stored row by row in a one dimensional array. Thus, the $(i,j)^{th}$ entry of a matrix is stored in the k th entry of the array, where $k = j + ix (i - 1) / 2$. The assembly procedure follows the

standard form given in the first Chapter of reference (58). Then according to the information supplied with each node, the system matrices are either condensed as explained in Section 2.4, or left as they were.

A.1.1.4. Eigenvalue Solution.

For the solution of the eigenvalue problem of equation (2.1.18) a subroutine available within the Engineering Science Department was used. The routine is based on the technique given in references (56,7). It solves the eigenvalue problem provided that $[K]$ and $[M]$ are real symmetric matrices, and $[M]$ is positive definite.

The routine, first factorizes matrix $[M]$ using Choleski decomposition method, and then combines the $[K]$ and $[M]$ into one matrix, reducing the problem into a standard form. Evaluation of eigenvalues are done by tribisection, and of the eigenvectors by inverse iteration method.

A.1.2. List of Sub-programs

In addition to the main part, the following function and subroutines take part in the program.

- GENTOP:** Generates the mesh for the finite element idealisation.
- GENOD:** Calculates the coordinates of the nodes.
- DOF:** Determines the master degrees of freedoms as instructed by the input data.
- CONEL:** Constructs the elasticity matrix [D]
- INTEGR:** Determines the abscissae and weight coefficients for the Gaussian quadrature according to the number of integration points given in the input data.
- VECI:** Determines the local orthogonal axes v_{1i} , v_{2i} , v_{3i} .
- JAKOB:** Calculates the inverse Jacobian matrix and/or the determinant of the Jacobian.
- THETA:** Determines the direction cosine matrix.
- SHODEN:** This is the only function used in the program. It gives the contribution of node i to the shape function (or its derivative) of any point (ξ, η) .
- REDUCE:** Performs the assembly and condensation of the system matrices (See Section A.1.1.3).
- MAPRIN:** Prints the system matrices which are stored in vector form, after converting them into matrix form.
- MATUNI:** Any matrix given as argument is put into unit form.
- MATCOP:** Copies the first argument matrix [A] into the second argument matrix [B]
- MATNUL:** Matrix [A] given as argument is returned as a zero matrix.
- TRNPOZ:** Takes the transpose of the first argument matrix [A] and stores it into the second argument matrix [B]
- FORMT:** Constructs the transformation matrix, for coordinate transformation.

- TRNSFR:** Performs the necessary transformation when top and bottom coordinates of a thickness vector are defined interchangeably.
- DEIGS:** Solves the eigenvalue problem (See Section A.1.1.4).
- FO1CKF:** A routine from *NAG library of NUMAC. Three matrices A,B,C are given as arguments. Performs the matrix multiplication $[A] \times [B] = [C]$
- FO1AAF:** Another routine from *NAG Library. Calculates the inverse of a matrix.

A.1.3. CPU Time and Storage

The program was written in "FORTRAN IV" computer language using double precision. The machine used for the analysis was IBM 370/168 under MTS. Since the system was using virtual pages for storage, it was unnecessary to use dynamic allocation. The dimensions of the system matrices allow them to expand up to a size of 300 x 300 during the assembly process. The program uses about 250 virtual memory pages which corresponds to just over 1000 K bytes.

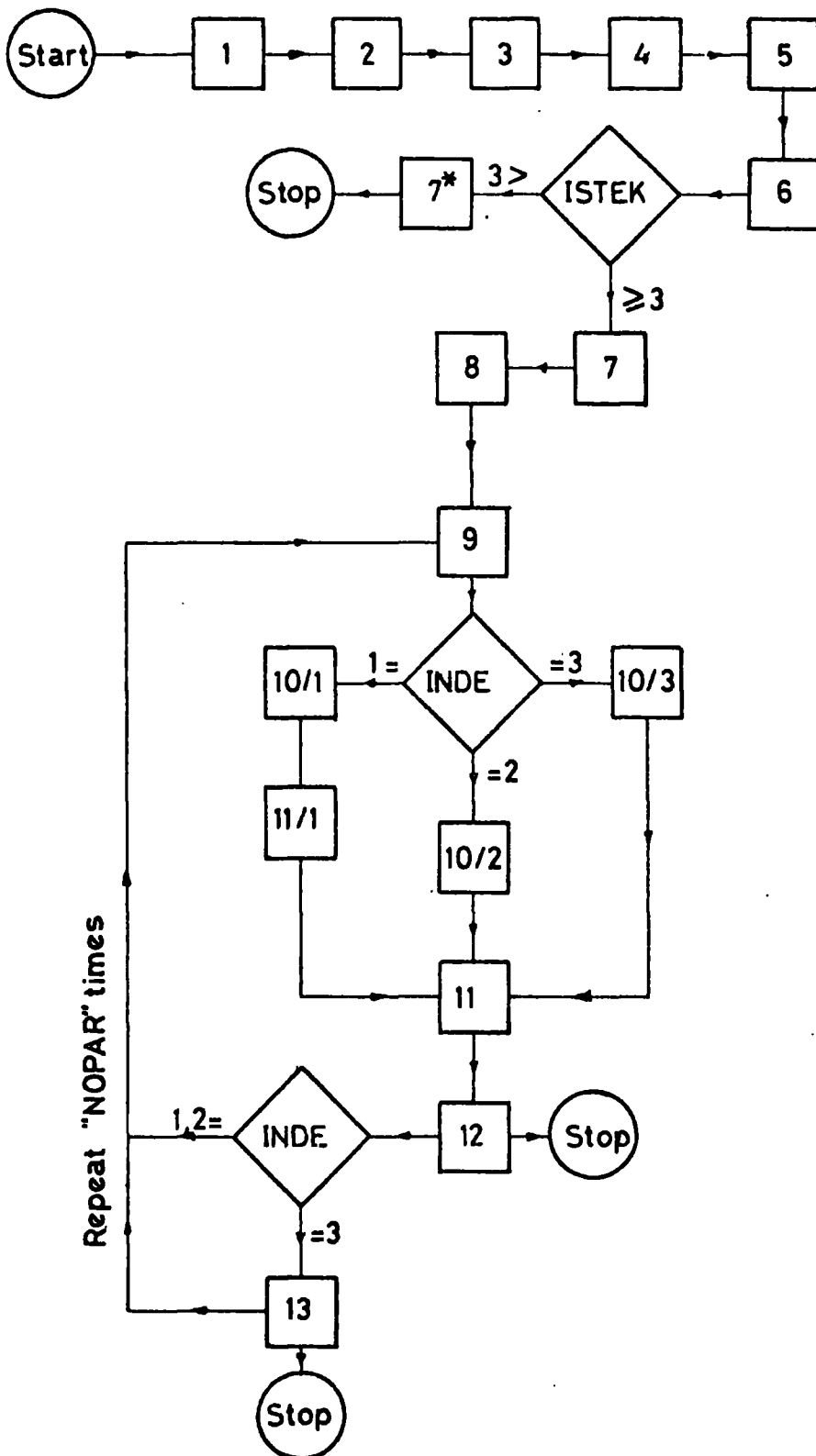
The run time of the program is affected by many variables. In general most of the CPU time is used for the evaluation of element matrices and for the reduction of the size of the system matrices. The time required for reduction depends upon their instantaneous size. By employing a careful node numbering system this size might be reduced considerably.

The following table is included to give a rough idea about the allocation of CPU time used within the program.

CPU Time Used sec.	Operation Performed
1.30	Evaluation of an element stiffness matrix for an 8-node element. Integration points 2 x 2 x 2.
1.92	Evaluation of a mass matrix for an 8-node element. 3 x 3 x 2 integration points were used.
1.67	Transformation of mass and stiffness matrices of one element to use for another element.
2.22	Evaluation of a stiffness matrix for a 10-node element integrated at 3 x 2 x 2 points.
2.16	Evaluation of a mass matrix for a 10-node element integrated at 3 x 3 x 2 points.
0.01	Assembly of one element into system matrices.
0.25	Reduction of the size of the system matrices from 85 to 80
0:42	Reduction of the size of the system matrices from 105 to 100
1.66	Reduction of the size of the system matrices from 180 to 175
3.52	Reduction of the size of the system matrices from 180 to 165
2.86	Reduction of the size of the system matrices from 220 to 215
4.70	Reduction of the size of the system matrices from 220 to 207
10.19	Reduction of the size of the system matrices from 255 to 233

TABLE A.1.1. CPU Time used during difference operations.

APPENDIX 2



DATA PREPARATION CHART

DATA PREPARATION

- 1) Job description card; 1 card; FORMAT (20A4); Columns 1-80.
- 2) Control integers; 1 card; FORMAT (9I5)

Col:	5	10	15	20	25	30	35	40	45
	NO	NONO	NNC	M1	M2	ISTEK	IPRINT	MAPRO	NEY

NO: If $ISTEK < 3$ NO= total number of elements in the structure
If $ISTEK \geq 3$ NO= total number of elements in the first row
of the mesh

NONO: Total number of nodes in the structure.

NNC: Total number of nodes that are to be constrained.

M1: Number of the first eigenvalue to be found.

M2: Number of the last eigenvalue to be found.

ISTEK: Control integer for data generation program (GENDAT)

If $ISTEK < 3$ GENDAT not to be used. All the coordinates of the nodes and master d.o.f. are to be given explicitly.

If $ISTEK \geq 3$ GENDAT is to be used and;

when $ISTEK = 5$, for the elements beyond the first row
ITYPE= 2

when $ISTEK > 5$, for the elements beyond the first row
ITYPE= 3

when $ISTEK < 5$, for the elements beyond the first row
ITYPE is the same as the elements in
the first row.

IPRINT: Output control integer.

- = 1 instantaneous size, frequencies and eigenvectors printed.
- = 2 1 + topology and coordinates printed.
- = 3 2 + element information printed.
- = 4 3 + restrained nodes are printed
- = 5 4 + physical and geometrical properties printed

- 9 > IPRINT \geq 7 5 + Element coordinates, topology, boundary condition, degrees of freedom printed
- = 9 7 + Structural stiffness matrix printed
- > 10 9 + Stiffness matrix for stiffening rods printed.

Recommended value for IPRINT = 8

MAPRO: Number of different material property sets.
 1 \leq MAPRO \leq 5

NEY: Number of elements along the shell.

3) Material properties; MAPRO cards; FORMAT (3 E 10.5)

Col: 10 20 30
 E RO PR

E = Modulus of elasticity

RO = Density

PR = Poisson's ratio

This card is also used to give the sectional properties of stiffening rods. When stiffeners are used, the value of MAPRO is increased by 1 and on the additional card following information is given:

E = Modulus of elasticity of the rods

RO = 0.0 (since the rods are assumed massless)

PR = Cross-sectional area of the rods.

4) Geometric properties; 1 card; FORMAT (2F 10.4)

Col: 10 20
 RLEN THIC

RLEN: Length of Shell

THIC: Thickness of the shell at a reference point.

5) Restrained nodes; NNC/16 cards; FORMAT (1615)

(used to display the restrained nodes, also in GENDAT)

6) Topology and element information; NO cards; FORMAT (14I5,4I2)

Col:	50	55	60	65	70	80
	TOPOLOGY	INTEG	LTIP	ITYP	I PRO	INVER

TOPOLOGY: Topology of the element is given. Nodal numbers must be entered in the order shown in figure 2.1.

INTEG: Three digit integer. Each digit shows the number of integration points for the stiffness matrix of the element in ξ, η, ζ directions respectively. In case of a 2-node element $INTEG = INC$ (see 7)

LTIF: Shows the type of element

- = 2 2-node stiffening bar
- = 8 8-node quadratic element
- = 10 10-node cubic-quadratic element.

ITYPE: Makes repetitive use of the element matrices by transformation.

- = 1 Matrices are calculated independently.
- = 2 Element has the same orientation and geometry and the same material properties as the previous one.
- = 3 Element has different orientation but the same geometry and material properties as the previous one.

I PRO: Element material property set number. $I PRO \leq MAPRO$

INVER: An integer array of dimension 4. It is used when one node is defined in two different ways for two adjacent elements, i.e. when top and bottom coordinates are interchanged. Nodal numbers in the element numbering system corresponding to the nodes which are to be inverted are given in this array. When not in use it is left blank.

7*) Coordinates and d.o.f. of the nodes; NONO cards;
FORMAT(6E10.4, 2I5)

Col:	10	20	30	40	50	60	65	70
	XT	XB	YT	YB	ZT	ZB	NDOF	IDIS

XT, YT, ZT: Top coordinates of the node.

XB, YB, ZB: Bottom coordinates of the node.

NDOF: Total number of degrees of freedom at that node
 $0 \leq NDOF \leq 5$

IDIS: Five digit integer. Each digit represents the d.o.f. u, v, w, α, β respectively. Each digit can take any value between 0-3. Their meanings are:

- 0: Slave d.o.f. It is to be eliminated when contributions of the elements to this node are completed.
- 1: Master d.o.f., to be kept in the eigenvalue problem.
- 2: Restrained d.o.f., not to be included in the structural matrix.
- 3: Slave d.o.f. Indicates that a hinge is assumed at that node.

7) Control integers for GENDAT; 1 card; FORMAT (13I5)

Col:	5	10	15	20	25	30	35	40	45	50-70
	NOPAR	INC	ID	IR	IBC1	IBC2	IS2	MIS1	MIS2	ISAR

NOPAR: Total number of substructures used for the determination of coordinates and topology by GENDAT. (In one substructure there can not be more than 10 nodes along the side).

INC: Increment between first and second level nodes.

ID: Number of levels where there are master d.o.f.

IR: Master degrees of freedoms. IR=IDIS (see 7*) at that level. It can only contain 0's and 1's. 2's and 3's are given separately.

IBC1 : NDOF (see 7*) of constrained end. If all restrained
IBC1 = 0.

IBC2: IDIS of constrained nodes. If all d.o.f. are
constrained then IBC2 = 22222.

IS2: IDIS of hinged nodes. It can not contain 1 or 2.
If no hinges, leave blank.

MIS1: Only needed when making use of the symmetry. Then
it equals to NDOF of the nodes along the axis of
symmetry.

MIS2: IDIS of nodes along the axis of symmetry. It can
only contain 0's and 2's.

ISAR: An integer array of dimension 4. If the number of
any node along the reference side of the substructure
is put in this array, its XT and XB coordinates are
corrected as $XT = XB = (XT+XB) / 2.0$.

8) Angle of pretwist and master d.o.f. levels; (1+ID)/8 cards;
FORMAT (8 E 10.4)

Col.	10	20	80
	TOTAN	WS(1)	WS(7)
		WS(8).....	WS(ID)

TOTAN: Total angle of pretwist in degrees. The structure is
twisted linearly along the length around the coordinate
centre. If no pretwist TOTAN = 0.0.

WS: An array containing the y-coordinates of the levels
(along the length) where there are master d.o.f.

RAD: Outer radius of curvature of the segment.

THIC: Thickness of the segment.

AINI: Angle(in degrees) between the positive X-axis and the radius which connects the first node on the segment to the centre. Angle is measured +ve from X to Z following the shortest path.

AFIN: Angle (in degrees) between X-axis and the radius which connects the final nodal point to the centre.

NOPT: Number of nodal points between AINI and AFIN.

NTERS: See 11/1.

10/3) Control integers for general cylindrical geometry;
1 card; FORMAT (2I5).

Col: 5 10

 NOPT NTERS

NOPT: Number of nodal points along the reference line of the prismatic section.

NTERS: See 11/1.

11) Node numbers; 1 card; FORMAT (10I5)

Node numbers, given in order, between the initial and final nodes of substructure. There are $NOPT \leq 10$ such points.

12) Topology indication number; 1 card; FORMAT (10I5)

Gives the midside nodes, next to the nodes given in 11.

0 : no midside node next to the node given in 11.

n : number of the midside node next to the node given in 11.

-n : nodes along this node are hinged or along the axis of symmetry.

-1 : coordinates of this node are calculated in another substructure.

13) Geometry of the general cylindrical section; NOPT cards; FORMAT (4 E 10.4).

Col:	10	20	30	40
	XATI	XABI	YATI	YABI

coordinates of each node in the reference plane of substructure are given in the order from initial to final.

Care must be taken on the following points:

- 1) The coordinate system chosen must be right handed.
- 2) When data generation subroutines are to be used, choose global y-direction along the length of the structure.
- 3) Top and bottom points of a node are decided according to ζ and corresponding global coordinates must be in accordance.
- 4) Symmetry and hinge facilities in GENDAT can only be used if the axis of symmetry or the hinges are on a line parallel to one of the global coordinates.
- 5) To use hinge facilities, element numbering system must go faster in the direction of hinged line.
- 6) Symmetry and hinge condition in GENDAT can not be used for the same run of the program.

APPENDIX 3

This appendix contains some additional figures to Chapter 5. Also a comparison of these figures is included.

Some abbreviations are used for the identification of difference mode shapes. In this convention IG stands for idealized geometry, similarly RG and EX for real geometry and experimental. For instance, IG-8 indicates the 8th mode given in idealized geometry finite element analysis.

List of Figures

<u>Figure No:</u>	<u>Notes:</u>
A.3.1 (a), A.3.1. (b)	Tip deflections for two frequencies coupled at mode IG-6.
A.3.2.	Key to the deflection curves. Shows the location of the numbered lines of the deflection curves in X-Z plane (X-section).
A.3.3 (a) - A.3.3 (c).	Sketches of the mode shapes obtained using I.G., with 102 "w" master d.o.f. (92 u + 92 w) were used for mode 6.
A.3.4 - A.3.19.	Deflection curves for some of the ideal modes obtained by using 102 master d.o.f. The curves are drawn only for the nodes on the top face. The deflection curves of the corresponding nodes on the bottom face are either the same (S) as the top ones, or the mirror image (MI) of them. This is indicated in the figures with the abbreviations (S) and (MI).
A.3.20(a) - A.3.20(c).	Sketches of some of the mode shapes obtained for real geometry, using 90 (w), 129 (w) and 146 (w + u) master degrees of freedoms.
A.3.21 - A.3.31.	Deflection curves for some of the mode shapes of figure A.3.20, using 129 master degrees of freedoms. The nodes of the top and bottom faces are shown separately.
A.3.32(a) - A.3.32(b)	Sketches of the mode shapes determined experimentally.

Comparison of the Modes

<u>I.G. Mode No:</u>	<u>Comparison with R.G. and EX.</u>
1 - 4	Mode shapes match in all cases perfectly.
5	Very little coupling with 2nd torsion is seen in experimental and RG finite element analyses. Frequency increases with respect to the I.G. analysis.
6	Experimental mode determination was difficult, probably requires a better excitation. Effect of imperfect geometry can be seen in R.G. mode shape. Frequencies agree in all three analyses.
7	Difference between the frequencies of I.G. and R.G. analyses is similar to the difference in first bending mode. Experimentally a frequency was detected at 708.5 Hz, identification of the mode shape was not possible, accepted as 2nd bending mode.
8/9	Frequencies are very close. EX-7/8 and R.G.-8/9/10 have similar mode shapes, possibly coupling of I.G.8 and I.G.9. Deflection curves of R.G.9 are very similar to the deflection curve of I.G.9.
10	A remarkable resemblance between EX-10 and R.G.13/14. Equivalent I.G. mode is I.G.10. Effect of choice of degrees of freedom can be seen by comparing R.G.13 and R.G.14.
11	R.G. 11/12 are quite similar to EX - 9, and their equivalent is I.G.11. R.G.11 and R.G.12 represents the same mode, slight difference is due to the different master degrees of freedoms.

I.G.
Mode No.

Comparison with R.G. and EX

- 12/13 EX-11 and EX-12 have more or less the same shape which is similar to R.G.15. Although sketches of I.G.12 resemble to this shape, deflection curve for I.G.13 and R.G.15 show the same pattern. A coupling between the two might have taken place.
- 15 EX-13 and I.G.15 are quite similar, no corresponding mode in R.G. group.
- 16 Deflection curves are similar to R.G.18's, no corresponding mode identified experimentally.
- 18 Examination of the deflection curves show a resemblance with R.G. 17.
- 19 EX-15 and R.G.19 have similar mode shapes. The only possible mode corresponding to them in I.G. group is I.G.19.

Similar frequencies pointed above are listed on table A.3.1.

Ideal Geometry		Real Geometry		Experimental	
Mode No.	Frequency Hz	Mode No.	Frequency Hz	Mode No.	Frequency Hz.
1	196.5	1	158	1	142
2	302.5	2	306.8	2	297
3	335.5	3	371.5	3	355
4	562	4	588.4	4	561
5	553	5	648	5	598
6	821	6	837	6	812
7	804.5	7	739		708
8/9	954/950	8/9/10	979/1027/1004	7/8	905/935
10	1042	13/14	1155/1249	10	1096
11	1047	11/12	1052/1109	9	980
12/13	1062/1191	15	1299	11/12	1136/1157
15	1419	-	-	13	1313
16	1502	18	1638	-	-
18	1598	17	1541	-	-
19	1733	19	1818	15	1633

Table A.3.1. Frequencies corresponding to similar mode shapes in three different analysis.

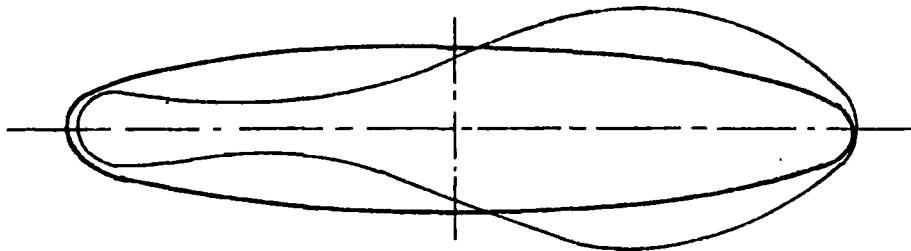


Figure A.3.1.(a): Tip deflection for mode 6 at frequency = 809 Hz.

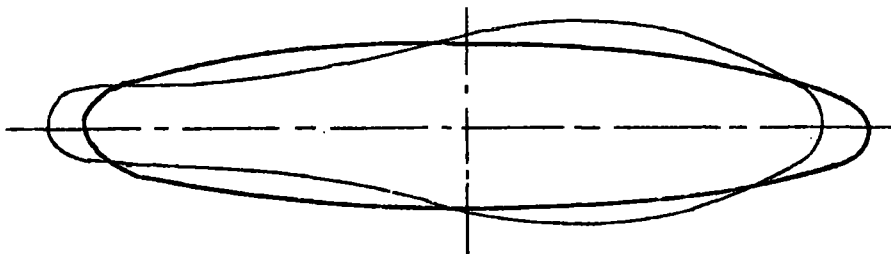


Figure A.3.1.(b): Tip deflection for mode 6 at frequency = 879 Hz.

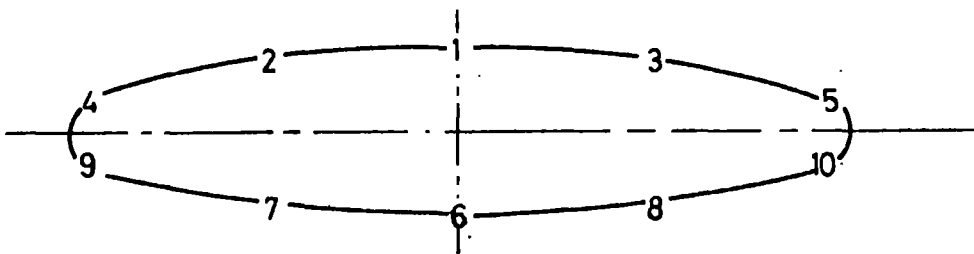


Figure A.3.2: Key to the deflection curves.

Mod. No.	Freq. Hz.	Top Face	Bottom Face	Tip section	NNC
1	196.5				0
2	302.5				0
3	335.5				0
4	562				1
5	553				1
6	809 / 879				0
7	804.5				1

Figure A.3.3.(a): Predicted ideal mode shapes of the Oval X-Section blade using the idealized geometry. NNC stands for number of nodal circles.

Mod. No.	Freq. Hz.	Top Face	Bottom Face	Tip section	NNC
8	954				0
9	950				1
10	1042				2
11	1047				2
12	1062				1
13	1191				2
14	1359				2

Figure A.3.3.(b).

Mod No.	Freq Hz.	Top Face	Bottom Face	Tip section	NNC
15	1419				0
16	1502				3
17	1583				3
18	1598				3
19	1733				0
20	1513				?

Figure A.3.3.(c).

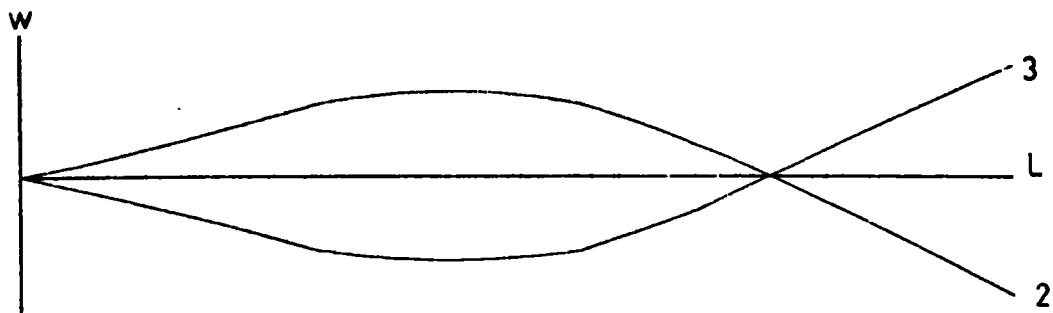


Figure A.3.4: Mode 4, (S), $f = 562$ Hz.

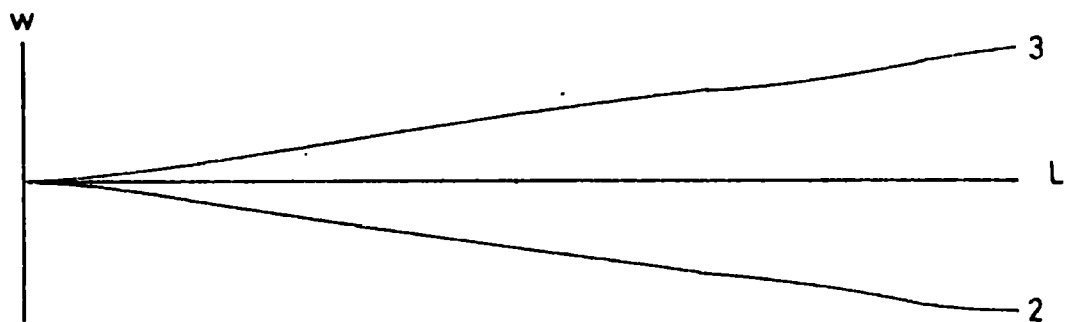


Figure A.3.5: Mode 6, (MI), $f = 821.5$ Hz.

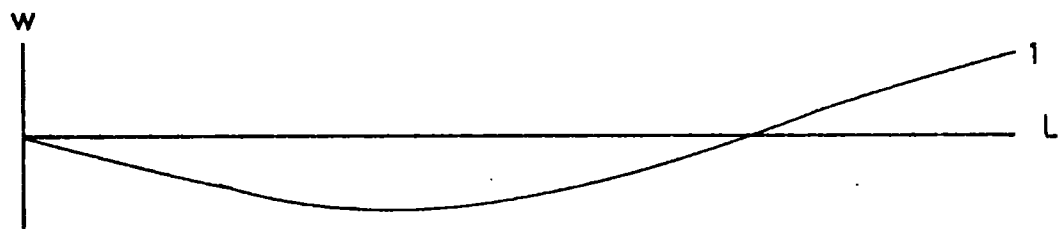


Figure A.3.6: Mode 7, (S), $f = 804.5$ Hz.

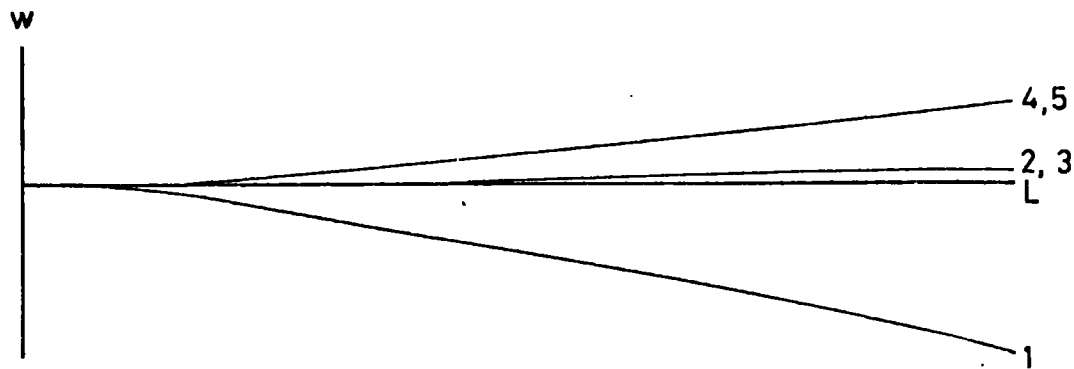


Figure A.3.7: Mode 8, (MI), $f = 954$ Hz.

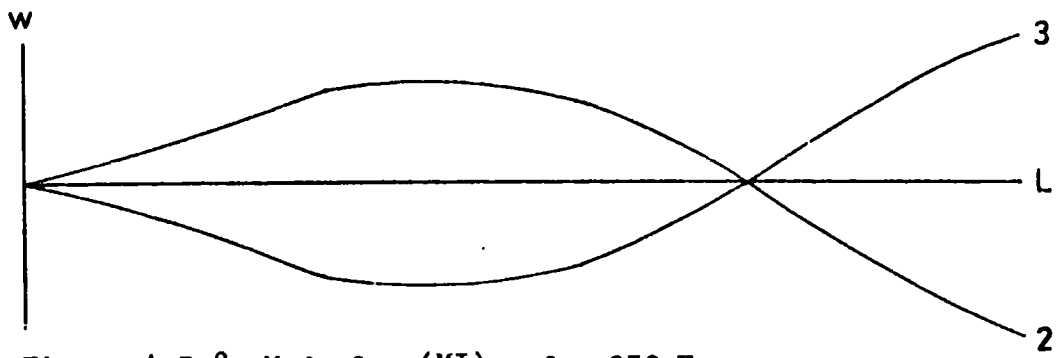


Figure A.3.8: Mode 9, (MI), $f = 950$ Hz.

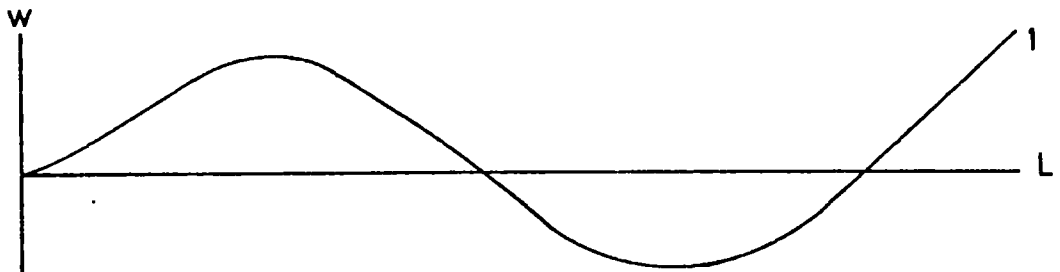


Figure A.3.9: Mode 10, (MI), $f = 1042$ Hz.

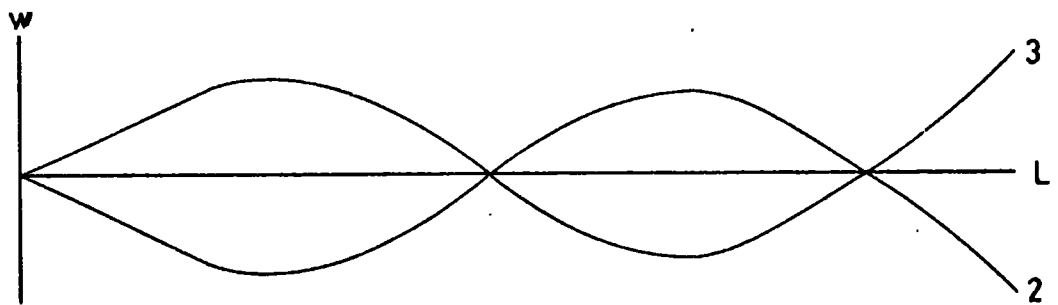


Figure A.3.10: Mode 11, (S), $f = 1047$ Hz.

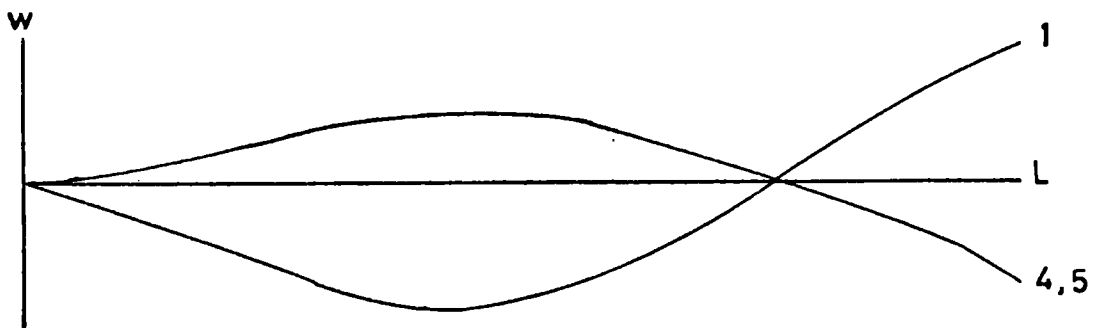


Figure A.3.11: Mode 12, (S), $f = 1062$ Hz.

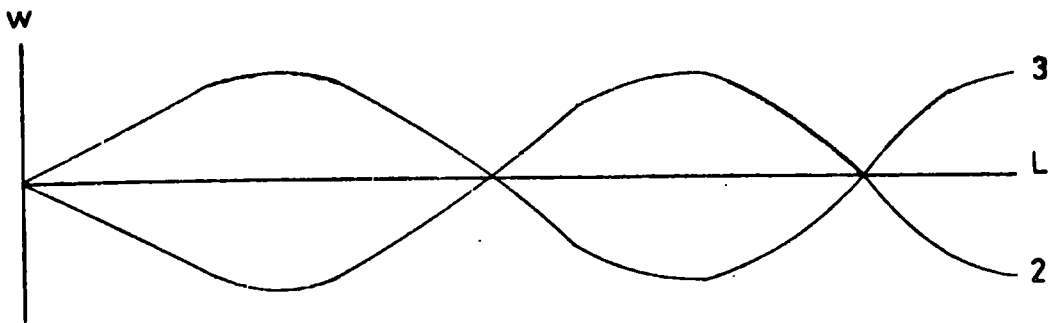


Figure A.3.12: Mode 13, (MI), $f = 1191$ Hz.

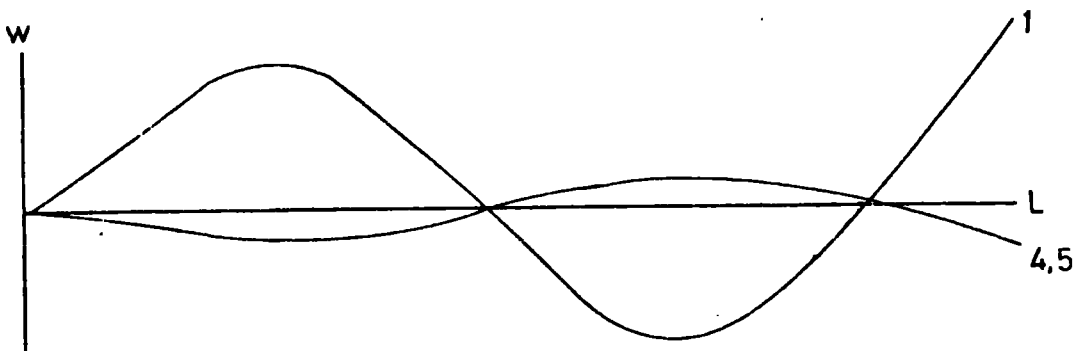


Figure A.3.13: Mode 14, (S), $f = 1359$ Hz.

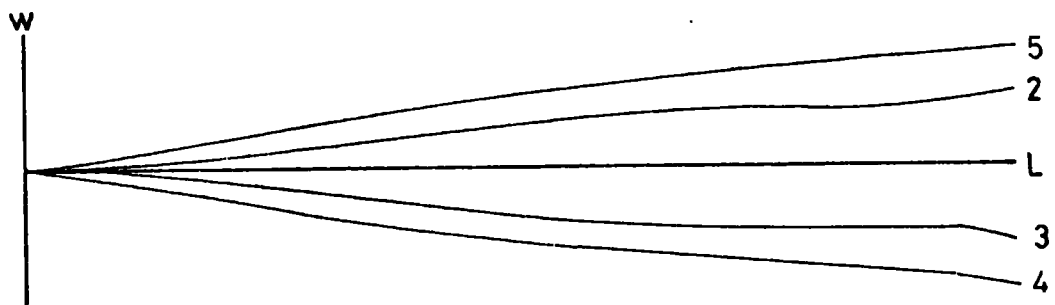


Figure A.3.14: Mode 15, (S), $f = 1419$ Hz.

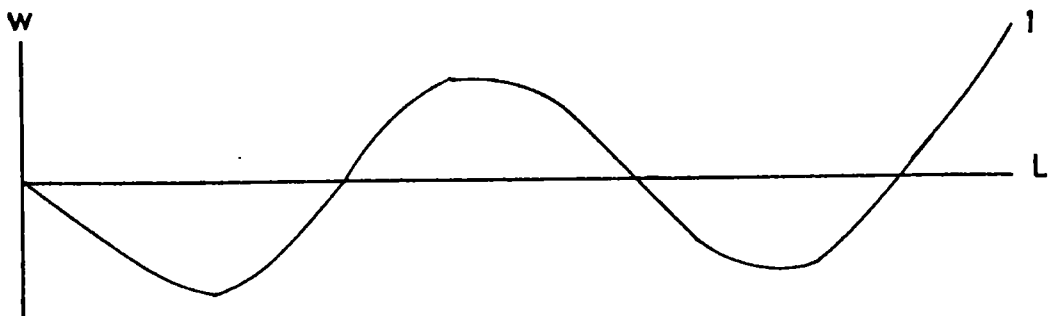


Figure A.3.15: Mode 16, (MI), $f = 1502$ Hz.

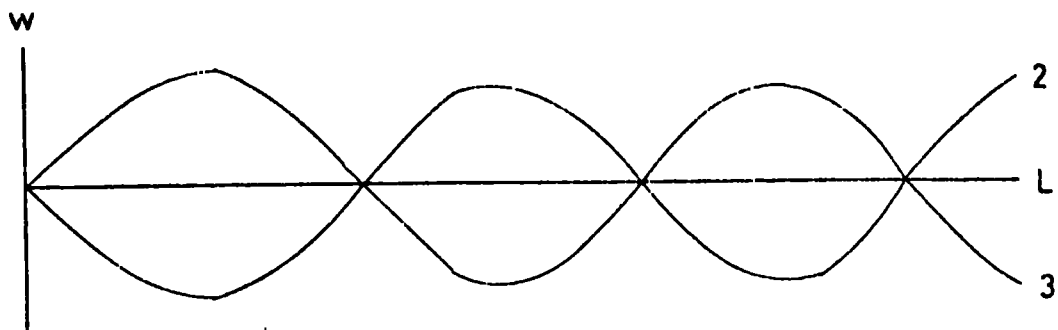


Figure A.3.16: Mode 17, (MI), $f = 1583$ Hz.

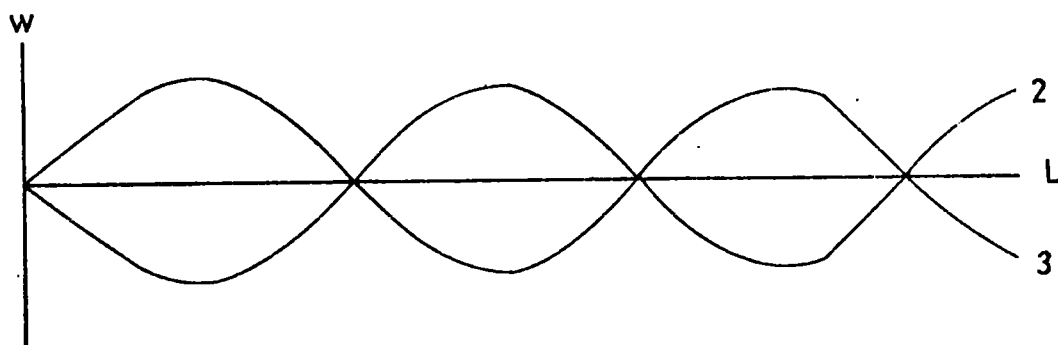


Figure A.3.17: Mode 18, (S), $f = 1598$ Hz.

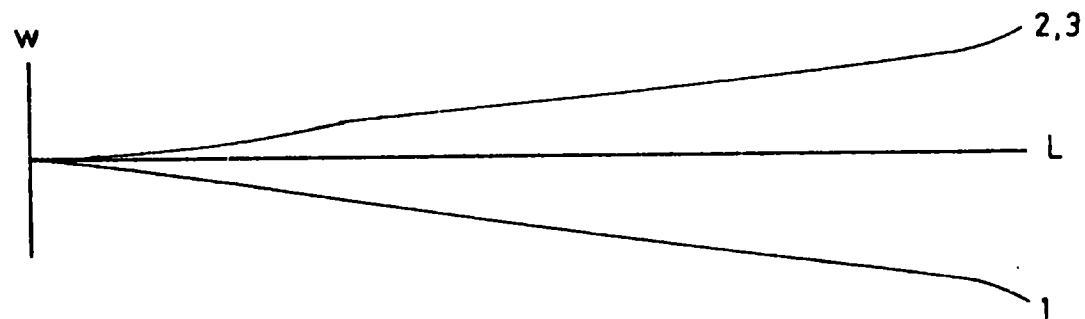


Figure A.3.18: Mode 19, (MI), $f = 1733$ Hz.

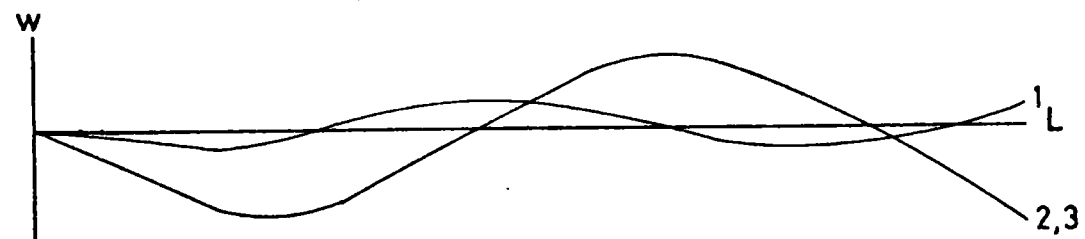


Figure A.3.19: Mode 20, (S), $f = 1513$ Hz.

Mod. No.	Top Face	Bottom Face	90 d.o.f.	146 d.o.f.	129 d.o.f.
1			158	158	158
2			307	306.8	306.8
3			372	371.5	371.6
4			597	596	588.4
5			661	660	648
6			839	836	837
7			759	758	738
8			-	-	97.9

Figure A. 3.20 (a): Mode shapes obtained from the finite element analysis of the real geometry.

Mod. No.	Top Face	Bottom Face	90 d.o.f.	146 d.o.f.	129 d.o.f.
9			1043	1036	1027
10			1016	1004	-
11			-	-	1052
12			-	1109	-
13			-	-	1155
14			-	1249	-
15			-	-	1299
16			-	-	1456

Figure A.3.20 (b):

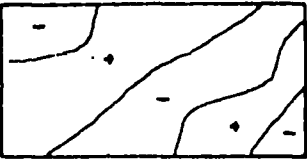
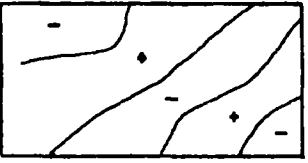

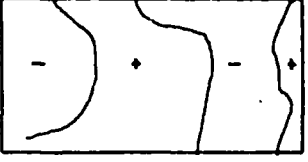
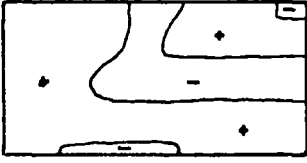

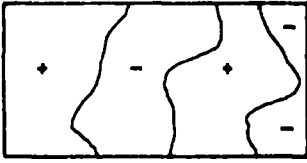
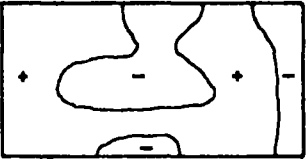
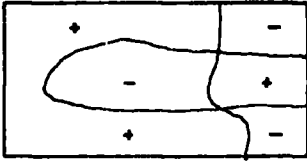
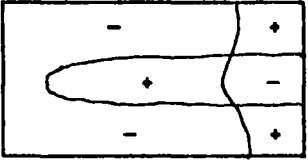
Mod. No.	Top Face	Bottom Face	129 d.o.f.
17	 Contour plot for Mod. No. 17 Top Face. The plot shows a diagonal contour line from the top-left to the bottom-right. The region above and to the left of the line is marked with a minus sign (-), and the region below and to the right is marked with a plus sign (+).	 Contour plot for Mod. No. 17 Bottom Face. The plot shows a diagonal contour line from the top-left to the bottom-right. The region above and to the left of the line is marked with a minus sign (-), and the region below and to the right is marked with a plus sign (+).	1541
18	 Contour plot for Mod. No. 18 Top Face. The plot shows a vertical contour line. The region to the left of the line is marked with a plus sign (+), and the region to the right is marked with a minus sign (-).	 Contour plot for Mod. No. 18 Bottom Face. The plot shows a vertical contour line. The region to the left of the line is marked with a minus sign (-), and the region to the right is marked with a plus sign (+).	1638
19	 Contour plot for Mod. No. 19 Top Face. The plot shows a horizontal contour line. The region above the line is marked with a plus sign (+), and the region below is marked with a minus sign (-).	 Contour plot for Mod. No. 19 Bottom Face. The plot shows a horizontal contour line. The region above the line is marked with a minus sign (-), and the region below is marked with a plus sign (+).	1818
20	 Contour plot for Mod. No. 20 Top Face. The plot shows a vertical contour line. The region to the left of the line is marked with a plus sign (+), and the region to the right is marked with a minus sign (-).	 Contour plot for Mod. No. 20 Bottom Face. The plot shows a vertical contour line. The region to the left of the line is marked with a plus sign (+), and the region to the right is marked with a minus sign (-).	1876
21	 Contour plot for Mod. No. 21 Top Face. The plot shows a vertical contour line. The region to the left of the line is marked with a plus sign (+), and the region to the right is marked with a minus sign (-).	 Contour plot for Mod. No. 21 Bottom Face. The plot shows a vertical contour line. The region to the left of the line is marked with a minus sign (-), and the region to the right is marked with a plus sign (+).	2005

Figure A.3.20 (c):

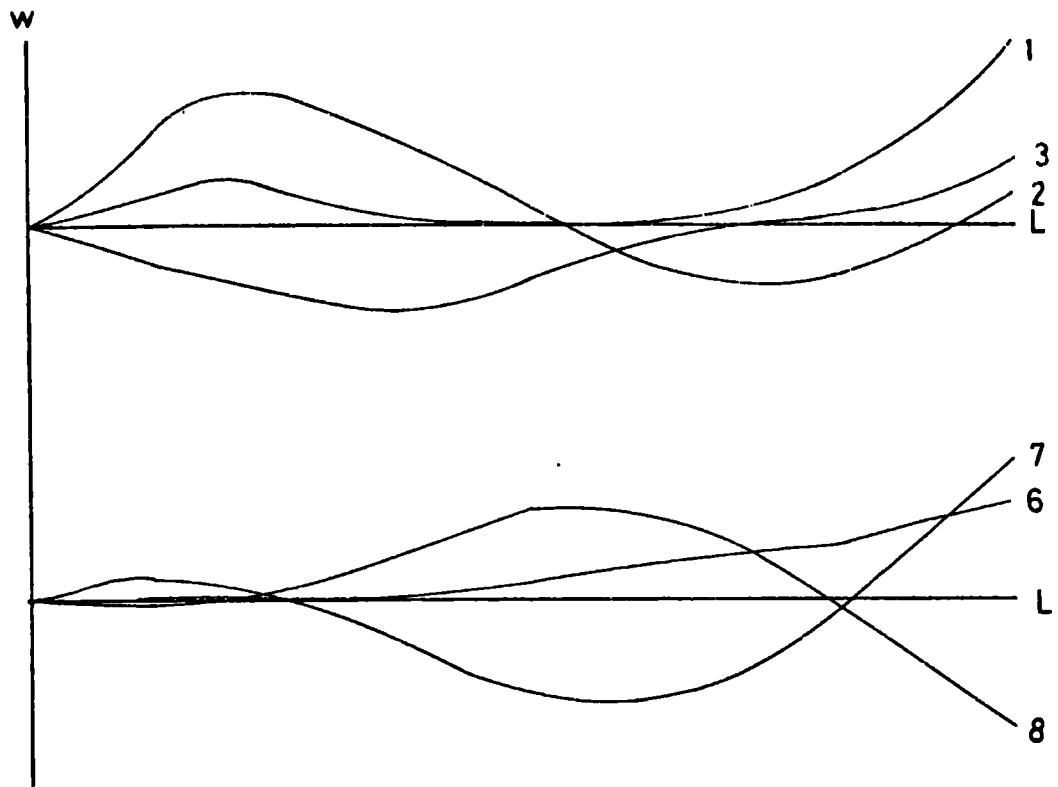


Figure A.3.21: Mode 8, $f = 979$ Hz.

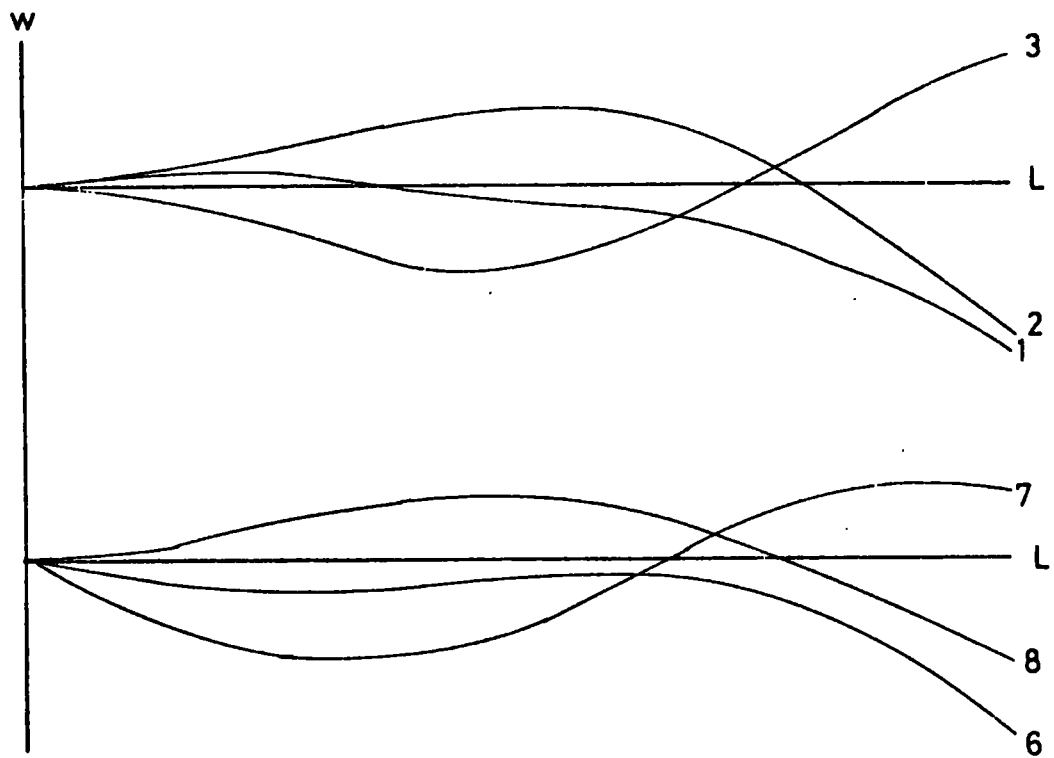


Figure A.3.22: Mode 9, $f = 1027$ Hz.

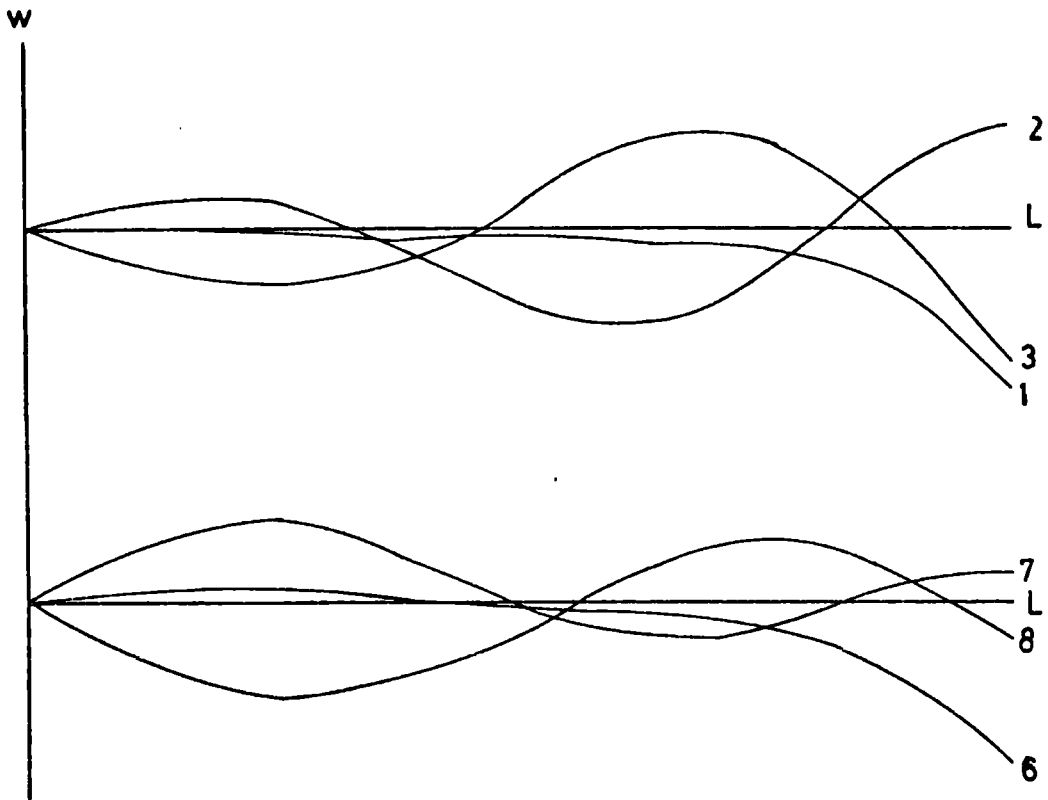


Figure A.3.23: Mode 11, $f = 1052$ Hz.

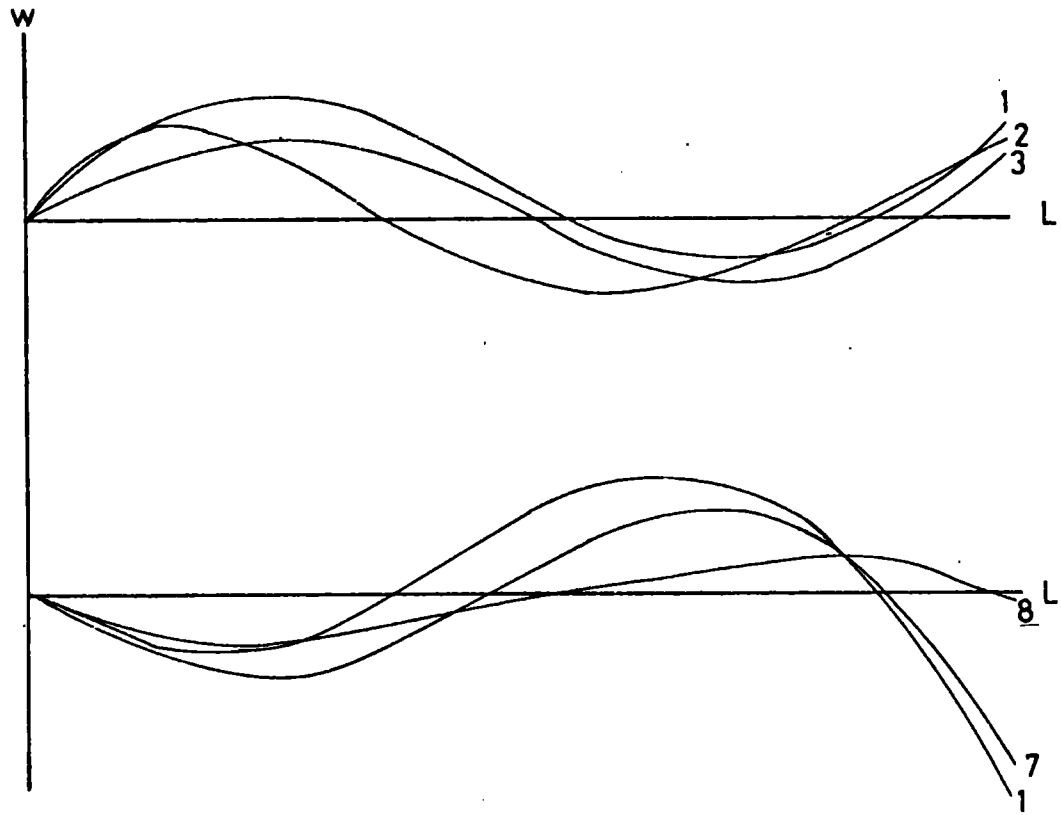


Figure A.3.24: Mode 13, $f = 1155$ Hz.

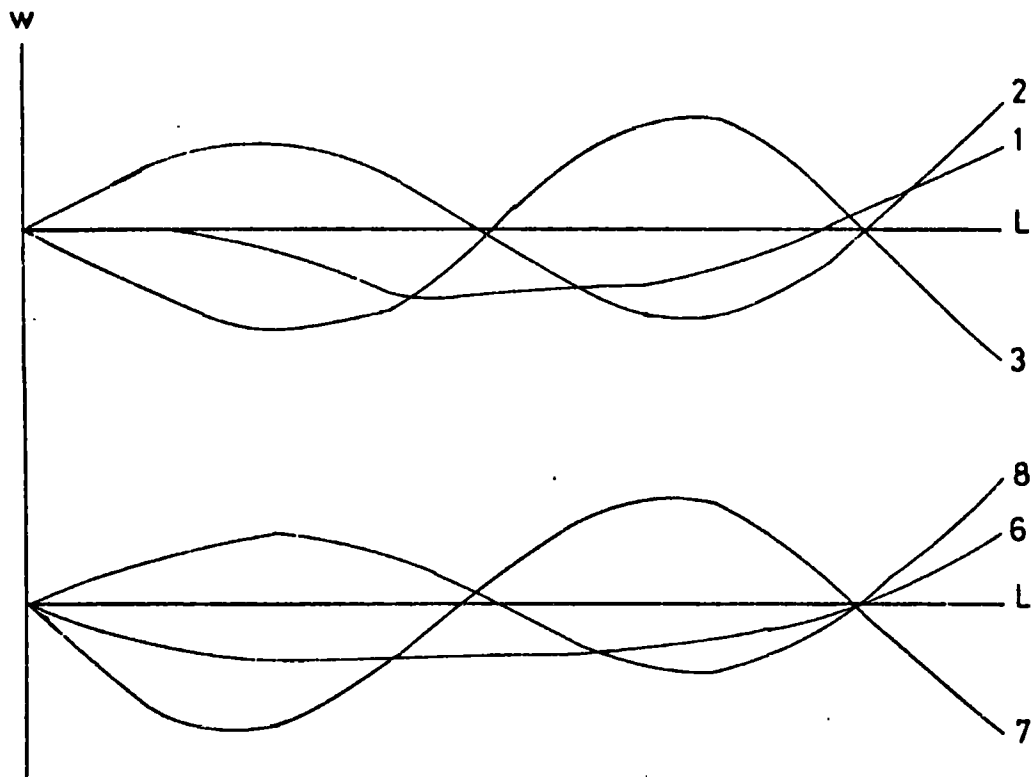


Figure A.3.25: Mode 15, $f = 1299$ Hz

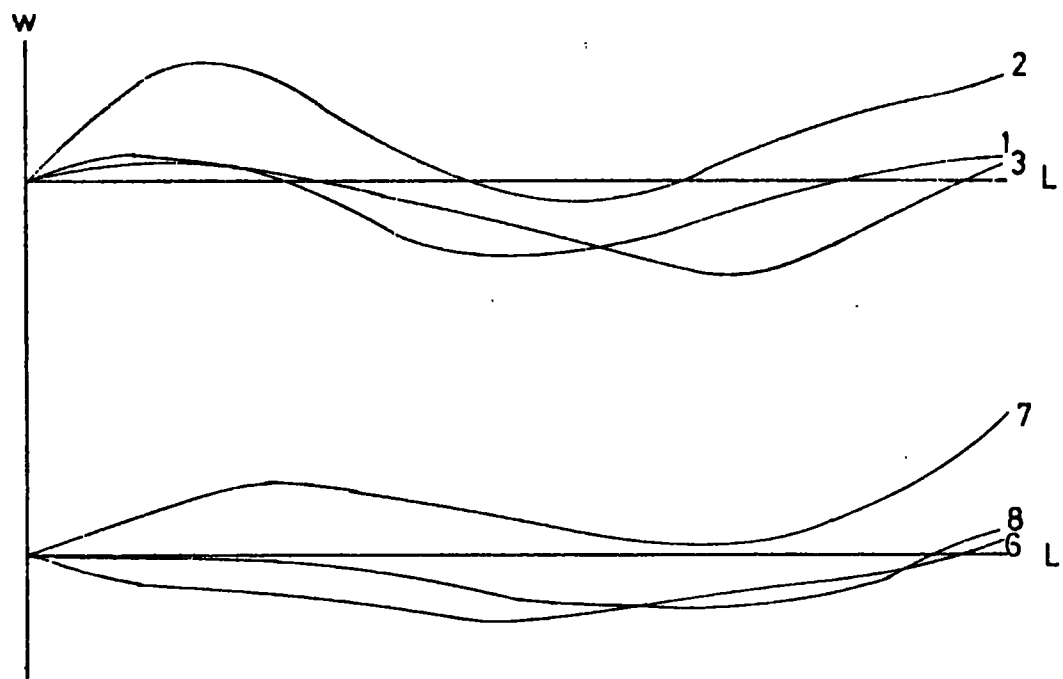


Figure A.3.26: Mode 16, $f = 1456$ Hz.

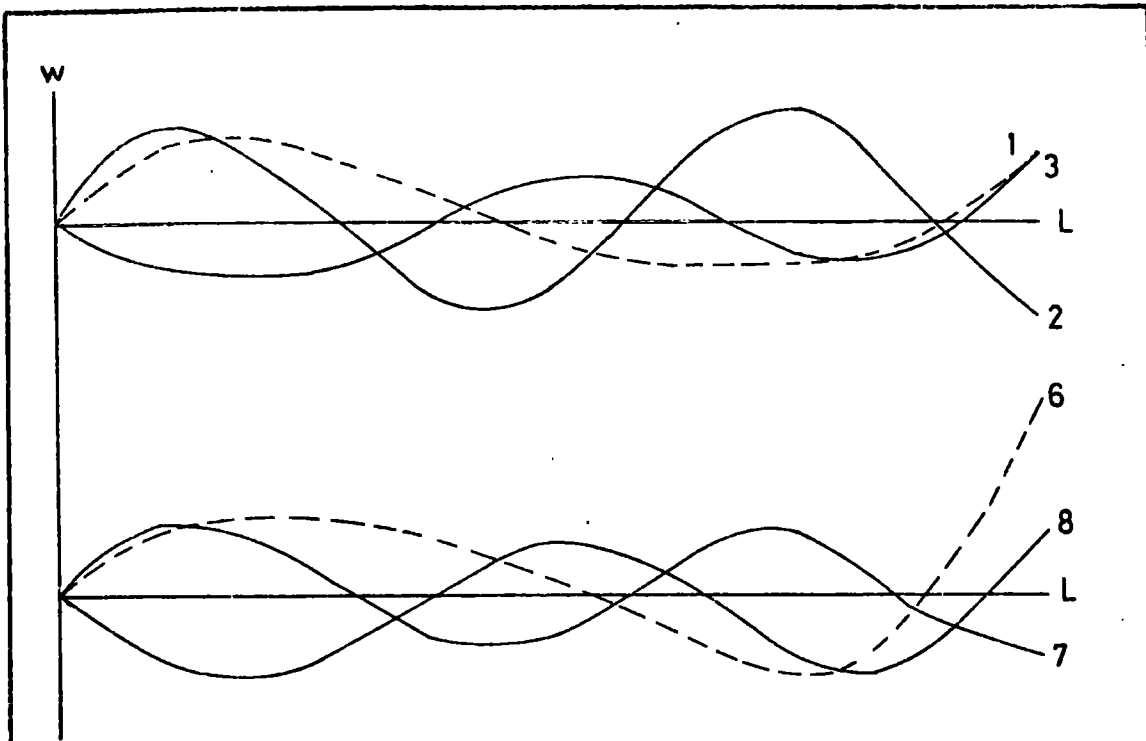


Figure: A.3.27: Mode 17, $f = 1541$ Hz.

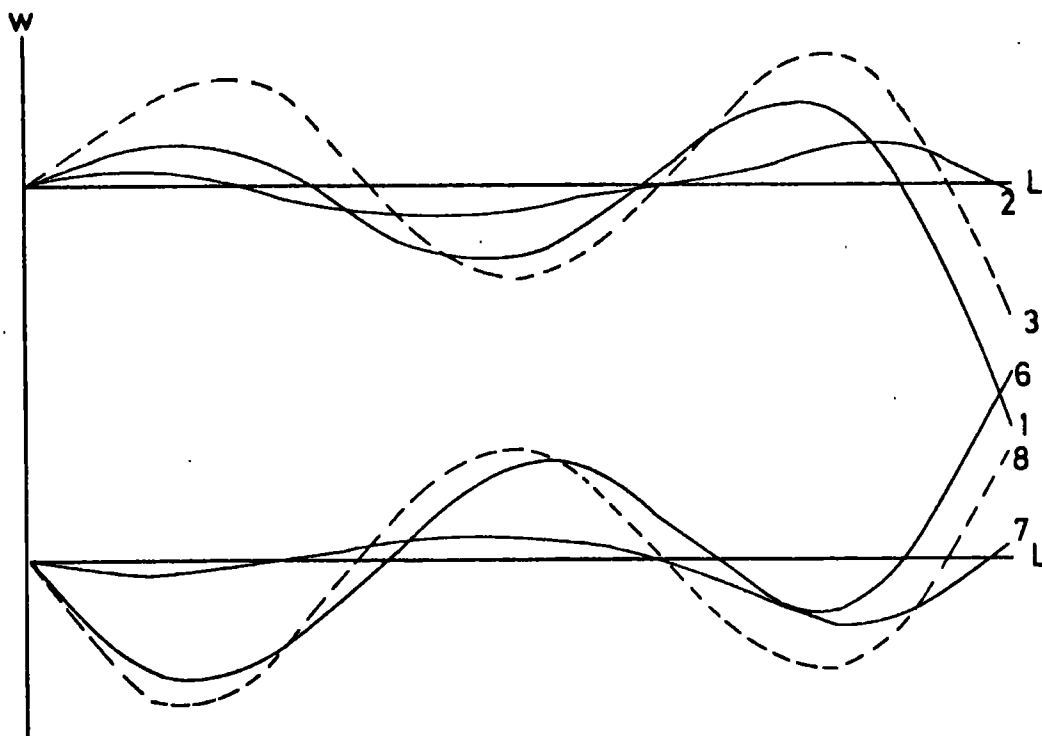


Figure A.3.28: Mode 18, $f = 1638$ Hz.

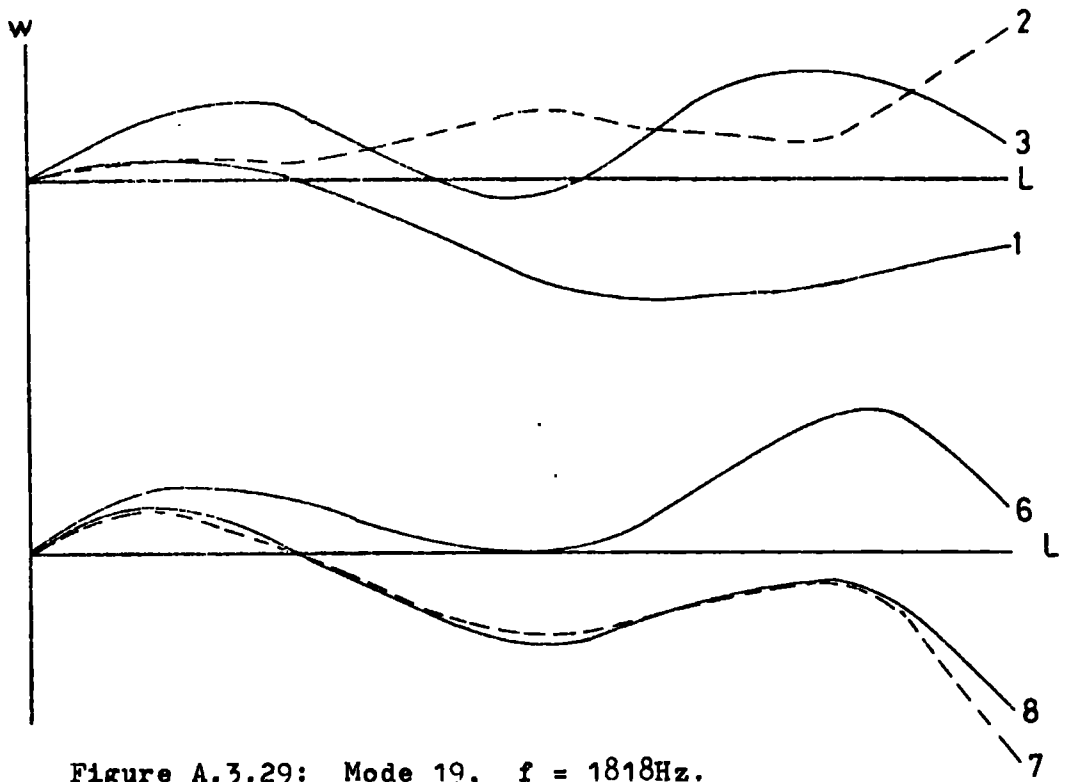


Figure A.3.29: Mode 19, $f = 1818\text{Hz}$.

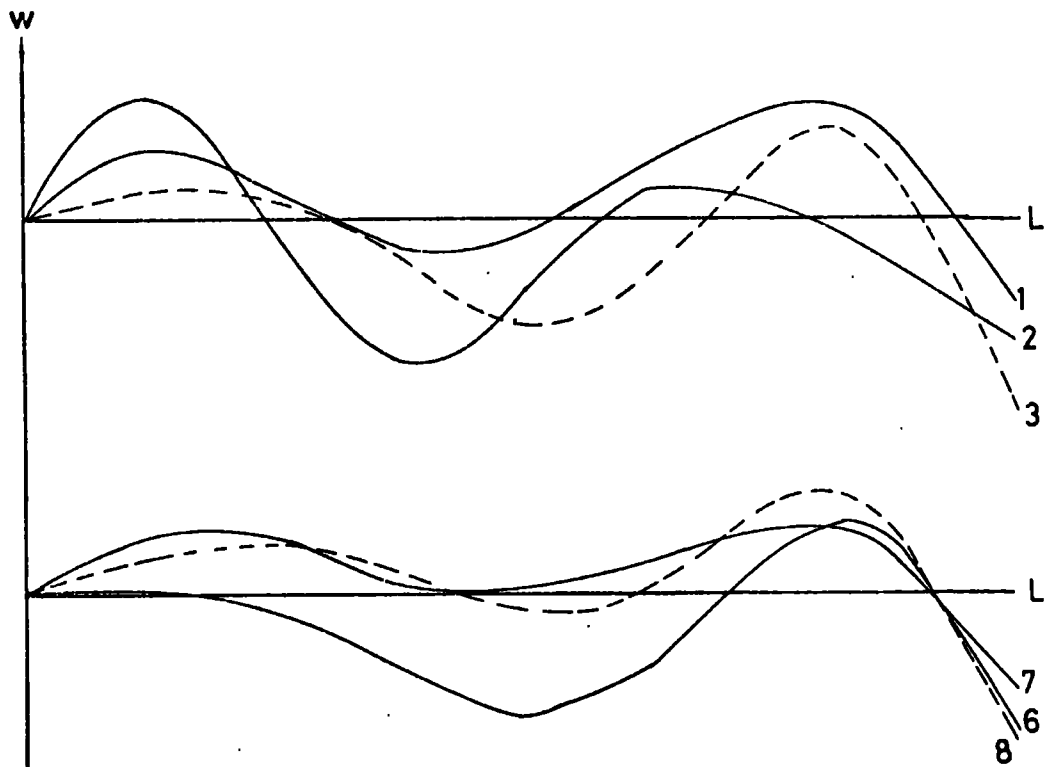


Figure A.3.30: Mode 20. $f = 1876\text{ Hz}$.

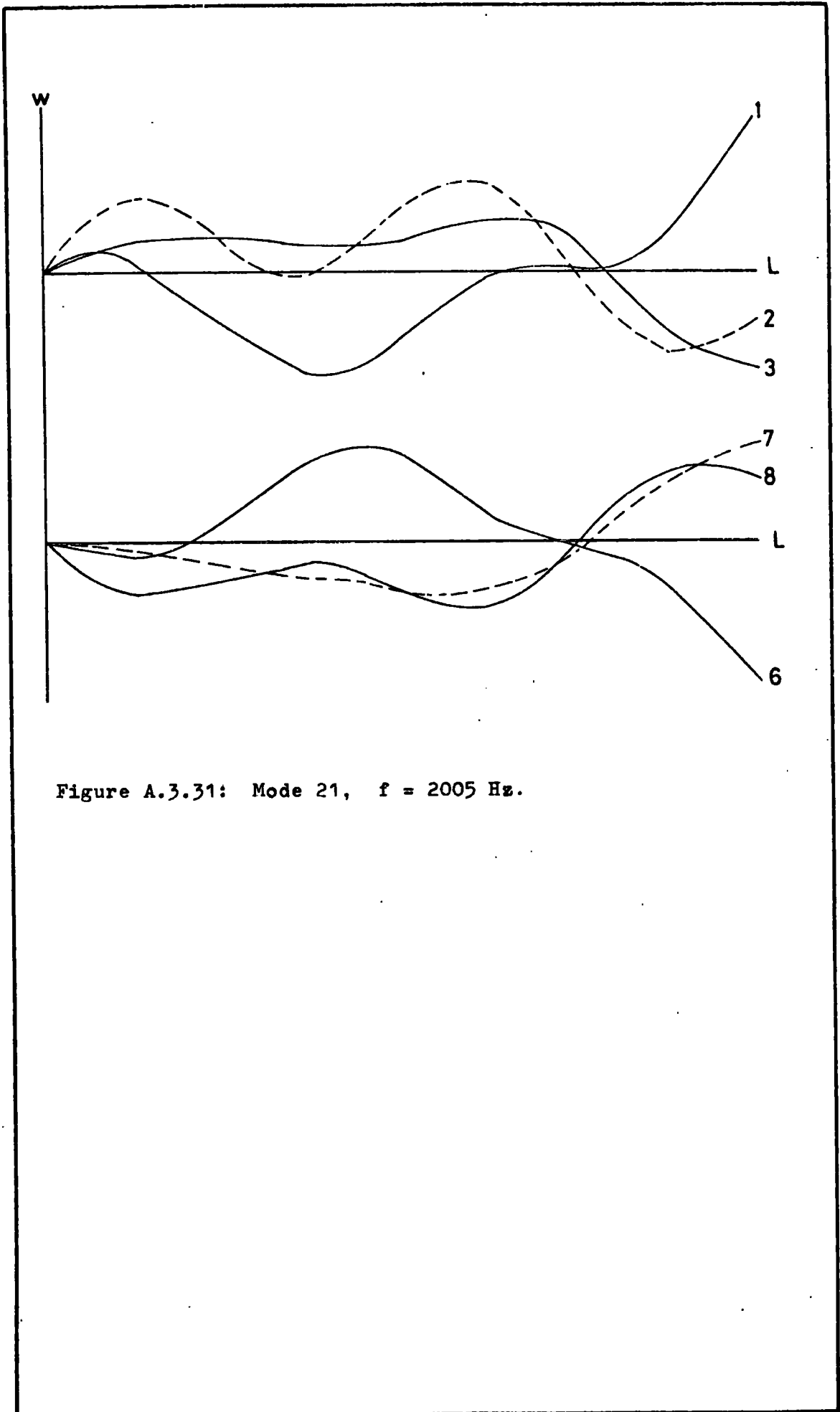


Figure A.3.31: Mode 21, $f = 2005$ Hz.

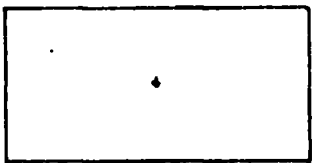
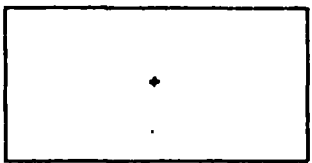
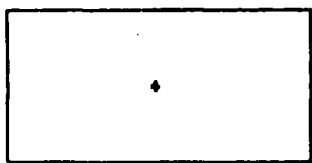
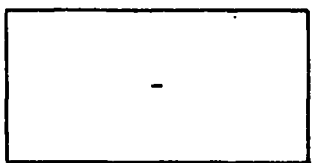
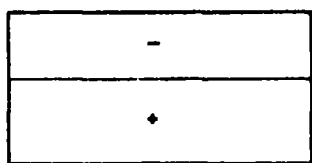
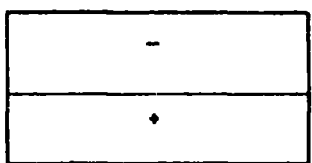
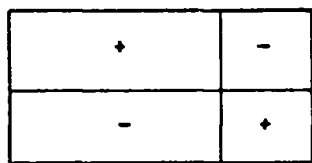
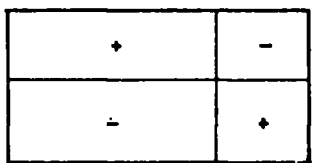
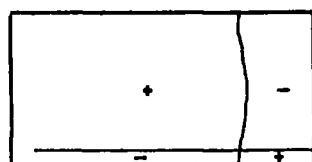
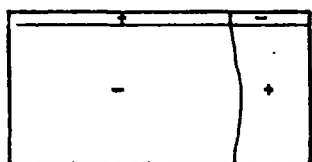
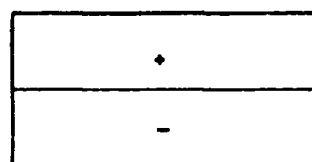
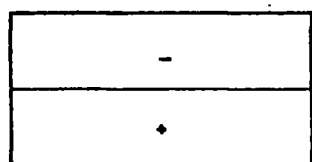
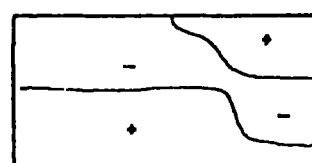
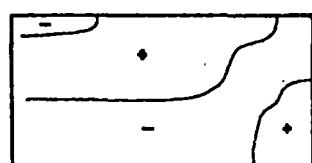


Mod. No.	Top Face	Bottom Face	Freq. Hz.
1			142
2			297
3			355
4			561.5
5			598
6			812
7			905
8			935

Figure A.3.32(a): Experimental mode shapes and natural frequencies.

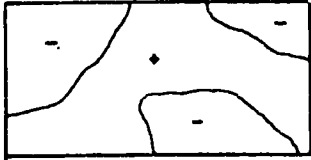
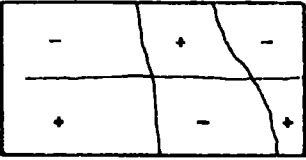
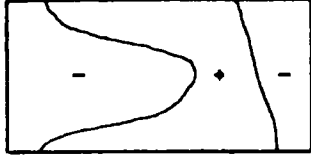


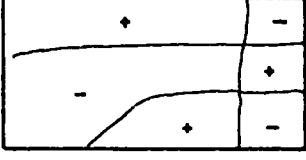
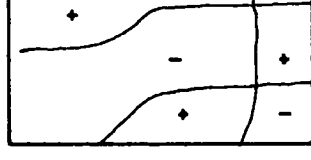


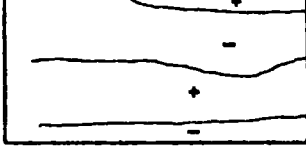


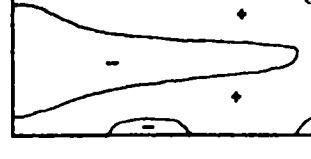
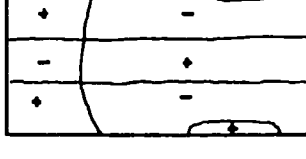
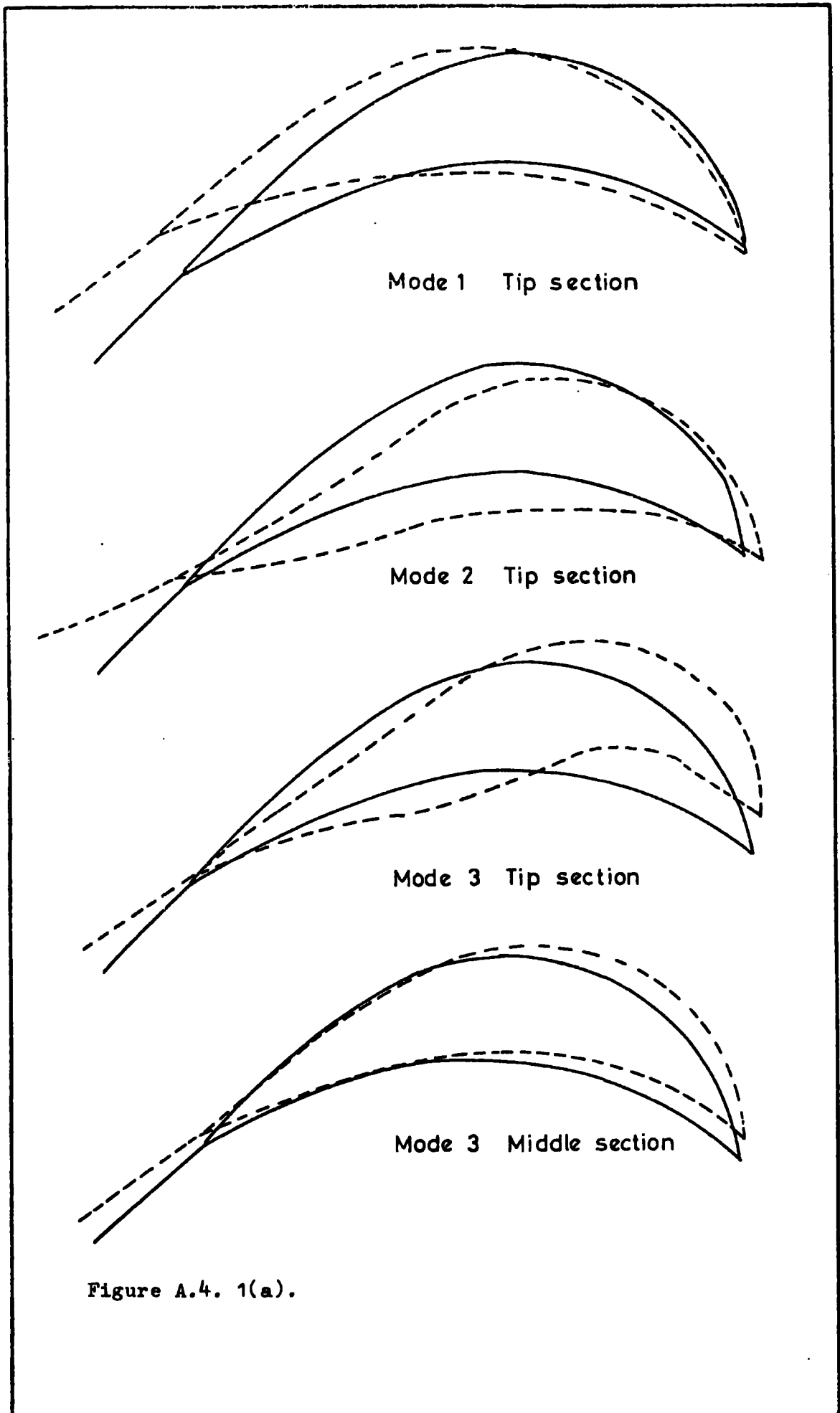
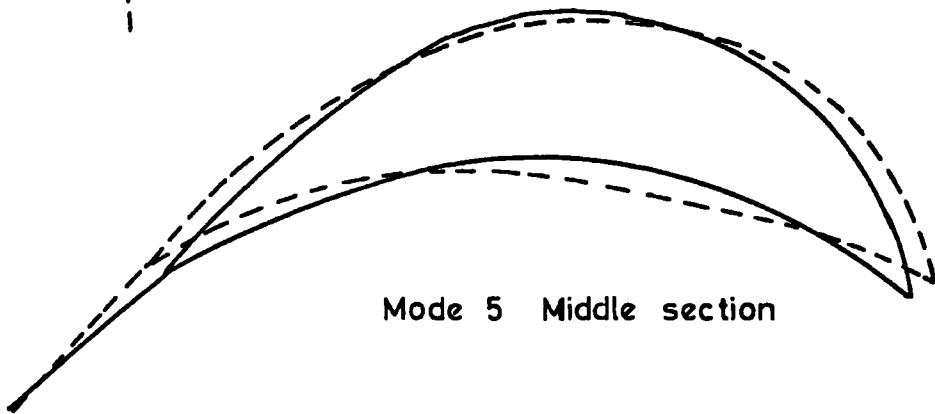
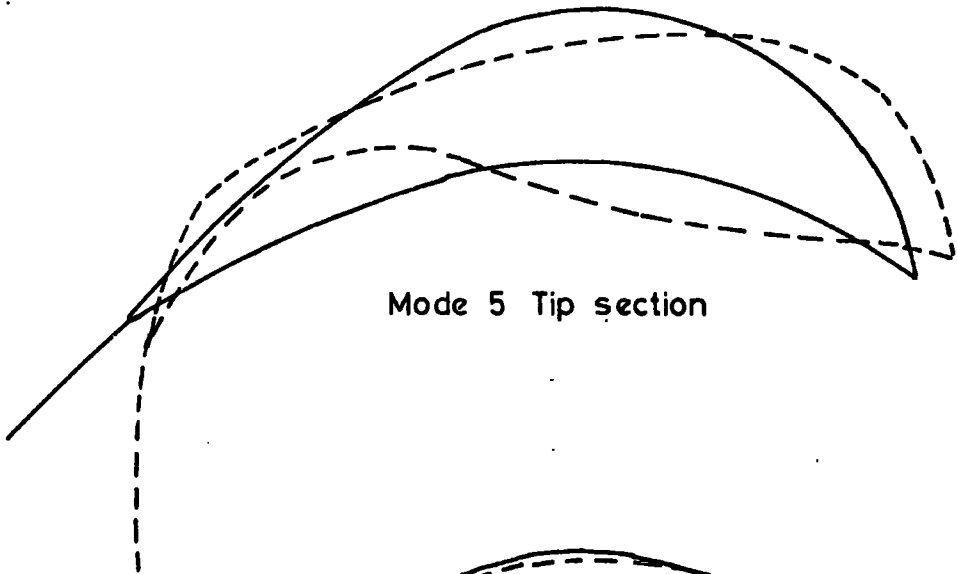
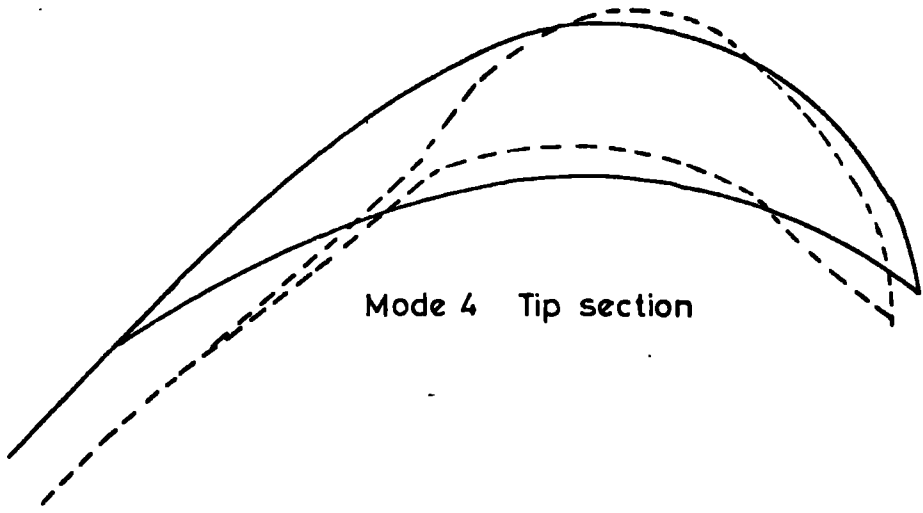
Mod. No.	Top Face	Bottom Face	Freq. Hz.
9			980
10			1096
11			1136
12			1157
13			1314
14			1387
15			1633

Figure A.3.32 (b): Experimental mode shapes and natural frequencies.

APPENDIX 4

Deflection of the cross-section of the G.E.C.
turbine blade analysed in Section 6.4.





A.4.1. (b):

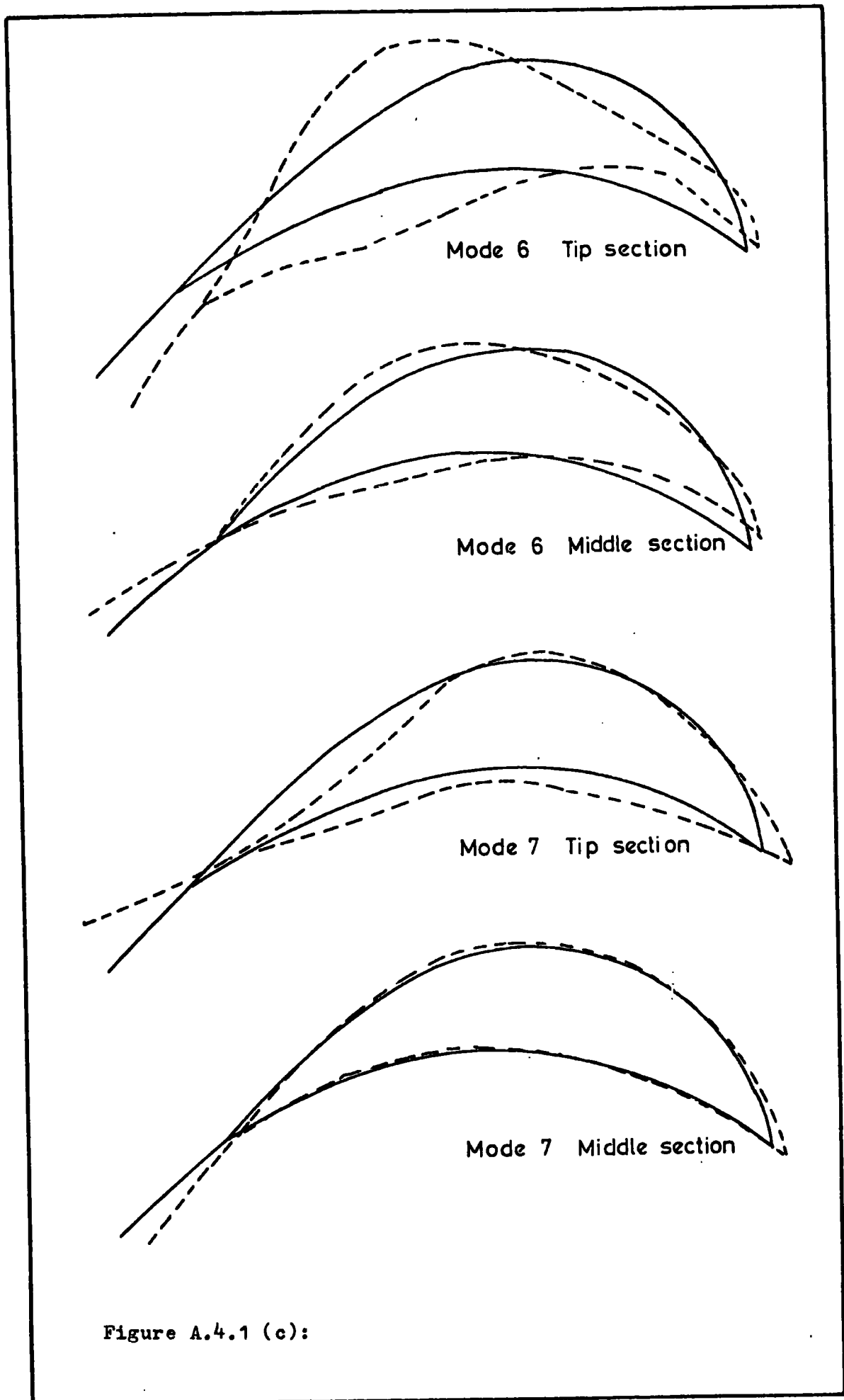


Figure A.4.1 (c):

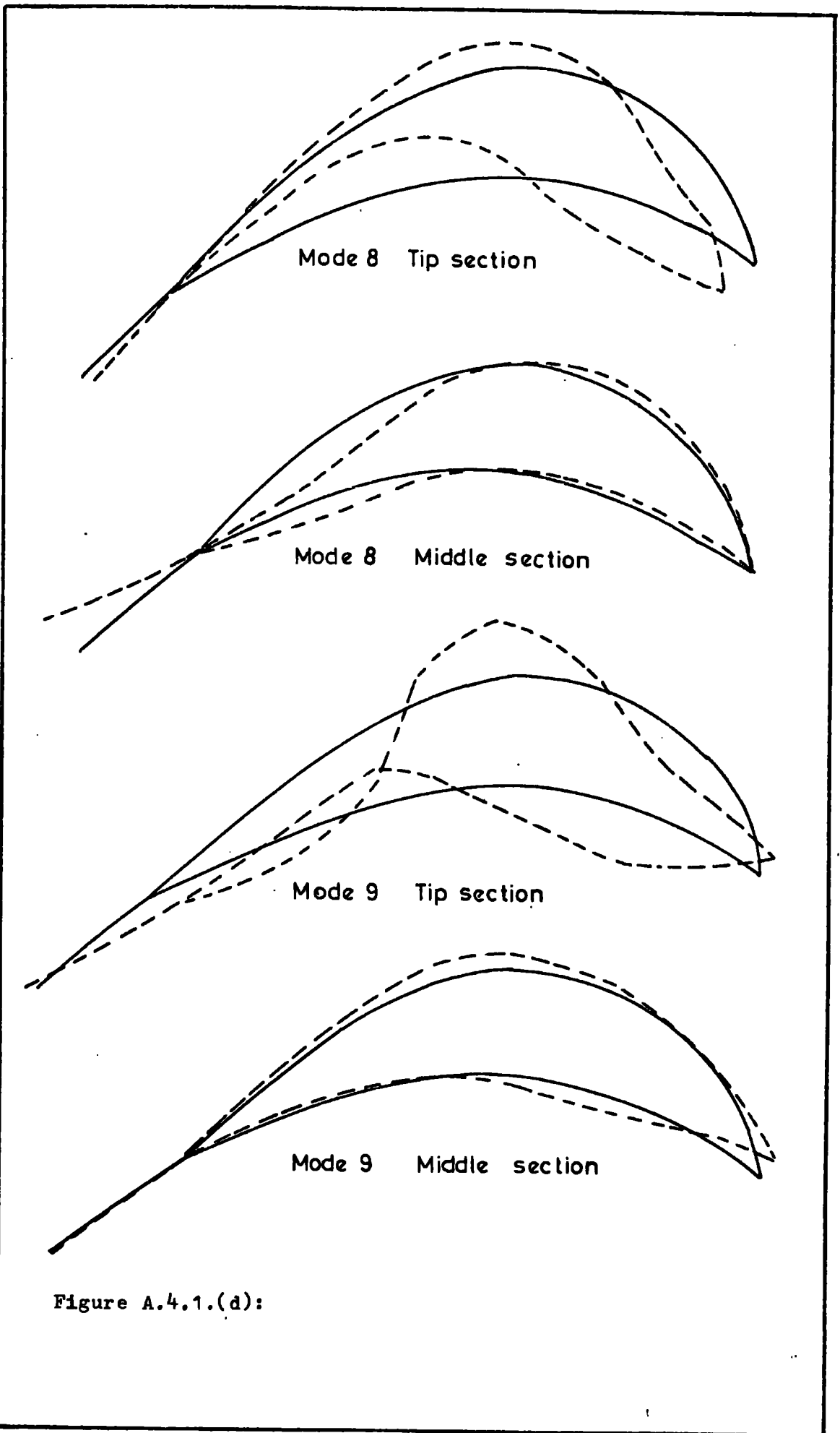


Figure A.4.1.(d):

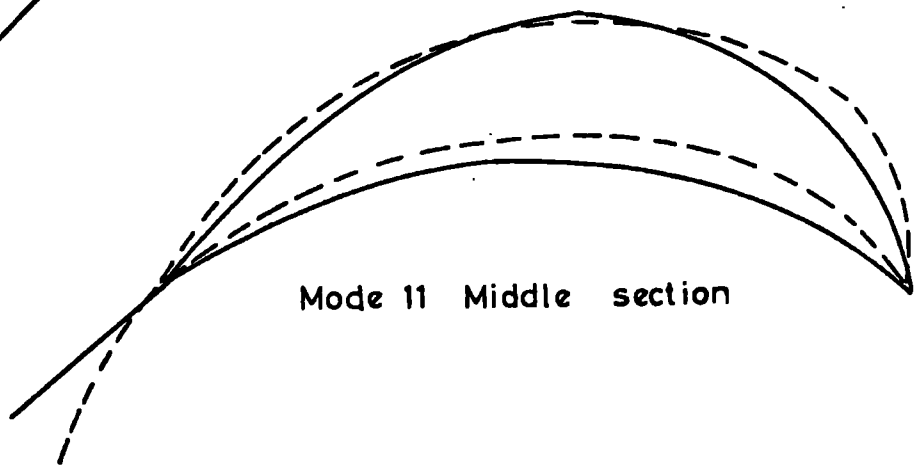
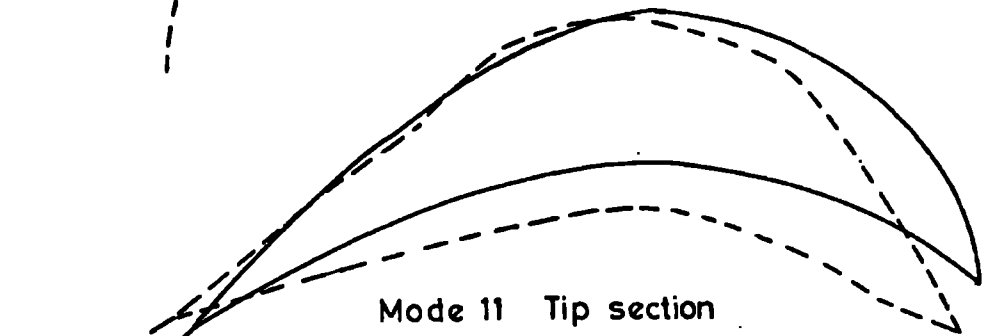
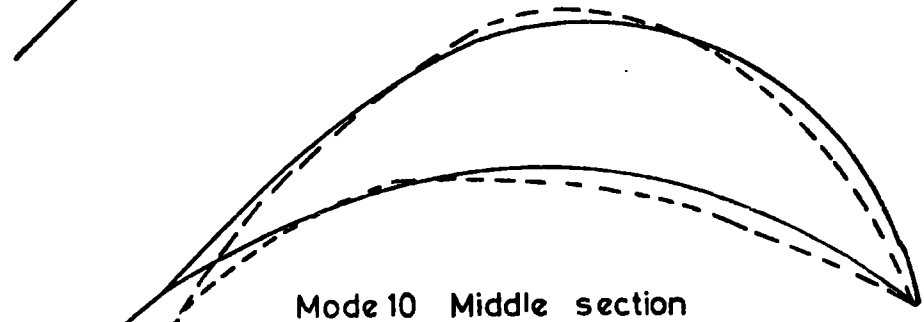
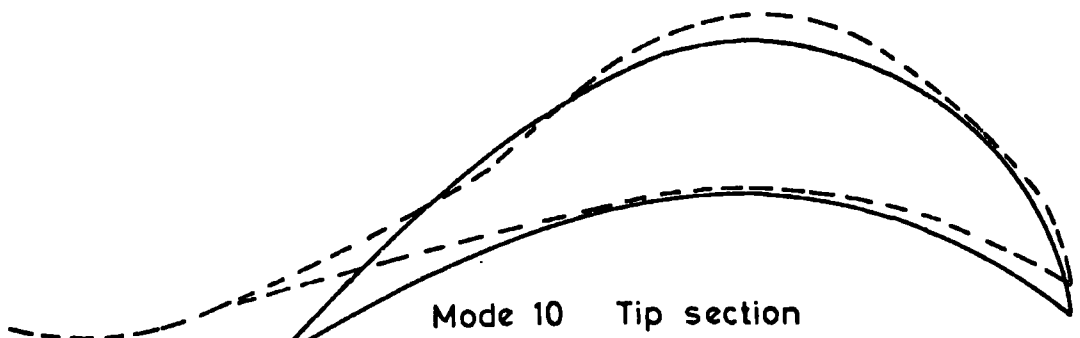


Figure A.4.1.(e):

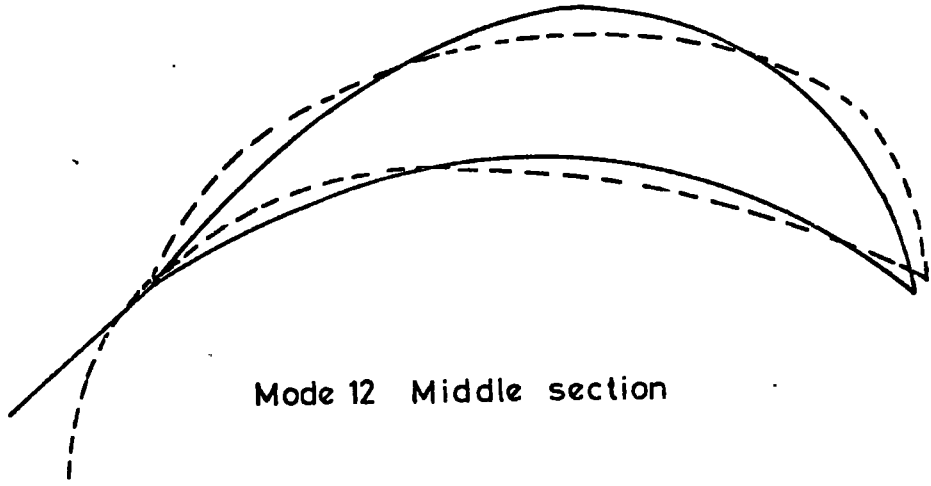
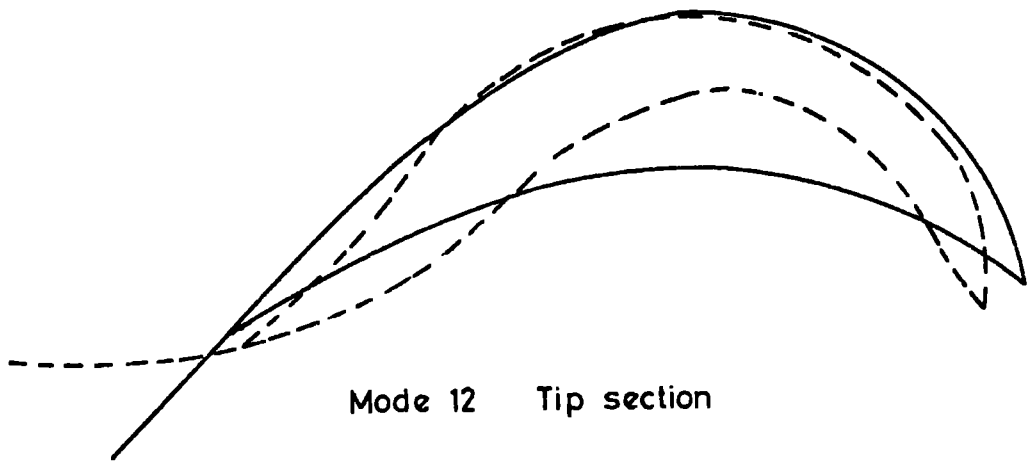


Figure A.4.1. (f):

APPENDIX 5

Listing of the computer program used in this thesis.

```

1.000 C PLATE AND SHELL VIBRATICN WITH FINITE ELEMENTS MU 0001
2.000 C PSVFE MU 0002
3.000 C MU 0003
4.000 C MU 0004
5.000 C MU 0005
6.000 C MU 0006
7.000 C NATUREL FREQUENCY ANALYSIS CF ARBITRARY SHELL STRUCTURES MU 0007
8.000 C USING AHMAD'S THICK-SHELL ELEMENT MU 0008
9.000 C MU 0009
10.000 C MU 0010
11.000 C MEHMET UCMAKLICGLU MU 0011
12.000 C MU 0012
13.000 C DEPARTMENT OF ENGINEERING SCIENCE MU 0013
14.000 C UNIVERSITY OF DURHAM MU 0014
15.000 C MU 0015
16.000 C 1975-1978 MU 0016
17.000 C MU 0017
18.000 C MU 0018
19.000 C MU 0019
20.000 C MU 0020
21.000 C MU 0021
22.000 C MU 0022
23.000 C MU 0023
24.000 C MU 0024
25.000 C COMMON/COR/XT(10),XB(10),YI(10),YB(10),ZT(10),ZB(10) MU 0025
26.000 C COMMON/JAK/RJ3(3,3),RJ2(3,3),AJNT(3) MU 0026
27.000 C COMMON/INTG/PNT4(4),PNT3(3),PNT2(2),H4(4),H3(3) MU 0027
28.000 C COMMON/STHE/V1(3),V2(3),V3(3) MU 0028
29.000 C COMMON/UNI/UNIT(3) MU 0029
30.000 C COMMON/GEN/WS(50),IPT(50),IBC1,IBC2,IS1,IS2,MIS1,MIS2 MU 0030
31.000 C COMMON/EPRO/E(5),RC(5),PR(5) MU 0031
32.000 C COMMON/RED/SLAVES(300),SLAVERM(300),NGDE(10,2),NDIS(10,5), MU 0032
33.000 C &KAYVEC(500),ICCM(500) MU 0033
34.000 C MU 0034
35.000 C DIMENSION TITLE(20),RJI(3,3),RJ(3,3), MU 0035
36.000 C &DUDG(3,50,3),DUDC(3,50,3),B(5,50),BT(50,5),S(5C,50),Z(5),D(5,5), MU 0036
37.000 C &UVW(3,50),UVWT(50,3),EFEM(50,50),SS(50,50),ELMAS(50,50) MU 0037
38.000 C &,DCM(3,3),DCMT(3,3),PMM(45150),PSM(45150), MU 0038
39.000 C &RCOT(50),VEC(7500),V1(3),V2(3),V3(3),V3I(3),INVER(4) MU 0039
40.000 C MU 0040

```


LINE NUMBER

TEXT

```

81.000 C READ THE TOPOLOGY AND THE COORDINATES
82.000 C PERFORM FIRST CALCULATIONS
83.000 C PRINT OUT THE RESULTS
84.000 C COPY THEM ON DEVICE NO:7 AS FURTHER DATA
85.000 C
86.000 C 701 DO 311 I=1,NGNO
87.000 C ICCM(I)=0
88.000 C 311 CONTINUE
89.000 C IF(IPRINT.LT.2) GO TO 332
90.000 C WRITE(6,315)
91.000 C 315 FORMAT(1H1,///20X,'**TOPOLOGY OF THE ELEMENTS**',
92.000 C &,///20X,'ELEM. NO:',22X,'TOPOLOGY',40 X,'INT.PNTS.'/)
93.000 C
94.000 C READ AND WRITE TOPOLOGY
95.000 C
96.000 C 332 IF(ISTEK.LT.3) GO TO 702
97.000 C CALL GENTOP(NODE,ICCM,NO,NEY,ISTEK)
98.000 C DO 350 I=1,NO
99.000 C LINE=I*1000
100.000 C READ(7,313) (NCDE(J,1),J=1,10),INTEG,LTIP,ITYP,IPTRO,INVER
101.000 C 350 WRITE(6,314) I,(NCDE(J,1),J=1,10),INTEG,LTIP,ITYP,IPTRO,INVER
102.000 C GO TO 334
103.000 C
104.000 C 702 DO 312 I=1,NO
105.000 C READ(5,313) (NCDE(J,1),J=1,10),INTEG,LTIP,ITYP,IPTRO,INVER
106.000 C WRITE(7,313) (NCDE(J,1),J=1,10),INTEG,LTIP,ITYP,IPTRO,INVER
107.000 C IF(IPRINT.LT.2) GO TO 703
108.000 C WRITE(6,314) I,(NCDE(J,1),J=1,10),INTEG,LTIP,ITYP,IPTRO,INVER
109.000 C 703 DO 312 K=1,LTIP
110.000 C LOC=NODE(K,1)
111.000 C ICCM(LOC)=ICCM(LOC)+11
112.000 C 312 CONTINUE
113.000 C 313 FORMAT(14I5,4I2)
114.000 C 314 FFORMAT(20X,I4,10X,10(2X,I4),5X,I3,5X,7I2)
115.000 C 334 WRITE(6,327)
116.000 C
117.000 C READ AND WRITE THE COORDINATES
118.000 C
119.000 C IF(ISTEK.LT.3) GO TO 333
120.000 C CALL GENOD(KLEN,NODE,NC,NEY)

```

MU 0081
 MU 0082
 MU 0083
 MU 0084
 MU 0085
 MU 0086
 MU 0087
 MU 0088
 MU 0089
 MU 0090
 MU 0091
 MU 0092
 MU 0093
 MU 0094
 MU 0095
 MU 0096
 MU 0097
 MU 0098
 MU 0099
 MU 0100
 MU 0101
 MU 0102
 MU 0103
 MU 0104
 MU 0105
 MU 0106
 MU 0107
 MU 0108
 MU 0109
 MU 0110
 MU 0111
 MU 0112
 MU 0113
 MU 0114
 MU 0115
 MU 0116
 MU 0117
 MU 0118
 MU 0119
 MU 0120


```

121.000 DO 351 I=1,NGND
122.000 LINE=(NO+I)*1000
123.000 READ(7,LINE,325)NDGF, IDIS,XC,XXC,YC,YYC,ZC,ZZC,ICOM(I),NDCF, IDIS
124.000 351 WRITE(6,326)I,XC,XXC,YC,YYC,ZC,ZZC,ICOM(I),NDCF, IDIS
125.000 GC TO 335
126.000 C
127.000 333 DO 321 I=1,NGND
128.000 READ(5,324)XC,XXC,YC,YYC,ZC,ZZC,NDGF, IDIS
129.000 WRITE(7,325)NDCF, IDIS,XC,XXC,YC,YYC,ZC,ZZC
130.000 IF(IPRINT.LT.2) GC TO 321
131.000 WRITE(6,326) I,XC,XXC,YC,YYC,ZC,ZZC,ICOM(I),NDCF, IDIS
132.000 321 CONTINUE
133.000 324 FORMAT(6E10.4,2I5)
134.000 325 FORMAT(2I5,6E17.9)
135.000 326 FORMAT(6X,I4,4X,6(2X,F10.4),5X,I2,6X,I3,5X,I5)
136.000 327 FORMAT(1H1,///20X,***COORDINATES CF THE NCDES***,
137.000 &///4X,'NODE NO:',10X,'XT',10X,'XB',1CX,'YT',1CX,'YB',10X,'ZT',
138.000 &,10X,'ZB',7X,'NGCE',5X,'B.C.',4X,'DOF',/)
139.000 335 CONTINUE
140.000 MAX=0
141.000 MBD=NGND*5
142.000 C
143.000 C INITIAL VALUES CF KAYVEC
144.000 KAYMA=0
145.000 DU 400 I=1,NGND
146.000 LINE=(NO+I)*1000
147.000 READ(7,LINE,401) NDOF
148.000 KAYMA=KAYMA+(5-NDOF)
149.000 KAYVEC(I)=KAYMA
150.000 400 CONTINUE
151.000 401 FORMAT(I5)
152.000 C
153.000 IF(MAPRO.EQ.1) CALL CONEL(D,I)
154.000 C
155.000 WRITE(6,805)
156.000 805 FORMAT(1H1,10X,'ELEMENT INFORMATION:')
157.000 C
158.000 DO 441 I=1,45150
159.000 PMM(I)=0.0
160.000 PSM(I)=0.0
MU 0121
MU 0122
MU 0123
MU 0124
MU 0125
MU 0126
MU 0127
MU 0128
MU 0129
MU 0130
MU 0131
MU 0132
MU 0133
MU 0134
MU 0135
MU 0136
MU 0137
MU 0138
MU 0139
MU 0140
MU 0141
MU 0142
MU 0143
MU 0144
MU 0145
MU 0146
MU 0147
MU 0148
MU 0149
MU 0150
MU 0151
MU 0152
MU 0153
MU 0154
MU 0155
MU 0156
MU 0157
MU 0158
MU 0159
MU 0160

```

```

161.000 441 CONTINUE
162.000 CG 222 NOEL=1,NG
163.000 WRITE(6,806) NCEL
164.000 806 FORMAT(/,5X,'ELEM. NO=',I2/)
165.000 LINEN=NOEL*1000
166.000 READ(7,LINEN,89) (NODE(I,1),I=1,10),INT1,INT2,INT3,LTIP,ITYP,IPRO
167.000 &,INVER
168.000 89 FORMAT(10I5,2X,3I1,3I5,4I2)
169.000 DO 62 IM=1,LTIP
170.000 LINE=(NODE(IM,1)+NG)*1000
171.000 DO 50 ININ=1,4
172.000 IF(IM.EQ.INVER(ININ)) GO TO 51
173.000 50 CONTINUE
174.000 READ(7,LINE,88) NODE(IM,2),(NDIS(IM,J),J=1,5)
175.000 &, XT(IM),XB(IM),YT(IM),YB(IM),ZT(IM),ZB(IM)
176.000 GO TO 52
177.000 51 READ(7,LINE,88) NODE(IM,2),(NDIS(IM,J),J=1,5)
178.000 &, XB(IM),XT(IM),YB(IM),YT(IM),ZB(IM),ZT(IM)
179.000 88 FORMAT(I5,5I1,6E17.9)
180.000 52 INDEX=0
181.000 DO 63 MI=1,5
182.000 63 IF(NDIS(IM,MI).EQ.3) INDEX=1
183.000 LCC=NCDE(IM,1)
184.000 IF(INDEX.NE.0) GO TO 64
185.000 ICOM(LOC)=ICOM(LOC)-1
186.000 GO TO 62
187.000 64 IC=ICCM(LOC)
188.000 ICOM(LOC)=-(IABS(IC))+1
189.000 62 CONTINUE
190.000 LT5=5*LTIP
191.000 IF(IPRINT.LT.7) GO TO 777
192.000 WRITE(6,807)(XT(I),I=1,LTIP)
193.000 807 FORMAT(2X,'XT:',10(2X,F10.5))
194.000 WRITE(6,808)(XB(I),I=1,LTIP)
195.000 808 FORMAT(2X,'XB:',10(2X,F10.5))
196.000 WRITE(6,809)(YT(I),I=1,LTIP)
197.000 809 FORMAT(2X,'YT:',10(2X,F10.5))
198.000 WRITE(6,820)(YB(I),I=1,LTIP)
199.000 820 FORMAT(2X,'YB:',10(2X,F10.5))
200.000 WRITE(6,821)(ZT(I),I=1,LTIP)

```

```

MU 0161
MU 0162
MU 0163
MU 0164
MU 0165
MU 0166
MU 0167
MU 0168
MU 0169
MU 0170
MU 0171
MU 0172
MU 0173
MU 0174
MU 0175
MU 0176
MU 0177
MU 0178
MU 0179
MU 0180
MU 0181
MU 0182
MU 0183
MU 0184
MU 0185
MU 0186
MU 0187
MU 0188
MU 0189
MU 0190
MU 0191
MU 0192
MU 0193
MU 0194
MU 0195
MU 0196
MU 0197
MU 0198
MU 0199
MU 0200

```

```

201.000 821 FORMAT(2X,'ZT:',10(2X,F10.5))
202.000      WRITE(6,822)(ZB(I),I=1,LTIP)
203.000 822 FORMAT(2X,'ZB:',10(2X,F10.5))
204.000      WRITE(6,90) (NCDE(I,1),I=1,LTIP)
205.000 90  FORMAT(/4X,'TOPOLOGY:',5X,10I12)
206.000      WRITE(6,92) (NGCE(I,2),I=1,LTIP)
207.000 92  FORMAT(4X,'NO. CF C.O.F.:',5X,10I12)
208.000      WRITE(6,67)((NDIS(I,J),J=1,5),I=1,LTIP)
209.000 67  FORMAT(/4X,'DEGREES OF FREEDOMS(1:MDF, 2:REST., 3:HINGED
210.000      &):',5X,1C(7X,5I1))
211.000 C
212.000 C
213.000 C
214.000 777 IF(LTIP.GT.2) GG TC 7777
215.000      CALL MATNUL(STF,10,10,10,10)
216.000      CALL MATNUL(ELMAS,10,10,50,50)
217.000      AREA=PR(IPRO)
218.000      CALL THETA(DCM,DCMT,RJ,ELEN,2)
219.000      IF(IPRINT.GE.7)WRITE(6,652)ELEN,PR(IPRO)
220.000 652 FORMAT(' LENGTH OF THE ELEMENT =',E16.6/
221.000      & , AREA OF THE ELEMENT =',E16.6)
222.000      IF(IPRINT.GE.10)WRITE(6,653)((DCM(I,J),I=1,3),J=1,3)
223.000 653 FORMAT(' DCM-MATRIX',/,3(5X,E16.6))
224.000      STF(3,3)=E(IPRO)*AREA/ELEN
225.000      STF(8,8)=STF(3,3)
226.000      STF(3,8)=-STF(3,3)
227.000      STF(8,3)=STF(3,8)
228.000      IF(IPRINT.GE.10)WRITE(6,651)((STF(I,J),I=1,10),J=1,10)
229.000 651 FORMAT(' STF-MATRIX',/,10(1X,E10.4))
230.000      CALL FORMT(TRAN,10,DCM,LTIP)
231.000      CALL TRNPCZ(TRAN,TRANT,10,10,10)
232.000      CALL FOICKF(STF,STF,TRAN,10,10,10,WS,50,2,C)
233.000      CALL FOICKF(STF,TRANT,STF,10,10,10,WS,50,3,0)
234.000      IF(IPRINT.GE.10)WRITE(6,654)((TRAN(I,J),I=1,1C),J=1,10)
235.000 654 FDMAT(' TRAN-MATRIX',/,10(1X,E10.4))
236.000      IF(IPRINT.GE.10)WRITE(6,651)((STF(I,J),I=1,10),J=1,10)
237.000      DO 655 I=1,10
238.000      DO 655 J=1,10
239.000 655 SS(I,J)=STF(I,J)
240.000      CALL REDUCE(ELMAS,SS,PMM,PSM,NDNC,MAX,LTIP)
MU 0201
MU 0202
MU 0203
MU 0204
MU 0205
MU 0206
MU 0207
MU 0208
MU 0209
MU 0210
MU 0211
MU 0212
MU 0213
MU 0214
MU 0215
MU 0216
MU 0217
MU 0218
MU 0219
MU 0220
MU 0221
MU 0222
MU 0223
MU 0224
MU 0225
MU 0226
MU 0227
MU 0228
MU 0229
MU 0230
MU 0231
MU 0232
MU 0233
MU 0234
MU 0235
MU 0236
MU 0237
MU 0238
MU 0239
MU 0240

```



```

281.000 107 CONTINUE MU 0281
282.000 DO 109 I=4,LT5,5 MU 0282
283.000 J=(I+1)/5 MU 0283
284.000 CALL VECI(VII,V2I,V3I,V3IL,J) MU 0284
285.000 PARA=SHODEN(XI,ETA,J,1,1,LTIP)*V3IL/2.0 MU 0285
286.000 DUDC(1,I,3)=PARA*VII(1) MU 0286
287.000 DUDC(2,I,3)=PARA*VII(2) MU 0287
288.000 DUDC(3,I,3)=PARA*VII(3) MU 0288
289.000 DUDC(1,I+1,3)=-{PARA*V2I(1)} MU 0289
290.000 DUDC(2,I+1,3)=-{PARA*V2I(2)} MU 0290
291.000 DUDC(3,I+1,3)=-{PARA*V2I(3)} MU 0291
292.000 109 CONTINUE MU 0292
293.000 C MU 0293
294.000 C MULTIPLY RJI BY DUDC MU 0294
295.000 C MU 0295
296.000 CALL JAKOB(RJ,RJI,V1I,V2I,V3I,DETJ,XI,ETA,ZETA,2,LTIP) MU 0296
297.000 DO 202 L=1,3 MU 0297
298.000 DO 202 I=1,3 MU 0298
299.000 DO 202 J=1,LT5 MU 0299
300.000 DUDG(L,J,I)=0.0 MU 0300
301.000 DO 202 K=1,3 MU 0301
302.000 DUDG(L,J,I)=DUDG(L,J,I)+RJI(I,K)*DUDC(L,J,K) MU 0302
303.000 C DUDC BECAME DERIV.OF DISP.W.R.T. CARTESIAN COORDINATES MU 0303
304.000 202 CONTINUE MU 0304
305.000 C CALL THETA(DCM,DCMT,RJ,V3L,1) MU 0305
306.000 C MULTIPLY DUDG BY DCM, STORE IN DUDC MU 0306
307.000 DO 226 L=1,3 MU 0307
308.000 DO 226 I=1,3 MU 0308
309.000 DO 226 J=1,LT5 MU 0309
310.000 DUDC(I,J,L)=0.0 MU 0310
311.000 DO 226 K=1,3 MU 0311
312.000 DUDC(I,J,L)=DUDC(I,J,L)+DCM(K,I)*DUDG(K,J,L) MU 0312
313.000 226 CONTINUE MU 0313
314.000 C MULTIPLY DCMT BY DUDC, STORE IN DUDG MU 0314
315.000 DO 227 L=1,3 MU 0315
316.000 DO 227 I=1,3 MU 0316
317.000 DO 227 J=1,LT5 MU 0317
318.000 DUDG(L,J,I)=0.0 MU 0318
319.000 DO 227 K=1,3 MU 0319
320.000 DUDG(L,J,I)=DUDG(L,J,I)+DUDC(L,J,K)*DCMT(I,K) MU 0320

```

```

321.000      227 CONTINUE
322.000      C
323.000      CCNSTRUCTION OF B MATRIX
324.000      C
325.000          DO 204 I=1,LT5
326.000              B(1,I)=DUDG(1,I,1)
327.000              B(2,I)=DUDG(2,I,2)
328.000              B(3,I)=DUDG(1,I,2)+DUDG(2,I,1)
329.000              B(4,I)=DUDG(3,I,1)+DUDG(1,I,3)
330.000              B(5,I)=DUDG(3,I,2)+DUDG(2,I,3)
331.000          204 CONTINUE
332.000      C      TRANSPOSE THE B MATRIX, FIND BT
333.000          CALL TRNPDZ(B,BT,5,50,5,50)
334.000      C
335.000      CCNSTRUCTION OF S MATRIX (S= BT*D*B )
336.000      C
337.000          CALL FOICKF(B,D,B,5,50,5,Z,5,3,0)
338.000          CALL FOICKF(S,BT,B,50,50,5,Z,5,1,C)
339.000          VLUM=VDLUM+HI*HJ*HM*DETJ
340.000      C
341.000      C      PERFORM THE INTEGRATIONS
342.000      C
343.000          DO 1111 I=1,LT5
344.000          DO 1111 J=1,LT5
345.000          SS(I,J)=SS(I,J)+HI*HJ*HM*DETJ*S(I,J)
346.000          1111 CONTINUE
347.000          WRITE(6,1C00) VLUM
348.000          VOL=0.000
349.000      C
350.000      CCNSTRUCTION OF THE MASS MATRIX
351.000      C
352.000          DC 3333 K3=1,2
353.000          DO 3333 K2=1,3
354.000          DC 3333 K1=1,3
355.000          CALL INTEGR(XI,ETA,ZETA,HI,HJ,HM,3,3,2,K1,K2,K3)
356.000      C
357.000      CCNSTRUCTION OF NBAR(UVW) FOR THE MASS MATRIX
358.000      C
359.000          CALL MATNUL(UVW,3,50,3,50)
360.000          DO 211 I=1,LT5,5
MU 0321
MU 0322
MU 0323
MU 0324
MU 0325
MU 0326
MU 0327
MU 0328
MU 0329
MU 0330
MU 0331
MU 0332
MU 0333
MU 0334
MU 0335
MU 0336
MU 0337
MU 0338
MU 0339
MU 0340
MU 0341
MU 0342
MU 0343
MU 0344
MU 0345
MU 0346
MU 0347
MU 0348
MU 0349
MU 0350
MU 0351
MU 0352
MU 0353
MU 0354
MU 0355
MU 0356
MU 0357
MU 0358
MU 0359
MU 0360

```

TEXT

```

361.000 J=(I+4)/5 MU 0361
362.000 UVW(1,I)=SHODEN(XI,ETA,J,1,1,LTIP) MU 0362
363.000 UVW(2,I+1)=UVW(1,I) MU 0363
364.000 UVW(3,I+2)=UVW(1,I) MU 0364
365.000 211 CONTINUE MU 0365
366.000 DO 212 I=4,LT5,5 MU 0366
367.000 J=(I+1)/5 MU 0367
368.000 CALL VECI(V1I,V2I,V3I,V3IL,J) MU 0368
369.000 PARA=SHODEN(XI,ETA,J,1,1,LTIP)*ZETA*V3IL/2.C MU 0369
370.000 UVW(1,I)=PARA*V1I(1) MU 0370
371.000 UVW(2,I)=PARA*V1I(2) MU 0371
372.000 UVW(3,I)=PARA*V1I(3) MU 0372
373.000 UVW(1,I+1)=-(PARA*V2I(1)) MU 0373
374.000 UVW(2,I+1)=-(PARA*V2I(2)) MU 0374
375.000 UVW(3,I+1)=-(PARA*V2I(3)) MU 0375
376.000 212 CONTINUE MU 0376
377.000 CALL JAKOB(RJ,RJI,V1I,V2I,V3I,DEIJ,XI,ETA,ZETA,1,LTIP) MU 0377
378.000 CALL TRNPOZ(UVW,UVWT,3,50,3,50) MU 0378
379.000 C MU 0379
380.000 VOL=VOL+HI*HJ*HM*DETJ MU 0380
381.000 CCNSTRUCTION OF EFEM MATRIX(EFEM=UVWT*UVW) MU 0381
382.000 C MU 0382
383.000 CALL FOICKF(EFEM,UVWT,UVW,50,50,3,Z,5,1,0) MU 0383
384.000 C PERFORM THE INTEGRATIONS MU 0384
385.000 DO 3333 I=1,LT5 MU 0385
386.000 DU 3333 J=1,LT5 MU 0386
387.000 ELMAS(I,J)=ELMAS(I,J)+HI*HJ*HM*DETJ*EFEM(I,J)*RC(IPRO) MU 0387
388.000 3333 CONTINUE MU 0388
389.000 WRITE(6,1000) VOL MU 0389
390.000 1000 FORMAT(5X,'VOLUME OF THE ELEMENT=',E16.6) MU 0390
391.000 IF(INVER(1).GT.0) CALL TRNSFR(ELMAS,SS,INVER,LTIP) MU 0391
392.000 GO TO 166 MU 0392
393.000 165 CALL MATCCP(DCM,RJ,3,3,3) MU 0393
394.000 CALL THETA(DCM,DCMT,RJI,V3L,2) MU 0394
395.000 CALL FOICKF(RJ,RJ,DCMT,3,3,3,WS,50,2,0) MU 0395
396.000 CALL FORMT(S,50,RJ,LTIP) MU 0396
397.000 CALL TRNPOZ(S,EFEM,50,50,50,50) MU 0397
398.000 CALL FOICKF(SS,SS,S,50,50,50,WS,50,2,C) MU 0398
399.000 CALL FOICKF(ELMAS,ELMAS,S,50,50,50,WS,50,2,0) MU 0399
400.000 CALL FOICKF(SS,EFEM,SS,50,50,50,WS,50,3,0) MU 0400

```

```

401.000      CALL FO1CKF(ELMAS,EFEM,ELMAS,50,50,50,WS,50,3,C)
402.000      CONTINUE
403.000      CALL REDUCE(ELMAS,SS,PMM,PSM,NDNO,MAX,LTIP)
404.000      IF(I1YP.NE.1) GO TO 2222
405.000      CALL THETA(DCM,DCMT,RJ,V3L,2)
406.000      CONTINUE
407.000      C
408.000      C
409.000      C
410.000      NCN=MBD-KAYVEC(NCNC)
411.000      WRITE(6,1282) MBD,NCN
412.000      1282 FORMAT(//4X,'ORIGINAL NO OF DEGREES OF FREEDOMS =',I5/
413.000      &      4X,'NO OF DEGREES OF FREEDOMS RETAINED =',I5)
414.000      C
415.000      C EIGEN SOLUTION
416.000      C
417.000      IF(IPRINT.GE.9) CALL MAPRIN(PSM,NCN,6)
418.000      CALL DEIGS(PSM,PMM,NCN,M1,M2,RCCT,VEC,IFAIL)
419.000      MDIM=NCN+NON*(NCN-1)/2
420.000      IF(IFAIL) GO TO 199
421.000      C
422.000      C
423.000      WRITE(6,1583)
424.000      1583 FORMAT(1H1,///20X,'***SQUARES OF THE NATUREL FREQUENCIES***')
425.000      GO TO 198
426.000      199 CALL MAPRIN(PSM,NCN,6)
427.000      STOP
428.000      198 NEV=M2-M1+1
429.000      DO 1 I=1,NEV
430.000      K=M2+1-I
431.000      WRITE(6,608)K,ROOT(I)
432.000      1 CONTINUE
433.000      608 FORMAT(25X,'EIGEN VALUE NO:',I2,'=',E15.6)
434.000      C
435.000      WRITE(6,17)
436.000      17 FORMAT(///)
437.000      PFR=E(I)*THIC**3/(12.0*(1.-PR(I)**2))
438.000      DO 607 I=1,NEV
439.000      ROOT(I)=DSQRT(DABS(RCCT(I)))
440.000      FREQ=ROOT(I)/6.28318531

```

```

MU 0401
MU 0402
MU 0403
MU 0404
MU 0405
MU 0406
MU 0407
MU 0408
MU 0409
MU 0410
MU 0411
MU 0412
MU 0413
MU 0414
MU 0415
MU 0416
MU 0417
MU 0418
MU 0419
MU 0420
MU 0421
MU 0422
MU 0423
MU 0424
MU 0425
MU 0426
MU 0427
MU 0428
MU 0429
MU 0430
MU 0431
MU 0432
MU 0433
MU 0434
MU 0435
MU 0436
MU 0437
MU 0438
MU 0439
MU 0440

```


TEXT

```

441.000 PHI=ROCT(I)/DSQRT(PFR/(RO(I)*THIC#RLLEN#4))
442.000 K=M2+1-I
443.000 WRITE(6,1975) K,ROCT(I),FREQ,PHI
444.000 1975 FORMAT(/10X,'FREQ.NO:',I2,5X,'CMEGA=',E15.6,5X,'FREQUENCY=',
445.000      &E15.6,5X,'PHI=',E15.6)
446.000 6C7 CONTINUE
447.000 WRITE(6,812)
448.000 812 FORMAT(1H1,///20X,'EIGEN VECTORS'//)
449.000 DO 810 I=1,NEV
450.000 IZA=NON*(I-1)+1
451.000 IZB=NON#I
452.000 K=M2+1-I
453.000 810 WRITE(6,811) K,(VEC(J),J=IZA,IZB)
454.000 811 FORMAT(5X,'VECTOR(',I2,')',/5(8X,E14.5))
455.000 STOP
456.000 END
457.000
458.000
459.000
460.000
461.000
462.000
463.000
464.000
465.000
466.000
467.000
468.000
469.000
470.000
471.000
472.000
473.000
474.000
475.000
476.000
477.000
478.000
479.000
480.000
C
C
C
C
FUNCTION DEFINITION
FUNCTION SHGDN(XI,ETA,I,N,NN,LTIP)
IMPLICIT REAL*8(A-H,O-Z)
IF(LTIP.EQ.10) GO TO(10,20,20,10,1,2,20,2,1,10),I
GO TO(10,20,20,10,30,20,30,10),I
10 XII=-1.0
GO TO 40
20 XII=1.0
GO TO 40
30 XII=0.0
GO TO 40
1 XII=-1.0/3.0
GO TO 40
2 XII=1.0/3.0
GO TO 40
40 IF(LTIP.EQ.10) GO TO(50,50,60,60,50,50,70,60,60,70),I
GO TO (50,60,60,50,70,60,70),I
50 ETAI=-1.0
GO TO 80
60 ETAI=1.0
GO TO 80
70 ETAI=0.0

```

LINE NUMBER

```

MU 0441
MU 0442
MU 0443
MU 0444
MU 0445
MU 0446
MU 0447
MU 0448
MU 0449
MU 0450
MU 0451
MU 0452
MU 0453
MU 0454
MU 0455
MU 0456
MU 0457
MU 0458
MU 0459
MU 0460
MU 0461
MU 0462
MU 0463
MU 0464
MU 0465
MU 0466
MU 0467
MU 0468
MU 0469
MU 0470
MU 0471
MU 0472
MU 0473
MU 0474
MU 0475
MU 0476
MU 0477
MU 0478
MU 0479
MU 0480

```

TEXT

481.000	80	XIC=XI*XI I	MU	0481
482.000		ETAO=ETA*ETAI	MU	0482
483.000		GO TO (85,95),N	MU	0483
484.000	85	IF(LTIP.EQ.10) GC TO(3,3,3,3,4,4,25,4,4,25),I	MU	0484
485.000		GO TO(90,90,90,90,15,25,15,25),I	MU	0485
486.000	90	SHODEN=0.25*(1.+ETAC)*{1.+XIC}*{XIC+ETAO-1.0}	MU	0486
487.000		RETURN	MU	0487
488.000	15	SHODEN=0.50*(1.0-XI**2)*(1.0+ETAO)	MU	0488
489.000		RETURN	MU	0489
490.000	25	SHODEN= 0.5*(1.0+XIO)*{1.0-ETA**2}	MU	0490
491.000		RETURN	MU	0491
492.000	3	SHODEN=1.0/32.0*(1.0+XIC)*{1.0+ETAO}*{8.0*ETAC-9.0* ε{1.0-XI**2)}	MU	0492
493.000		RETURN	MU	0493
494.000	4	SHODEN=9.0/32.0*(1.0+ETAO)*{1.0-XI**2}*{1.0+9.0*XIO}	MU	0494
495.000		RETURN	MU	0495
496.000	95	GO TO (96,11),NN	MU	0496
497.000	56	IF(LTIP.EQ.10) GO TO(5,5,5,5,6,6,55,6,6,55),I	MU	0497
498.000		GO TO (35,35,35,35,45,55,45,55),I	MU	0498
499.000	35	SHODEN=0.25*(1.0+ETAO)*XII*(2.*XIC+ETAO)	MU	0499
500.000		RETURN	MU	0500
501.000	45	SHODEN=-XI*(1.+ETAO)	MU	0501
502.000		RETURN	MU	0502
503.000	55	SHODEN=0.5*XII*(1.0-ETA**2)	MU	0503
504.000		RETURN	MU	0504
505.000	5	SHODEN=1.0/32.0*(1.0+ETAO)*{XII*(8.0*ETAC-9.0*(1.0-XI**2))+ ε18.0*XII*(1.0+XIC)}	MU	0505
506.000		RETURN	MU	0506
507.000	6	SHODEN=9.0/32.0*(1.0+ETAO)*{(-2.0*XI)*{1.0+9.0*XIO}+9*XII* ε (1.0-XI**2)}	MU	0507
508.000		RETURN	MU	0508
509.000	11	IF(LTIP.EQ.10) GO TO(7,7,7,7,8,8,41,8,8,41),I	MU	0509
510.000		GO TO(21,21,21,21,31,41,31,41),I	MU	0510
511.000	21	SHODEN=0.25*(1.0+XIO)*ETAI*(2.*ETAC+XIC)	MU	0511
512.000		RETURN	MU	0512
513.000	31	SHODEN=0.50*(1.0-XI**2)*ETAI	MU	0513
514.000		RETURN	MU	0514
515.000	41	SHODEN=-ETA*(1.0+XIO)	MU	0515
516.000		RETURN	MU	0516
517.000	7	SHODEN=1.0/32.0*(1.0+XIO)*{ETAI*(8.0*ETAC-9.0* ε	MU	0517
518.000			MU	0518
519.000			MU	0519
520.000			MU	0520

```

521.000      E(1.0-XI**2))+E.C*ETAI*(1.0+ETAC))
522.000      RETURN
523.000      8 SHODEN=9.0/32.0*(1.0-XI**2)*(1.0+9.0*XIO)*ETAI
524.000      RETURN
525.000      END
526.000      C
527.000      C
528.000      SUBROUTINE JAKOB(RJ,RJI,VII,V2I,V3I,DETI,XI,ETA,ZETA,KAYI,LTIP)
529.000      IMPLICIT REAL*8(A-H,O-Z)
530.000      COMMON/JAK/RJ3(3,3),RJ2(3,3),AINT(3)
531.000      COMMON/COR/XT(10),XB(10),YT(10),YE(10),ZT(10),ZB(10)
532.000      DIMENSION RJI(3,3),RJ(3,3)
533.000      DIMENSION VII(3),V2I(3),V3I(3)
534.000      CCNSTRUCTION GF JACCBIAN (RJ)
535.000      CALL MATNUL(RJ,3,3,3,3)
536.000      DO 102 NN=1,2
537.000      DO 102 I=1,LTIP
538.000      CALL VECI(VII,V2I,V3I,V3IL,I)
539.000      SHDER=SHODEN(XI,ETA,I,2,NN,LTIP)
540.000      PARSH=SHDER*ZETA/2.0*V3IL
541.000      RJ(NN,1)=RJ(NN,1)+SHDER*(XT(I)+XB(I))/2.0+PARSH*V3I(1)
542.000      RJ(NN,2)=RJ(NN,2)+SHDER*(YT(I)+YB(I))/2.0+PARSH*V3I(2)
543.000      RJ(NN,3)=RJ(NN,3)+SHDER*(ZT(I)+ZB(I))/2.0+PARSH*V3I(3)
544.000      102 CONTINUE
545.000      DO 103 I=1,LTIP
546.000      CALL VECI(VII,V2I,V3I,V3IL,I)
547.000      SHAPE=SHODEN(XI,ETA,I,1,LTIP)*C.5*V3IL
548.000      RJ(3,1)=RJ(3,1)+SHAPE*V3I(1)
549.000      RJ(3,2)=RJ(3,2)+SHAPE*V3I(2)
550.000      RJ(3,3)=RJ(3,3)+SHAPE*V3I(3)
551.000      103 CONTINUE
552.000      DO 1032 I=1,3
553.000      DO 1032 J=1,3
554.000      RJ2(I,J)=RJ(I,J)
555.000      RJ3(I,J)=RJ(I,J)
556.000      1032 CONTINUE
557.000      IF(KAYI.EQ.1) GO TO 100
558.000      C      INVERSION OF JACCBIAN (RJI)
559.000      NFAIL=0
560.000      CALL FOIAAF(RJ2,3,3,RJI,3,AINT,NFAIL)
MU 0521
MU 0522
MU 0523
MU 0524
MU 0525
MU 0526
MU 0527
MU 0528
MU 0529
MU 0530
MU 0531
MU 0532
MU 0533
MU 0534
MU 0535
MU 0536
MU 0537
MU 0538
MU 0539
MU 0540
MU 0541
MU 0542
MU 0543
MU 0544
MU 0545
MU 0546
MU 0547
MU 0548
MU 0549
MU 0550
MU 0551
MU 0552
MU 0553
MU 0554
MU 0555
MU 0556
MU 0557
MU 0558
MU 0559
MU 0560

```

TEXT

```

561.000 C DETERMINATION OF THE DETERMINANT CF THE JACOBIAN(DETJ) MU 0561
562.000 100 DETJ=RJ3(1,1)*(RJ3(2,2)*RJ3(3,3)-RJ3(3,2)*RJ3(2,3)) MU 0562
563.000 E -RJ3(1,2)*(RJ3(2,1)*RJ3(3,3)-RJ3(3,1)*RJ3(2,3)) MU 0563
564.000 E +RJ3(1,3)*(RJ3(2,1)*RJ3(3,2)-RJ3(3,1)*RJ3(2,2)) MU 0564
565.000 RETURN MU 0565
566.000 END MU 0566
567.000 C MU 0567
568.000 C BLOCK DATA MU 0568
569.000 IMPLICIT REAL*8(A-H,O-Z) MU 0569
570.000 COMMON/INTG/PNT4(4),PNT3(3),PNT2(2),H4(4),H3(3) MU 0570
571.000 COMMON/UNI/UNIT(3) MU 0571
572.000 DATA PNT4/-0.861136311594053,-C.329981043584856, MU 0572
573.000 E 0.339581043584856, 0.861136311594053/ MU 0573
574.000 DATA PNT3/-0.774596669241483,0.0E0,0.774596669241483/ MU 0574
575.000 DATA PNT2/-0.577350269189626,0.577350269189626/ MU 0575
576.000 DATA H3 /C.555555555555556,0.888888888888889,C.555555555555556/ MU 0576
577.000 DATA H4/0.347854845137454,0.652145154862546, MU 0577
578.000 E 0.652145154862546,0.347854845137454/ MU 0578
579.000 DATA UNIT/0.0D00,1.0D00,0.0D00/ MU 0579
580.000 END MU 0580
581.000 C MU 0581
582.000 C SUBROUTINE FOR THE DIRECTI CN CCSINES MU 0582
583.000 C MU 0583
584.000 C MU 0584
585.000 SUBROUTINE THETA(DCM,DCMT,RJ,V3L,NE) MU 0585
586.000 C NE=2 DIREC. COS. FOR ELEMENT MU 0586
587.000 IMPLICIT REAL*8(A-H,O-Z) MU 0587
588.000 COMMON/UNI/UNIT(3) MU 0588
589.000 COMMON/STHE/V1(3),V2(3),V3(3) MU 0589
590.000 COMMON/COR/XT(10),XB(10),YT(10),YB(10),ZT(10),ZB(10) MU 0590
591.000 DIMENSION RJ(3,3),DCM(3,3),DCMT(3,3) MU 0591
592.000 CCNSTRUCTION OF THE DIRECTI CN CCSINE MATRIX MU 0592
593.000 CCNSTRUCTING THE VECTOR V3 MU 0593
594.000 IF (NE.EQ.2) GO TO 1 MU 0594
595.000 V3(1)=RJ(1,2)*RJ(2,3)-RJ(2,2)*RJ(1,3) MU 0595
596.000 V3(2)=RJ(2,1)*RJ(1,3)-RJ(1,1)*RJ(2,3) MU 0596
597.000 V3(3)=RJ(1,1)*RJ(2,2)-RJ(2,1)*RJ(1,2) MU 0597
598.000 GO TO 2 MU 0598
599.000 1 V3(1)=(XT(2)+XB(2))/2.0-(YT(1)+XB(1))/2.0 MU 0599
600.000 V3(2)=(YT(2)+YB(2))/2.0-(Y1(1)+YB(1))/2.0 MU 0600

```

TEXT

```

601.000      V3(3)=(ZT(2)+ZB(2))/2.0-(ZT(1)+ZB(1))/2.0
602.000      CCNSTRUCT VECTOR V1
603.000      2 V1(1)=UNIT(2)*V3(3)-UNIT(3)*V3(2)
604.000      V1(2)=UNIT(3)*V3(1)-UNIT(1)*V3(3)
605.000      V1(3)=UNIT(1)*V3(2)-UNIT(2)*V3(1)
606.000      CCNSTRUCT V2 VECTOR
607.000      V2(1)=V3(2)*V1(3)-V3(3)*V1(2)
608.000      V2(2)=V3(3)*V1(1)-V3(1)*V1(3)
609.000      V2(3)=V3(1)*V1(2)-V3(2)*V1(1)
610.000      C      NORMALIZE THE VECTORS
611.000      V3L=DSQRT(V3(1)**2+V3(2)**2+V3(3)**2)
612.000      V1L=DSQRT(V1(1)**2+V1(2)**2+V1(3)**2)
613.000      V2L=DSQRT(V2(1)**2+V2(2)**2+V2(3)**2)
614.000      DO 220 I=1,3
615.000      V1(I)=V1(I)/V1L
616.000      V2(I)=V2(I)/V2L
617.000      V3(I)=V3(I)/V3L
618.000      220 CONTINUE
619.000      CCNSTRUCT THE DIRECTION COSINE MATRIX
620.000      C
621.000      DC 221 I=1,3
622.000      DCM(I,1)=V1(I)
623.000      DCM(I,2)=V2(I)
624.000      DCM(I,3)=V3(I)
625.000      221 CONTINUE
626.000      C      TRANSPOSE DCM MATRIX FIND DCMT
627.000      DO 225 I=1,3
628.000      DO 225 J=1,3
629.000      DCMT(I,J)=DCM(J,I)
630.000      225 CONTINUE
631.000      RETURN
632.000      END
633.000      C
634.000      C
635.000      SUBROUTINE VEC1(V1I,V2I,V3I,V3IL,I)
636.000      IMPLICIT REAL*8(A-H,O-Z)
637.000      COMMON/COR/XT(10),XB(10),YT(10),YB(10),ZT(10),ZB(10)
638.000      COMMON /UNI/UNIT(3)
639.000      DIMENSION VII(3),V2I(3),V3I(3)
640.000      CCNSTRUCT V3I VECTOR
MU 0601
MU 0602
MU 0603
MU 0604
MU 0605
MU 0606
MU 0607
MU 0608
MU 0609
MU 0610
MU 0611
MU 0612
MU 0613
MU 0614
MU 0615
MU 0616
MU 0617
MU 0618
MU 0619
MU 0620
MU 0621
MU 0622
MU 0623
MU 0624
MU 0625
MU 0626
MU 0627
MU 0628
MU 0629
MU 0630
MU 0631
MU 0632
MU 0633
MU 0634
MU 0635
MU 0636
MU 0637
MU 0638
MU 0639
MU 0640

```

TEXT

LINE NUMBER

```

641.000 V3I(1)=XT(I)-XB(I) MU 0641
642.000 V3I(2)=YT(I)-YB(I) MU 0642
643.000 V3I(3)=ZT(I)-ZB(I) MU 0643
644.000 CCONSTRUCT VECTOR V1I MU 0644
645.000 V1I(1)=UNIT(2)*V3I(3)-UNIT(3)*V3I(2) MU 0645
646.000 V1I(2)=UNIT(3)*V3I(1)-UNIT(1)*V3I(3) MU 0646
647.000 V1I(3)=UNIT(1)*V3I(2)-UNIT(2)*V3I(1) MU 0647
648.000 CCONSTRUCT V2I VECTOR MU 0648
649.000 V2I(1)=V3I(2)*V1I(3)-V3I(3)*V1I(2) MU 0649
650.000 V2I(2)=V3I(3)*V1I(1)-V3I(1)*V1I(3) MU 0650
651.000 V2I(3)=V3I(1)*V1I(2)-V3I(2)*V1I(1) MU 0651
652.000 C NORMALIZE THE VECTORS MU 0652
653.000 V3IL=DSQRT(V3I(1)**2+V3I(2)**2+V3I(3)**2) MU 0653
654.000 V1IL=DSQRT(V1I(1)**2+V1I(2)**2+V1I(3)**2) MU 0654
655.000 V2IL=DSQRT(V2I(1)**2+V2I(2)**2+V2I(3)**2) MU 0655
656.000 DO 1 K=1,3 MU 0656
657.000 V1I(K)=V1I(K)/V1IL MU 0657
658.000 V2I(K)=V2I(K)/V2IL MU 0658
659.000 V3I(K)=V3I(K)/V3IL MU 0659
660.000 1 CONTINUE MU 0660
661.000 RETURN MU 0661
662.000 END MU 0662
663.000 C MU 0663
664.000 C MATRIX PRINT SUBROUTINE MU 0664
665.000 SUBROUTINE MAPRIN(A,N,NPRIN) MU 0665
666.000 IMPLICIT REAL*8(A-H,O-Z) MU 0666
667.000 DIMENSION A(1) MU 0667
668.000 IF(NPRIN.EQ.3) GO TO 2 MU 0668
669.000 WRITE(NPRIN,1) N MU 0669
670.000 1 FORMAT(// 'MATRIX-DIMENSION=', I3//) MU 0670
671.000 2 M=1 MU 0671
672.000 DO 3 I=1,N MU 0672
673.000 M=M+I-1 MU 0673
674.000 MM=M+I-1 MU 0674
675.000 WRITE(NPRIN,5) I,(A(J),J=M,MM) MU 0675
676.000 3 CONTINUE MU 0676
677.000 5 FORMAT(2X, 'ROW NG:', I3/(9(2X,E12.5))) MU 0677
678.000 RETURN MU 0678
679.000 END MU 0679
680.000 C SUBROUTINE TO NULL A MATRIX MU 0680

```

```

661.000 C
662.000 SUBROUTINE MATNUL(A,M,N,MM,NN)
663.000 IMPLICIT REAL*8(A-H,O-Z)
664.000 DIMENSION A(MM,NN)
665.000 DO 1 I=1,M
666.000 DO 1 J=1,N
667.000 1 A(I,J)=0.000
668.000 RETURN
669.000 END
670.000 C
671.000 C SUB. TO PUT ZERO MATRIX INTO UNIT FORM
672.000 C
673.000 SUBROUTINE MATUNI(A,M,N)
674.000 IMPLICIT REAL*8(A-H,O-Z)
675.000 DIMENSION A(N,N)
676.000 DO 1 I=1,M
677.000 1 A(I,I)=1.000
678.000 RETURN
679.000 END
680.000 C
681.000 C SUB. TO FORM TRANSFORMATION MATRIX
682.000 C
683.000 SUBROUTINE FORMT(T,M,DCM,LTIP)
684.000 IMPLICIT REAL*8(A-H,O-Z)
685.000 DIMENSION T(M,M),DCM(3,3)
686.000 CALL MATNUL(T,M,M,M)
687.000 CALL MATUNI(T,M,M)
688.000 DO 1 I=1,LTIP
689.000 DO 1 J=1,3
690.000 K=(I-1)*5+J
691.000 DO 1 L=1,3
692.000 LL=(I-1)*5+L
693.000 1 T(K,LL)=DCM(J,L)
694.000 RETURN
695.000 END
696.000 C
697.000 C SUB. TO TRANSPOSE A MATRIX
698.000 C
699.000 SUBROUTINE TRNPOZ(A,B,M,N,MM,NN)
700.000 IMPLICIT REAL*8(A-H,O-Z)
701.000
702.000
703.000
704.000
705.000
706.000
707.000
708.000
709.000
710.000
711.000
712.000
713.000
714.000
715.000
716.000
717.000
718.000
719.000
720.000

```

MU 0681

MU 0682

MU 0683

MU 0684

MU 0685

MU 0686

MU 0687

MU 0688

MU 0689

MU 0690

MU 0691

MU 0692

MU 0693

MU 0694

MU 0695

MU 0696

MU 0697

MU 0698

MU 0699

MU 0700

MU 0701

MU 0702

MU 0703

MU 0704

MU 0705

MU 0706

MU 0707

MU 0708

MU 0709

MU 0710

MU 0711

MU 0712

MU 0713

MU 0714

MU 0715

MU 0716

MU 0717

MU 0718

MU 0719

MU 0720

TEXT

```

721.000 DIMENSION A(MM,NN),B(LA,MM) MU 0721
722.000 DO 1 I=1,M MU 0722
723.000 DO 1 J=1,N MU 0723
724.000 1 B(J,I)=A(I,J) MU 0724
725.000 RETURN MU 0725
726.000 END MU 0726
727.000 C MU 0727
728.000 C SUB. TO COPY A MATRIX MU 0728
729.000 SUBROUTINE MATCCP(A,B,M,N,M,M,NN) MU 0729
730.000 IMPLICIT REAL*8(A-H,O-Z) MU 0730
731.000 DIMENSION A(MM,NN), B(MM,NN) MU 0731
732.000 DO 1 I=1,M MU 0732
733.000 DO 1 J=1,N MU 0733
734.000 1 B(I,J)=A(I,J) MU 0734
735.000 RETURN MU 0735
736.000 END MU 0736
737.000 C MU 0737
738.000 C SUBROUTINE FOR INTEGRATION POINTS MU 0738
739.000 C MU 0739
740.000 SUBROUTINE INTEGR(XI,ETA,ZETA,HI,HJ,HM,INT1,INT2,INT3,I1,I2,I3) MU 0740
741.000 IMPLICIT REAL*8(A-H,O-Z) MU 0741
742.000 COMMON/INTG/PNT4(4),PNT3(3),PNT2(2),F4(4),H3(3) MU 0742
743.000 GO TO(1,2,3,4),INT1 MU 0743
744.000 1 XI=0.000 MU 0744
745.000 HI=2.000 MU 0745
746.000 GO TO 4 MU 0746
747.000 2 XI=PNT2(I1) MU 0747
748.000 HI=1.0000 MU 0748
749.000 GO TO 4 MU 0749
750.000 3 XI=PNT3(I1) MU 0750
751.000 HI=H3(I1) MU 0751
752.000 GO TO 4 MU 0752
753.000 40 XI=PNT4(I1) MU 0753
754.000 HI=H4(I1) MU 0754
755.000 4 GC TC (I1,I2,I3),INT2 MU 0755
756.000 11 ETA=0.000 MU 0756
757.000 HJ=2.0000 MU 0757
758.000 GO TO 14 MU 0758
759.000 12 ETA=PNT2(I2) MU 0759
760.000 HJ=1.0000 MU 0760

```



```

761.000      GO TO 14
762.000      13 ETA=PNT3(I2)
763.000      HJ=H3(I2)
764.000      14 GO TO(21,22,23),INT3
765.000      21 ZETA=0.0D0
766.000      HM=2.0D0
767.000      RETURN
768.000      22 ZETA=PNT2(I3)
769.000      HM=1.0D0D
770.000      RETURN
771.000      23 ZETA=PNT3(I3)
772.000      HM=H3(I3)
773.000      RETURN
774.000      END
775.000      C
776.000      C
777.000      C
778.000      SUBROUTINE FCR D MATRIX
779.000      SUBROUTINE CCNEL(D,I)
780.000      IMPLICIT REAL*8(A-H,O-Z)
781.000      COMMON/EPRO/E(5),RD(5),PR(5)
782.000      DIMENSION D(5,5)
783.000      CALL MATNUL(D,5,5,5,5)
784.000      C=E(I)/(1.0-PR(I)**2)
785.000      D(1,1)=C
786.000      D(2,2)=C
787.000      D(3,3)=(1.0-PR(I))/2.0*C
788.000      D(4,4)=(1.0-PR(I))/2.4*C
789.000      D(5,5)=D(4,4)
790.000      D(1,2)=C*PR(I)
791.000      D(2,1)=C*PR(I)
792.000      RETURN
793.000      END
794.000      C
795.000      C
796.000      SUBROUTINE REDUCE(ELMAS,SS,PSM,NCNC,MAX,LTIP)
797.000      IMPLICIT REAL*8(A-H,O-Z)
798.000      COMMON/RED/SLAVES(300),SLAVEM(300),NODE(10,2),NDIS(10,5),
799.000      &KAYVEC(500),ICCM(500)
800.000      DIMENSION ELMAS(50,50), SS(50,50), PMM(45150),PSM(45150)

```

```

MU 0761
MU 0762
MU 0763
MU 0764
MU 0765
MU 0766
MU 0767
MU 0768
MU 0769
MU 0770
MU 0771
MU 0772
MU 0773
MU 0774
MU 0775
MU 0776
MU 0777
MU 0778
MU 0779
MU 0780
MU 0781
MU 0782
MU 0783
MU 0784
MU 0785
MU 0786
MU 0787
MU 0788
MU 0789
MU 0790
MU 0791
MU 0792
MU 0793
MU 0794
MU 0795
MU 0796
MU 0797
MU 0798
MU 0799
MU 0800

```

TEXT

LINE NUMBER

```

C
801.000 ASSEMBLE THE SYSTEM MATRICES, INSERT THE BOUNDARY CONDITIONS
802.000 DO 403 I=1,LTIP
803.000 IF(NODE(I,2).EQ.0)GO TO 403
804.000 LCC1=NODE(I,1)
805.000 ILCP=LCC1*5-KAYVEC(LCC1)-NCDE(I,2)
806.000 ILCE=I*5-5
807.000 DO 404 J=1,LTIP
808.000 IF(NODE(J,2).EQ.0) GO TO 404
809.000 IF(NODE(J,1).LT.LCC1) GC TC 404
810.000 LCC2=NODE(J,1)
811.000 ILRP=LCC2*5-KAYVEC(LCC2)-NCDE(J,2)
812.000 ILRE=J*5-5
813.000 KS=NODE(J,2)
814.000 LS=NODE(I,2)
815.000 MK=0
816.000 DO 402 K=1,KS
817.000 IF(I.EQ.J) LS=K
818.000 MKF=MK+1
819.000 CO 8 MK=MKF,5
820.000 IF(NDIS(J,MK).NE.2) GC TO 9
821.000 8 CONTINUE
822.000 GO TO 404
823.000 9 ML=0
824.000 DO 402 L=1,LS
825.000 LCPM=ILCP+L
826.000 LRPJ=ILRP+K
827.000 LIVE=LCPM+LRPM*(LRPM-1)/2
828.000 MLF=ML+1
829.000 DO 5 ML=MLF,5
830.000 IF(NDIS(I,ML).NE.2) GO TO 6
831.000 5 CONTINUE
832.000 GO TO 403
833.000 6 LCEM=ILCE+ML
834.000 LREM=ILRE+MK
835.000 PMM(LIVE)=PMM(LIVE)+ELMAS(LREM,LCEM)
836.000 PSM(LIVE)=PSM(LIVE)+SS(LREM,LCEM)
837.000 402 CONTINUE
838.000 404 CONTINUE
839.000 403 CONTINUE
840.000 C DETERMINE THE INSTANTANEOUS SIZE OF SYSTEM MATRICES
MU CEC1
MU 0802
MU 0803
MU 0804
MU 0805
MU 0806
MU 0807
MU 0808
MU 0809
MU 0810
MU 0811
MU 0812
MU 0813
MU 0814
MU 0815
MU 0816
MU 0817
MU 0818
MU 0819
MU 0820
MU 0821
MU 0822
MU 0823
MU 0824
MU 0825
MU 0826
MU 0827
MU 0828
MU 0829
MU 0830
MU 0831
MU 0832
MU 0833
MU 0834
MU 0835
MU 0836
MU 0837
MU 0838
MU 0839
MU 0840

```

```

841.000 DO 100 IL=1,LTIP MU 0841
842.000 IF(NODE(IL,2).EQ.0) GO TO 100 MU 0842
843.000 IF(NGDE(IL,1).GT.MAX) MAX=NODE(IL,1) MU 0843
844.000 100 CONTINUE MU 0844
845.000 LIM1=5*MAX-KAYVEC(MAX) MU 0845
846.000 C MU 0846
847.000 C REDUCTION OF SCME VARIABLES MU 0847
848.000 C MU 0848
849.000 DO 10 I=1,LTIP MU 0849
850.000 LIM=5*MAX-KAYVEC(MAX) MU 0850
851.000 IF(NODE(I,2).EQ.0) GO TO 10 MU 0851
852.000 LOC=NUDE(I,1) MU 0852
853.000 IN1=ICCM(LOC)/10 MU 0853
854.000 IN2=ICCM(LCC)-10*IN1 MU 0854
855.000 IF(IN2) 19,17,10 MU 0855
856.000 19 IF(IN2+2) 10,21,22 MU 0856
857.000 22 IF(IN1.NE.-2)GC TO 10 MU 0857
858.000 21 DO 18 MEM=1,5 MU 0858
859.000 IF(NDIS(I,MEM).EQ.0)NDIS(I,MEM)=1 MU 0859
860.000 18 CONTINUE MU 0860
861.000 17 IKCL=LOC#5-KAYVEC(LOC) MU 0861
862.000 IND=0 MU 0862
863.000 JSON=NODE(I,2) MU 0863
864.000 JAN=0 MU 0864
865.000 DO 20 J=1,JSON MU 0865
866.000 JAN=JAN+1 MU 0866
867.000 16 IF(NDIS(I,JAN).EQ.1)GC TO 20 MU 0867
868.000 IF(NDIS(I,JAN).NE.2) GO TO 15 MU 0868
869.000 JAN=JAN+1 MU 0869
870.000 GO TO 16 MU 0870
871.000 15 KCLE=IKOL-JSON+J-IND MU 0871
872.000 ISAV=KOLE#(KCLE-1)/2 MU 0872
873.000 PIVGT=PSM(KOLE+ISAV) MU 0873
874.000 IF(PIVOT.EQ.0) GO TO 32 MU 0874
875.000 GO TO 31 MU 0875
876.000 32 WRITE(6,30) KCLE,NODE(I,1),KAYVEC(LOC),I,J MU 0876
877.000 30 FORMAT(/5X,'PIVOT FOR THE SLAVE',I5,'IS ZERC',5X,'NODE=',I5, MU 0877
878.000 &,'KAYMA=',I5,'I=',I5,3X,'J=',I5) MU 0878
879.000 STOP MU 0879
880.000 31 IF(LIM-KOLE) 101,101,102 MU 0880

```

TEXT

881.000	102	LSON=KCLE	MU	0881
882.000		GC TO 103	MU	0882
883.000	101	LSON=LIM	MU	0883
884.000	103	DC 1 L=1,LSON	MU	0884
885.000		LIVE=L+ISAV	MU	0885
886.000		SLAVES(L)=PSM(LIVE)	MU	0886
887.000		SLAVEM(L)=PMM(LIVE)	MU	0887
888.000		PSM(LIVE)=0.000	MU	0888
889.000		PMM(LIVE)=0.000	MU	0889
890.000	1	CONTINUE	MU	0890
891.000		IF(LSON.EQ.LIM) GO TO 104	MU	0891
892.000		ILK=LSCN+1	MU	0892
893.000		DO 4 L=ILK,LIM	MU	0893
894.000		LIVE=KOLE+L*(L-1)/2	MU	0894
895.000		SLAVES(L)=PSM(LIVE)	MU	0895
896.000		SLAVEM(L)=PMM(LIVE)	MU	0896
897.000		PSM(LIVE)=0.000	MU	0897
898.000		PMM(LIVE)=0.000	MU	0898
899.000	4	CONTINUE	MU	0899
900.000	104	IF(LIM-KOLE) 105,106,106	MU	0900
901.000	106	KSCN=KOLE-1	MU	0901
902.000		GO TO 107	MU	0902
903.000	105	KSON=LIM	MU	0903
904.000	107	DO 2 K=1,KSON	MU	0904
905.000		IPK=K*(K-1)/2	MU	0905
906.000		PARI=SLAVES(K)/PIVCT	MU	0906
907.000		DO 3 L=1,K	MU	0907
908.000		LIVE=L+IPK	MU	0908
909.000		PARA=SLAVES(L)*PARI	MU	0909
910.000		PMM(LIVE)=PMM(LIVE)-{SLAVEM(K)*SLAVES(L)	MU	0910
911.000	8	+SLAVEM(L)*SLAVES(K)-SLAVEM(KCLE)*PARA}/PIVCT	MU	0911
912.000		PSM(LIVE)=PSM(LIVE)-PARA	MU	0912
913.000	3	CONTINUE	MU	0913
914.000	2	CONTINUE	MU	0914
915.000		IF(LIM-KOLE) 108,108,201	MU	0915
916.000	201	KILK=KCLE+1	MU	0916
917.000		KSCN=LIM	MU	0917
918.000		IF(IN2.EQ.0) GO TO 202	MU	0918
919.000		DO 212 K=KILK,KSCN	MU	0919
920.000		IPK=K*(K-1)/2	MU	0920

921.000	PARI=SLAVES(K)/PIVCT	MU 0921
922.000	LSON=KOLE-1	MU 0922
923.000	DC 213 L=1,LSON	MU 0923
924.000	LIVE=L+IPK	MU 0924
925.000	PARA=SLAVES(L)*PARI	MU 0925
926.000	PMM(LIVE)=PMM(LIVE)-(SLAVEM(K)*SLAVES(L)	MU 0926
927.000	+SLAVEM(L)*SLAVES(K)-SLAVEM(KCLE)*PARA)/PIVCT	MU 0927
928.000	&	MU 0928
929.000	213 CONTINUE	MU 0929
930.000	DC 223 L=KILK,K	MU 0930
931.000	LIVE=L+IPK	MU 0931
932.000	PARA=SLAVES(L)*PARI	MU 0932
933.000	PMM(LIVE)=PMM(LIVE)-(SLAVEM(K)*SLAVES(L)	MU 0933
934.000	+SLAVEM(L)*SLAVES(K)-SLAVEM(KOLE)*PARA)/PIVCT	MU 0934
935.000	&	MU 0935
936.000	PSM(LIVE)=PSM(LIVE)-PARA	MU 0936
937.000	223 CONTINUE	MU 0937
938.000	212 CONTINUE	MU 0938
939.000	GO TO 108	MU 0939
940.000	202 DO 112 K=KILK,KSON	MU 0940
941.000	KK=K-1	MU 0941
942.000	IPK=K*(K-1)/2	MU 0942
943.000	IPKK=KK*(KK-1)/2	MU 0943
944.000	PARI=SLAVES(K)/PIVCT	MU 0944
945.000	LSON=KOLE-1	MU 0945
946.000	DO 113 L=1,LSON	MU 0946
947.000	LIVE1=L+IPKK	MU 0947
948.000	LIVE2=L+IPK	MU 0948
949.000	PARA=SLAVES(L)*PARI	MU 0949
950.000	PMM(LIVE1)=PMM(LIVE2)-(SLAVEM(K)*SLAVES(L)	MU 0950
951.000	+SLAVEM(L)*SLAVES(K)-SLAVEM(KCLE)*PARA)/PIVCT	MU 0951
952.000	&	MU 0952
953.000	PSM(LIVE1)=PSM(LIVE2)-PARA	MU 0953
954.000	113 CONTINUE	MU 0954
955.000	DO 123 L=KILK,K	MU 0955
956.000	LL=L-1	MU 0956
957.000	LIVE1=LL+IPKK	MU 0957
958.000	LIVE2=L+IPK	MU 0958
959.000	PARA=SLAVES(L)*PARI	MU 0959
960.000	PMM(LIVE1)=PMM(LIVE2)-(SLAVEM(K)*SLAVES(L)	MU 0960
	+SLAVEM(L)*SLAVES(K)-SLAVEM(KOLE)*PARA)/PIVCT	
	&	
	PSM(LIVE1)=PSM(LIVE2)-PARA	

TEXT

961.000	123	CONTINUE	MU 0961
962.000	112	CONTINUE	MU 0962
963.000	108	IF(IN2.NE.0) GO TO 20	MU 0963
964.000		IND=IND+1	MU 0964
965.000		LIM=LIM-1	MU 0965
966.000	20	CONTINUE	MU 0966
967.000		DO 11 M=1,NONC	MU 0967
968.000		IF(M.LT.NCDE(I,1)) GO TO 11	MU 0968
969.000		KAYVEC(M)=KAYVEC(M)+IND	MU 0969
970.000	11	CONTINUE	MU 0970
971.000	10	CONTINUE	MU 0971
972.000		WRITE(6,50) LIM1,LIM	MU 0972
973.000	50	FORMAT('//5X,' INSTANTANEOUS SIZE OF THE SYSTEM MATRICES =',	MU 0973
974.000		& 15,'/5X,' SIZE OF THE SYSTEM MATRICES AFTER REDUCTION=',15//)	MU 0974
975.000		LIM=LIM+1	MU 0975
976.000		IF(LIM.GT.LIM1) GO TO 52	MU 0976
977.000		DO 51 MM=LIM,LIM1	MU 0977
978.000		DO 51 MN=1,MM	MU 0978
979.000		LIVE3=MN+MM*(MM-1)/2	MU 0979
980.000		PMM(LIVE3)=0.C	MU 0980
981.000		PSM(LIVE3)=0.0	MU 0981
982.000	51	CONTINUE	MU 0982
983.000	52	CONTINUE	MU 0983
984.000		RETURN	MU 0984
985.000		END	MU 0985
986.000			MU 0986
987.000		DATA GENERATION SUBROUTINE	MU 0987
988.000			MU 0988
989.000			MU 0989
990.000		MAIN GENDAT	MU 0990
991.000		SUBROUTINE GENOD(UZUN,NCDE,NC,NMC,NEY)	MU 0991
992.000		IMPLICIT REAL*8(A-H,O-Z)	MU 0992
993.000		COMMON/GEN/ WS(50),IPT(50),IBC1,IBC2,IS1,IS2,MIS1,MIS2	MU 0993
994.000		DIMENSION NODE(10,2),ISAR(4)	MU 0994
995.000		READ(5,3) NCPAR,INC,IO,IR,IBC1,IBC2,IS2,MIS1,MIS2,ISAR	MU 0995
996.000		READ(5,35) TOTAN, (WS(I),I=1,IO)	MU 0996
997.000		PI=3.141592653589793	MU 0997
998.000		TANIN=TOTAN/(NEY*180)*PI	MU 0998
999.000		ISI=0	MU 0999
1000.000		ZINC=UZUN/NEY	MU 1000
		LEVEL=NEY+1	

```

1,001.000 L=LEVEL+NEY MU 1001
1,002.000 DO 1 J=1,NCPAR MU 1002
1,003.000 READ(5,7) INDE MU 1003
1,004.000 7 FORMAT(I5) MU 1004
1,005.000 GO TO (51,52,53),INDE MU 1005
1,006.000 51 READ(5,56) XT1,YTI,XBI,YBI,XTF,YTF,XBF,YBF,NOPT,NTERS MU 1006
1,007.000 56 FORMAT(8E10.4,/2I5) MU 1007
1,008.000 M1=0 MU 1008
1,009.000 M2=NOPT-1 MU 1009
1,010.000 GO TO 55 MU 1010
1,011.000 52 READ(5,4) XC,YC,RAD,THIC,AINI,AFIN,NOPT,NTERS MU 1011
1,012.000 AINI=AINI/180.0*PI MU 1012
1,013.000 AFIN=AFIN/180.0*PI MU 1013
1,014.000 AINC={AFIN-AINI}/(NOPT-1) MU 1014
1,015.000 A=AINI MU 1015
1,016.000 GO TO 55 MU 1016
1,017.000 53 READ(5,58)NOPT,NTERS MU 1017
1,018.000 58 FORMAT(2I5) MU 1018
1,019.000 55 READ(5,5) NCDE MU 1019
1,020.000 DO 1 I=1,NOPT MU 1020
1,021.000 GO TO (61,62,63) ,INDE MU 1021
1,022.000 61 XATI=(M1*XTF+M2*XTI)/FLOAT(M1+M2) MU 1022
1,023.000 XABI=(M1*XBF+M2*XBI)/FLOAT(M1+M2) MU 1023
1,024.000 YATI=(M1*YTF+M2*YTI)/FLCAT(M1+M2) MU 1024
1,025.000 YABI=(M1*YBF+M2*YBI)/FLCAT(M1+M2) MU 1025
1,026.000 GO TO 66 MU 1026
1,027.000 62 XATI=XC+RAD*DCCS(A) MU 1027
1,028.000 XABI=XC+{RAD-THIC}*DCCS(A) MU 1028
1,029.000 YATI=YC+RAD*DSIN(A) MU 1029
1,030.000 YABI=YC+{RAD-THIC}*DSIN(A) MU 1030
1,031.000 GO TO 66 MU 1031
1,032.000 63 READ(5,67)XATI,XABI,YATI,YABI MU 1032
1,033.000 67 FORMAT(4E10.4) MU 1033
1,034.000 66 DO 41 N=1,4 MU 1034
1,035.000 IF(ISAR(N).EQ.NODE(I,1)) GO TO 43 MU 1035
1,036.000 41 CONTINUE MU 1036
1,037.000 GO TO 42 MU 1037
1,038.000 43 XATI=(XATI+XABI)/2.0 MU 1038
1,039.000 XABI=XATI MU 1039
1,040.000 42 IF(NODE(I,2)) 10,20,30 MU 1040

```

LINE NUMBER

TEXT

PAGE

27

```

1,041.000 10 IF(NODE(I,2).LT.-1) GC TC 30
1,042.000 IF(NTERS.LT.0) WRITE(6,11)NDDF(I,1),XABI,XATI,YABI,YATI
1,043.000 IF(NTERS.GE.0) WRITE(6,11) NCDE(I,1),XATI,XABI,YATI,YABI
1,044.000 GC TO 2
1,045.000 20 DO 21 K=1,LEVEL
1,046.000 Z=ZINC*(K-1)
1,047.000 NODNO=NODE(I,1)+(K-1)*INC
1,048.000 CALL DOF(Z,NODNO,NNC, ID, IR, NDOF, IDIS)
1,049.000 TANG=(K-1)*TANIN
1,050.000 XAT=XATI*DCOS(TANG)-YATI*DSIN(TANG)
1,051.000 XAB=XABI*DCOS(TANG)-YABI*DSIN(TANG)
1,052.000 YAT=XATI*DSIN(TANG)+YATI*DCOS(TANG)
1,053.000 YAB=XABI*DSIN(TANG)+YABI*DCOS(TANG)
1,054.000 LINE=(NODNO+NG)*1000
1,055.000 IF(NTERS.LT.0)WRITE(7,LINE,22)NDOF, IDIS, XAB, XAT, Z, Z, YAB, YAT
1,056.000 IF(NTERS.GE.0) WRITE(7,LINE,22) NDOF, IDIS, XAT, XAB, Z, Z, YAT, YAB
1,057.000 21 CONTINUE
1,058.000 GO TO 2
1,059.000 30 DO 31 K=1,L
1,060.000 Z=ZINC/2.00*(K-1)
1,061.000 KK=2
1,062.000 IF(K/2*2.NE.K) KK=1
1,063.000 NODNO=IABS(NODE(I, KK))+(K-KK)*INC/2
1,064.000 IF(INC.NE.1) GO TO 40
1,065.000 NCDNO=NODE(I,1)+K-1
1,066.000 40 IF(NODE(I,2).LT.-1)IS1=99
1,067.000 CALL DOF(Z,NODNO,NNC, ID, IR, NDOF, IDIS)
1,068.000 TANG=(K-1)*TANIN/2.0
1,069.000 XAT=XATI*DCOS(TANG)-YATI*DSIN(TANG)
1,070.000 XAB=XABI*DCOS(TANG)-YABI*DSIN(TANG)
1,071.000 YAT=XATI*DSIN(TANG)+YATI*DCOS(TANG)
1,072.000 YAB=XABI*DSIN(TANG)+YABI*DCOS(TANG)
1,073.000 LINE=(NODNO+NG)*1000
1,074.000 IF(NTERS.LT.0)WRITE(7,LINE,22)NDOF, IDIS, XAB, XAT, Z, Z, YAB, YAT
1,075.000 IF(NTERS.GE.0) WRITE(7,LINE,22) NDOF, IDIS, XAT, XAB, Z, Z, YAT, YAB
1,076.000 31 CONTINUE
1,077.000 2 GO TO(71,72,1),INDE
1,078.000 71 M1=M1+1
1,079.000 M2=M2-1
1,080.000 GO TO 1
MU 1041
MU 1042
MU 1043
MU 1044
MU 1045
MU 1046
MU 1047
MU 1048
MU 1049
MU 1050
MU 1051
MU 1052
MU 1053
MU 1054
MU 1055
MU 1056
MU 1057
MU 1058
MU 1059
MU 1060
MU 1061
MU 1062
MU 1063
MU 1064
MU 1065
MU 1066
MU 1067
MU 1068
MU 1069
MU 1070
MU 1071
MU 1072
MU 1073
MU 1074
MU 1075
MU 1076
MU 1077
MU 1078
MU 1079
MU 1080

```


TEXT

```

72 A=A+AINC MU 1081
1 CONTINUE MU 1082
3 FORMAT(13I5) MU 1083
4 FORMAT(6E10.4,2I5) MU 1084
5 FORMAT(10I5) MU 1085
11 FORMAT(5X,'NODE:',I3,3X,'CO-OR.=',4E12.4) MU 1086
22 FORMAT(2I5,6E17.9) MU 1087
35 FORMAT(8E10.4) MU 1088
RETURN MU 1089
END MU 1090
C MU 1091
C MU 1092
SUBROUTINE DOF(Z,NCENO,NNC,LD,IR,NDOF,IDIS) MU 1093
IMPLICIT REAL*8(A-H,O-Z) MU 1094
COMMON/GEN/ WS(50),IPT(50),IBC1,IBC2,IS1,IS2,MIS1,MIS2 MU 1095
IDIS=0 MU 1096
NDOF=5 MU 1097
DO 1 I=1,NNC MU 1098
IF(IPT(I).EQ.NODNO) GO TO 3 MU 1099
1 CONTINUE MU 1100
IF(IS1.NE.95) GO TO 4 MU 1101
IS1=0 MU 1102
IF(MIS1.NE.0) GO TO 7 MU 1103
IDIS=IS2 MU 1104
4 DO 2 J=1,LD MU 1105
IF(DABS(WS(I)-Z).LE.0.01) GO TO 5 MU 1106
2 CONTINUE MU 1107
RETURN MU 1108
3 NDOF=IBC1 MU 1109
IDIS=IBC2 MU 1110
RETURN MU 1111
5 IF(NDOF.EQ.MIS1) GO TO 8 MU 1112
IF(IDIS.EQ.IS2) GO TO 6 MU 1113
IDIS=IR MU 1114
RETURN MU 1115
6 IDIS=IS2+IR MU 1116
RETURN MU 1117
7 NDOF=MIS1 MU 1118
IDIS=MIS2 MU 1119
GO TO 4 MU 1120

```

LINE NUMBER

TEXT

PAGE 29

```

1,121.000      8 IDIS=IDIS+IR
1,122.000      IF(IDIS/10000.EQ.3)IDIS=IDIS-10000
1,123.000      IF((IDIS/1000-IDIS/10000*10).EQ.3)IDIS=IDIS-1000
1,124.000      IF((IDIS/100-IDIS/1000*10).EQ.3)IDIS=IDIS-100
1,125.000      IF((IDIS/10-IDIS/100*10).EQ.3)IDIS=IDIS-10
1,126.000      IF((IDIS-IDIS/10*10).EQ.3)IDIS=IDIS-1
1,127.000      RETURN
1,128.000      END
1,129.000      C
1,130.000      C
1,131.000      SUBROUTINE GENTOP(NODE,ICCM,NC,NEY,ISTEK)
1,132.000      IMPLICIT REAL*8(A-H,O-Z)
1,133.000      DIMENSION NODE(10,2),ICCM(500),INVER(4)
1,134.000      DO 4 I=1,NO
1,135.000      READ(5,5) (NODE(J,1),J=1,10),INTEG,LTIP,ITYP,IIPC,INVER
1,136.000      INCI=0
1,137.000      INC2=0
1,138.000      GO 4 K=1,NEY
1,139.000      NGDE(1,1)=NGDE(1,1)+INCI
1,140.000      NODE(2,1)=NODE(2,1)+INCI
1,141.000      C
1,142.000      IF(LTIP.EQ.2)GC TC 2
1,143.000      C
1,144.000      NODE(3,1)=NGDE(3,1)+INCI
1,145.000      NODE(4,1)=NODE(4,1)+INCI
1,146.000      NODE(5,1)=NGDE(5,1)+INC2
1,147.000      IF(LTIP.EQ.10)GO TO 1
1,148.000      NGDE(6,1)=NGDE(6,1)+INCI
1,149.000      NODE(7,1)=NGDE(7,1)+INC2
1,150.000      NODE(8,1)=NGDE(8,1)+INCI
1,151.000      GO TO 2
1,152.000      1 NODE(6,1)=NGDE(6,1)+INC2
1,153.000      NODE(7,1)=NGDE(7,1)+INCI
1,154.000      NODE(8,1)=NGDE(8,1)+INC2
1,155.000      NGDE(9,1)=NGDE(9,1)+INC2
1,156.000      NODE(10,1)=NGDE(10,1)+INCI
1,157.000      2 DO 3 M=1,LTIP
1,158.000      LOC=NODE(M,1)
1,159.000      3 ICCM(LOC)=ICCM(LOC)+1
1,160.000      C
MU 1121
MU 1122
MU 1123
MU 1124
MU 1125
MU 1126
MU 1127
MU 1128
MU 1129
MU 1130
MU 1131
MU 1132
MU 1133
MU 1134
MU 1135
MU 1136
MU 1137
MU 1138
MU 1139
MU 1140
MU 1141
MU 1142
MU 1143
MU 1144
MU 1145
MU 1146
MU 1147
MU 1148
MU 1149
MU 1150
MU 1151
MU 1152
MU 1153
MU 1154
MU 1155
MU 1156
MU 1157
MU 1158
MU 1159
MU 1160

```

LINE NUMBER

TEXT

PAGE 30

```

1,161.000 C
1,162.000 IF(LTIP.NE.2) GO TC 20
1,163.000 INCI=INTEG
1,164.000 IF(INCI.LE.2) GO TC 6
1,165.000 GO TO 21
1,166.000 C
1,167.000 C
1,168.000 20 INCI=NODE(4,1)-NODE(1,1)
1,169.000 INC2=NODE(7,1)-NODE(5,1)
1,170.000 IF(LTIP.EQ.10) INC2=NODE(9,1)-NODE(5,1)
1,171.000 IF(INCI.NE.INC2) GOTO 6
1,172.000 21 LINE=(I+(K-1)*NO)*1000
1,173.000 C
1,174.000 IF(LTIP.EQ.2) GO TC 4
1,175.000 C
1,176.000 IF (K.GT.1) GO TO 11
1,177.000 GO TO 4
1,178.000 11 IF(ISTEK-5) 4,13,14
1,179.000 13 ITYP=2
1,180.000 GO TO 4
1,181.000 14 ITYP=3
1,182.000 GO TO 4
1,183.000 6 LINE=(K+(I-1)*NEY)*1000
1,184.000 IF(K.GT.1) ITYP=2
1,185.000 4 WRITE(7,LINE,5)(NODE(J,1),J=1,10),INTEG,LTIP,ITYP,IPRO,INVER
1,186.000 5 FORMAT(14I5,4I2)
1,187.000 NO=NO#NEY
1,188.000 RETURN
1,189.000 END
1,190.000 C
1,191.000 C
1,192.000 C
1,193.000 C
1,194.000 C
1,195.000 C
1,196.000 SUBROUTINE TRANSR(ELMAS,SS,INVER,LTIP)
1,197.000 IMPLICIT REAL*8(A-H,C-Z)
1,198.000 DIMENSION ELMAS(50,50),SS(50,50),INVER(4)
1,199.000 LT5=5#LTIP
1,200.000 DO 1 I=1,LTIP

```

LINE NUMBER

TEXT

PAGE

1,201.000	DO 2 K=1,4	MU 1201
1,202.000	IF(I.EQ.INVER(K)) GO TO 3	MU 1202
1,203.000	2 CONTINUE	MU 1203
1,204.000	GO TO 1	MU 1204
1,205.000	3 I5=I#5	MU 1205
1,206.000	DO 4 J=1,LT5	MU 1206
1,207.000	ELMAS(I5,J)=--ELMAS(I5,J)	MU 1207
1,208.000	ELMAS(J,I5)=--ELMAS(J,I5)	MU 1208
1,209.000	SS(I5,J)=--SS(I5,J)	MU 1209
1,210.000	SS(J,I5)=--SS(J,I5)	MU 1210
1,211.000	4 CONTINUE	MU 1211
1,212.000	1 CONTINUE	MU 1212
1,213.000	RETURN	MU 1213
1,214.000	END	MU 1214

